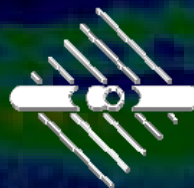


ICPEAC XXX
Cairns, Australia
2017-07-25

Tutorial:
Shaping atoms and molecules
with **strong** laser fields

Outline

- Introduction
(what is the shape of atoms?)
- What happens to atoms in strong laser fields?
(besides ionization and recollision)
- How to experimentally measure shape-changing atoms?
(methods viable for neutral species)
- What do we learn?
- Where do we go with it?



Thomas Pfeifer
MPIK Heidelberg, Germany

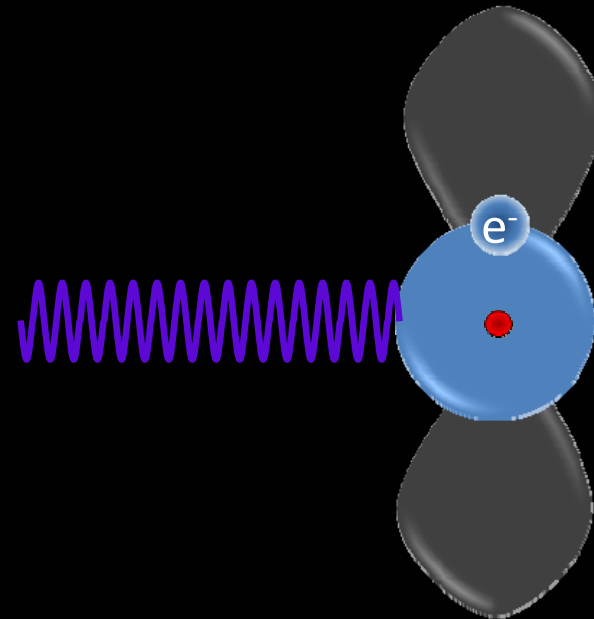
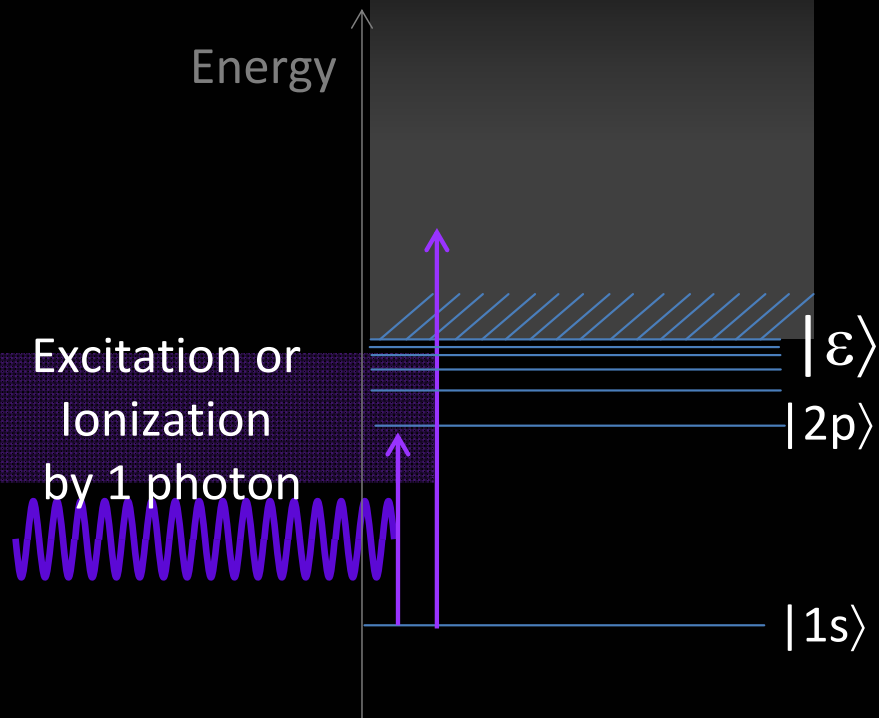
The “shape” of atoms ...

... in the Energy domain: The Spectrum

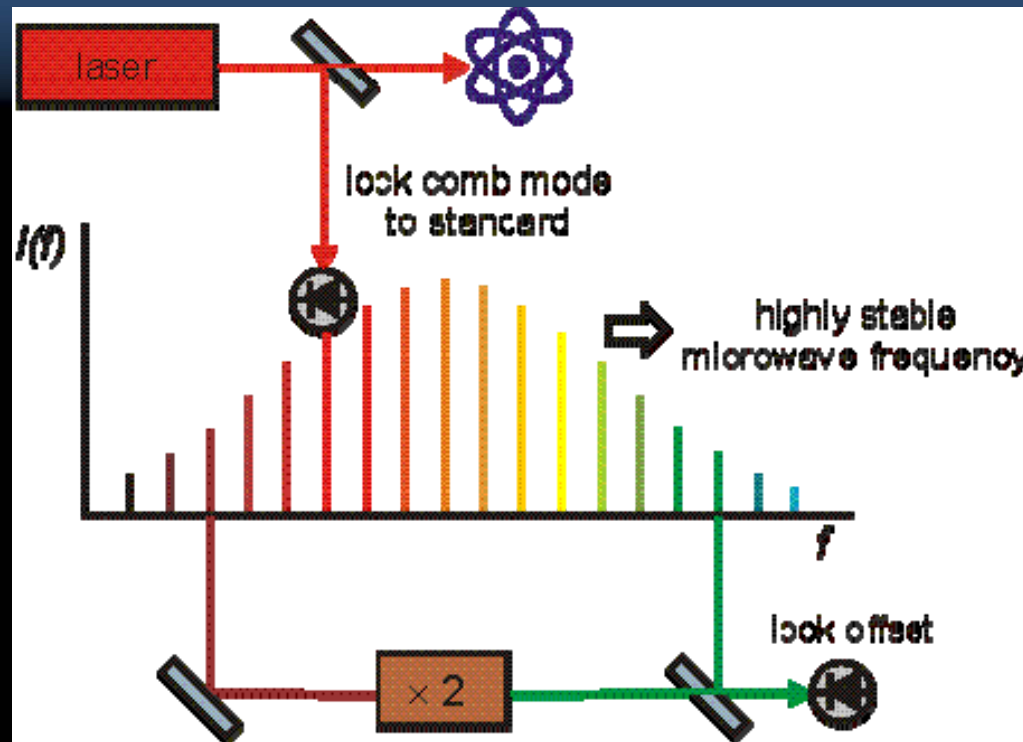
... in real space: Wavefunction/Orbital

... and their interaction with light

weak ...



Atomic clocks



www.npl.co.uk/

An Atomic Clock with 10^{-18} Instability

Yb atoms

N. Hinkley,^{1,2} J. A. Sherman,¹ N. B. Phillips,¹ M. Schioppo,¹ N. D. Lemke,¹ K. Beloy,¹
M. Pizzocaro,^{1,3,4} C. W. Oates,¹ A. D. Ludlow^{1*}

Science 341, 1215 (2013)

Sr atoms

An optical lattice clock with accuracy and stability at the 10^{-18} level

Nature 506, 71 (2014)

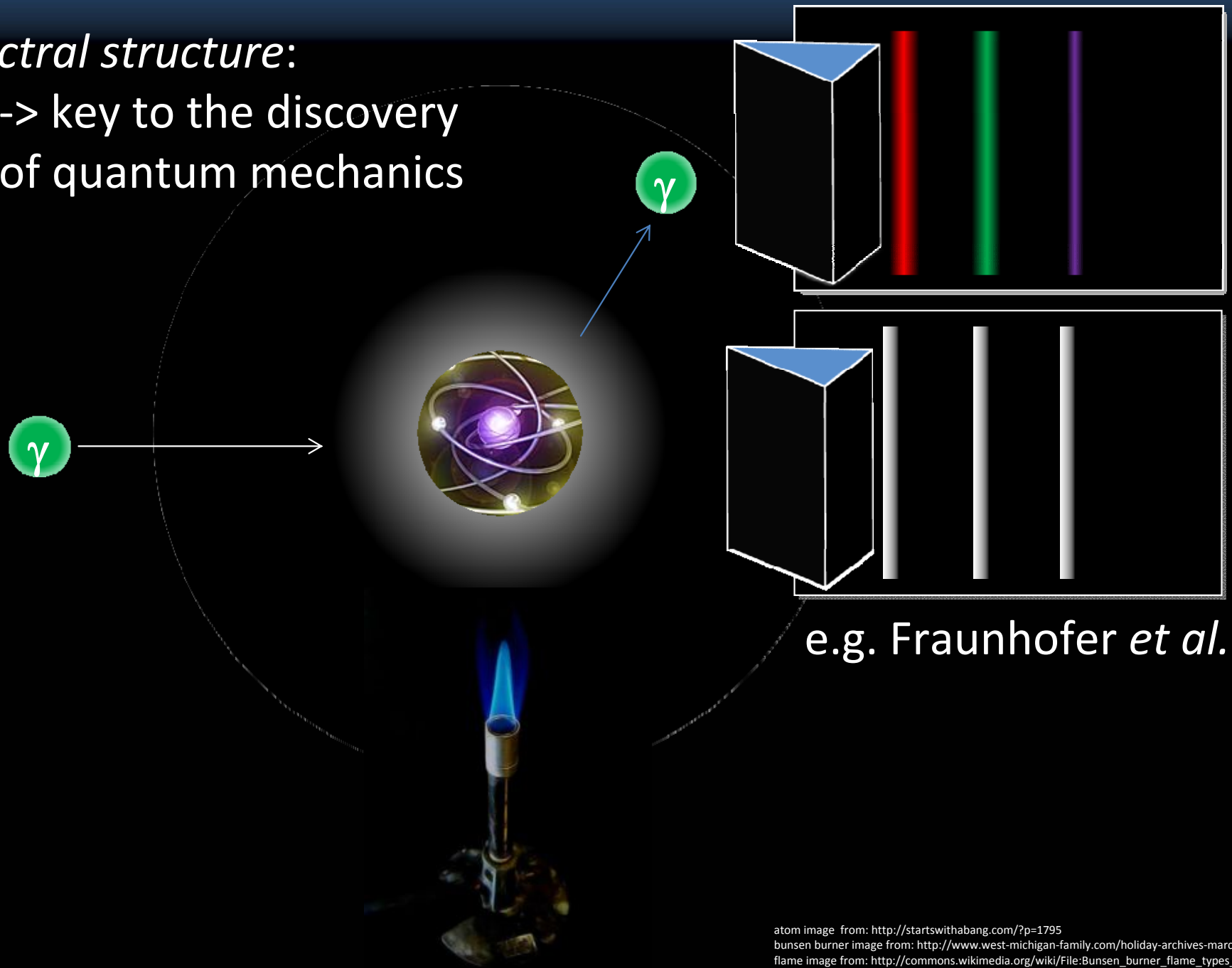
B. J. Bloom^{1,2*}, T. L. Nicholson^{1,2*}, J. R. Williams^{1,2†}, S. L. Campbell^{1,2}, M. Bishof^{1,2}, X. Zhang^{1,2}, W. Zhang^{1,2}, S. L. Bromley^{1,2} & J. Ye^{1,2}

traditional spectroscopy

(Kirchhoff, Bunsen, *et al.* @Heidelberg ~1860)

the *spectral structure*:

-> key to the discovery
of quantum mechanics



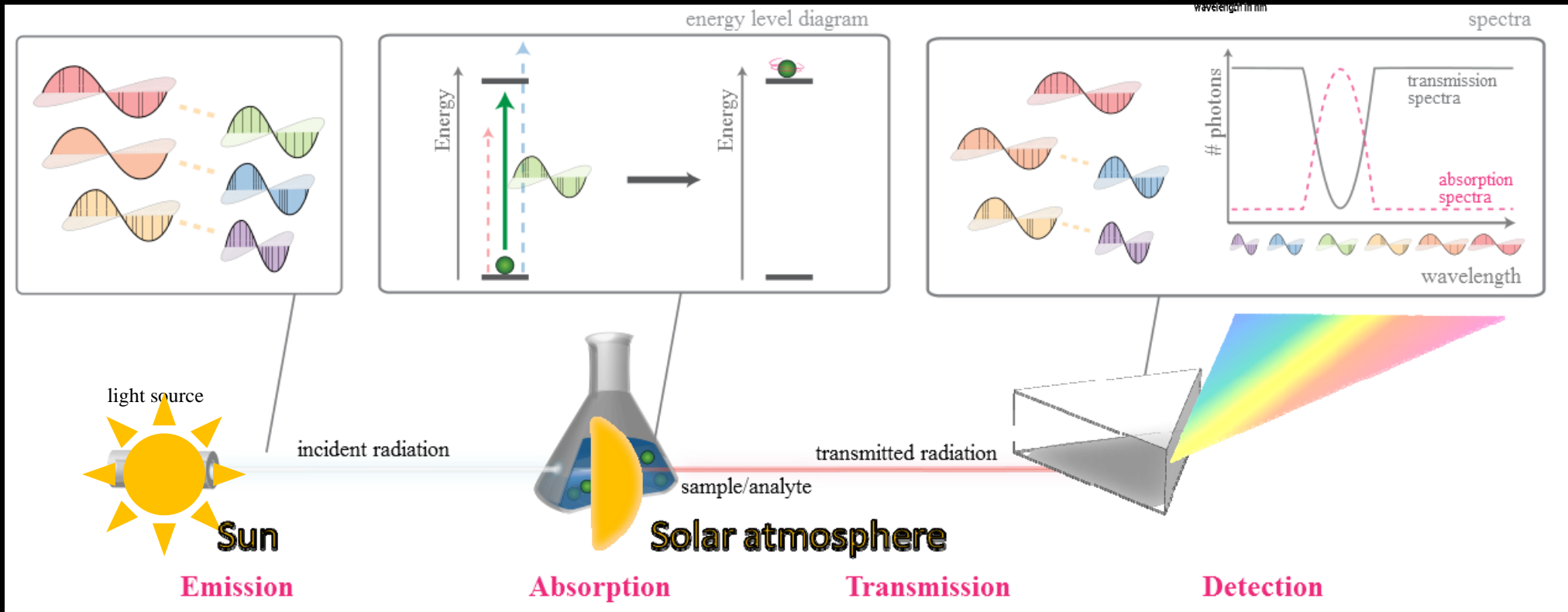
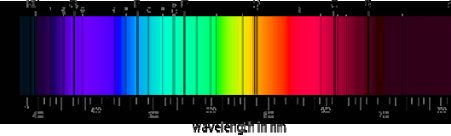
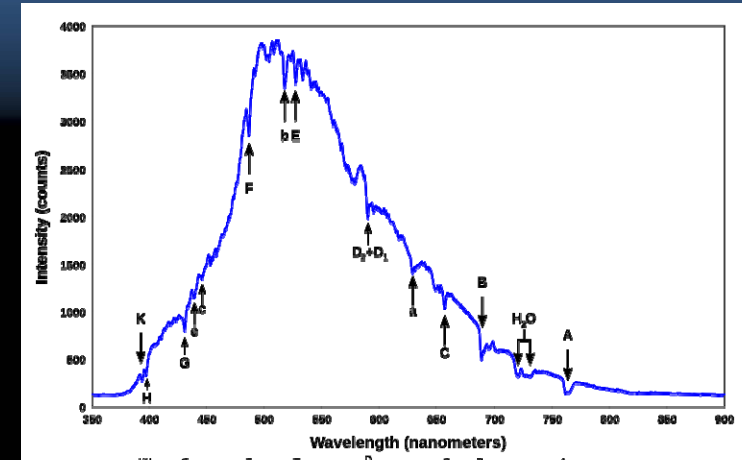
Absorption Spectroscopy

as explained on



Wollaston (1802), Fraunhofer (1814):
Solar spectrum
assigned to characteristic emission lines
Kirchhoff, Bunsen (1859)

Spectral line positions



Absorption Spectroscopy

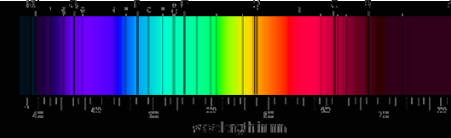
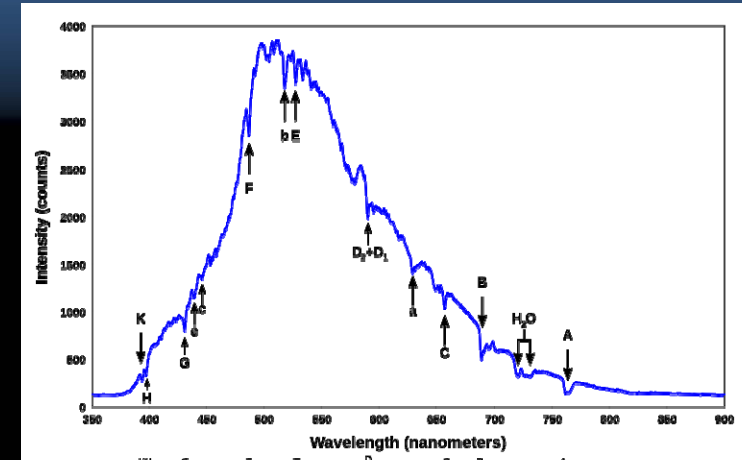
as explained on



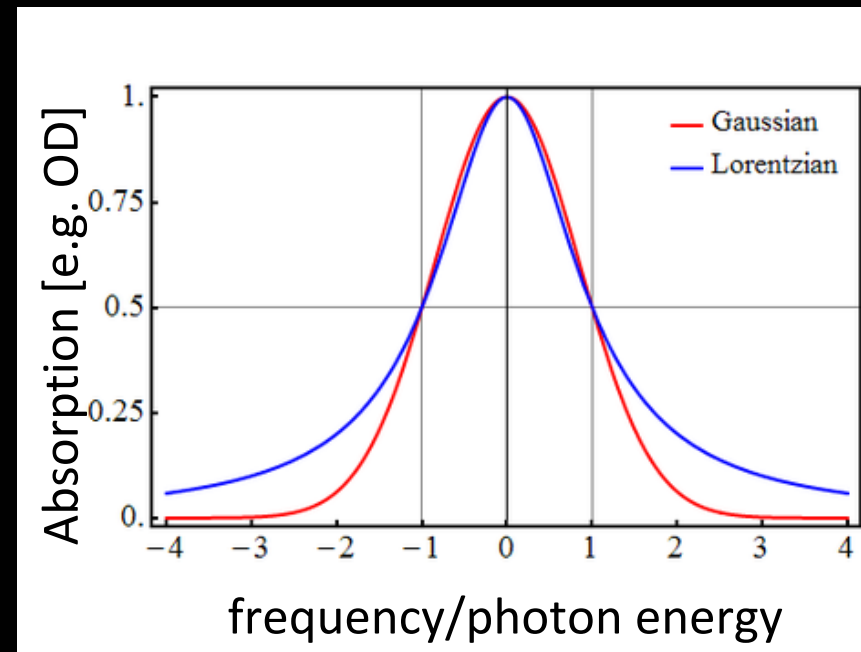
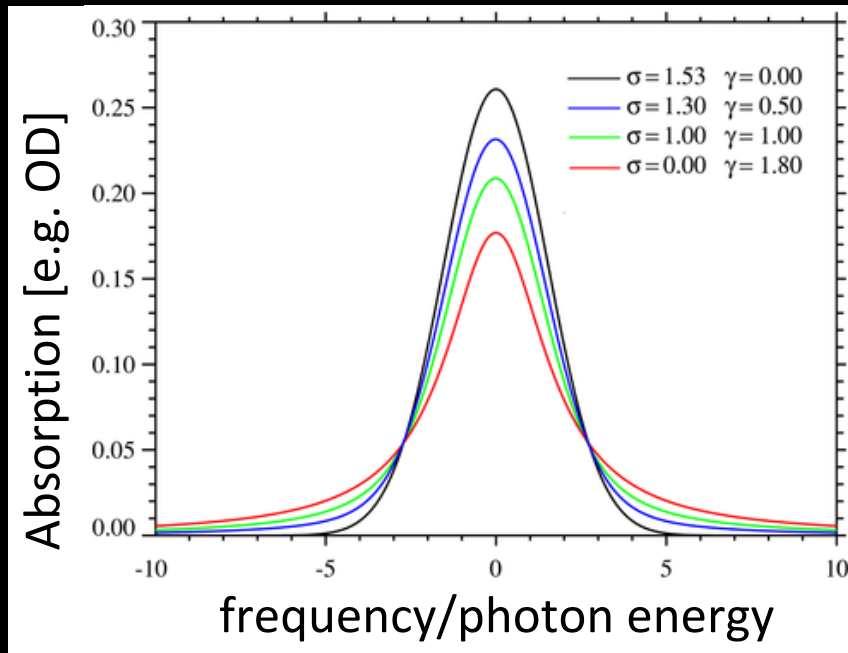
Wollaston (1802), Fraunhofer (1814):
Solar spectrum
assigned to characteristic emission lines
Kirchhoff, Bunsen (1859)

Spectral line positions

Spectral line widths



Spectral line shapes



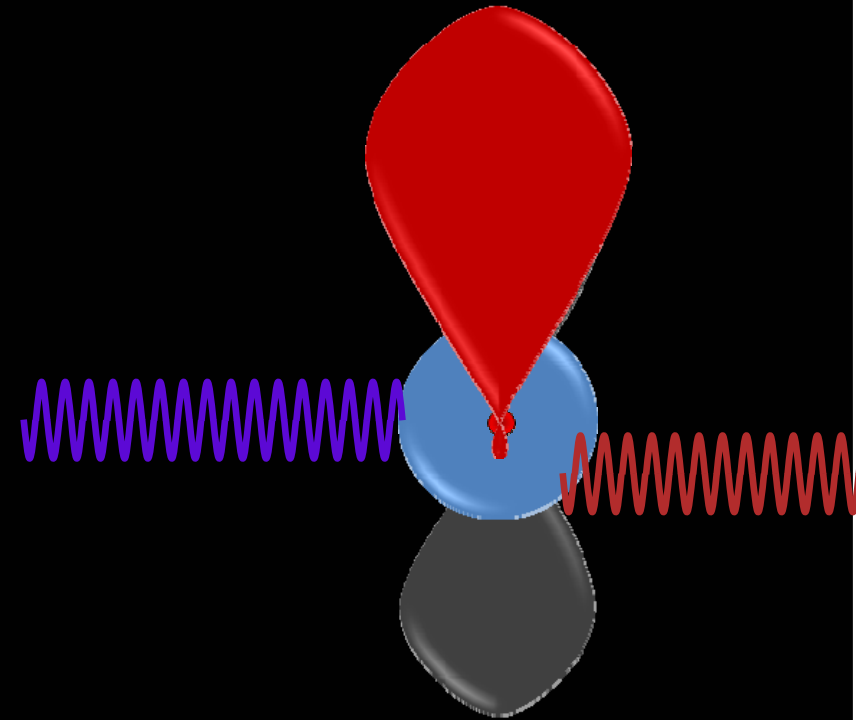
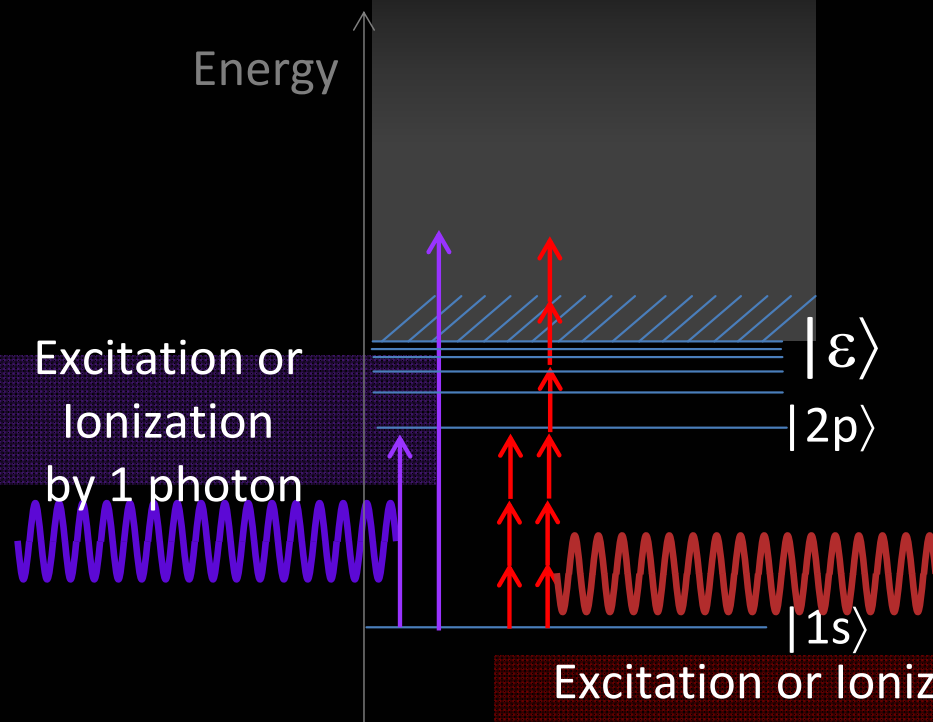
The „shape“ of atoms...

... gets changed by intense laser fields

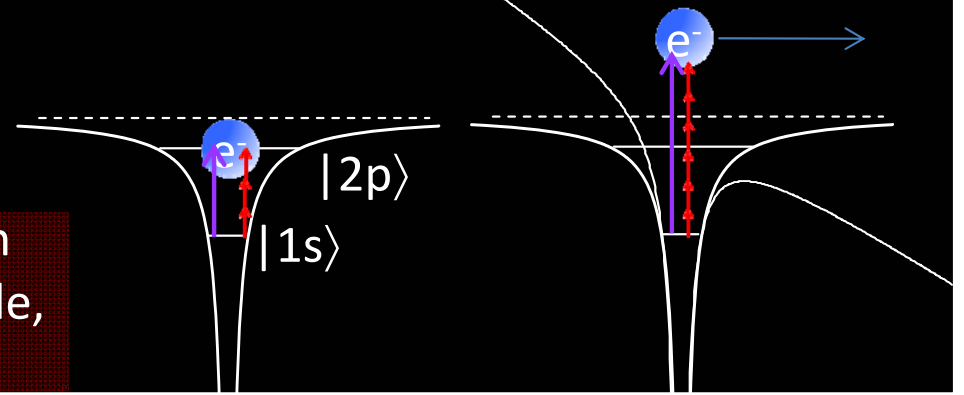
... and their interaction with light

weak ...

... and strong fields

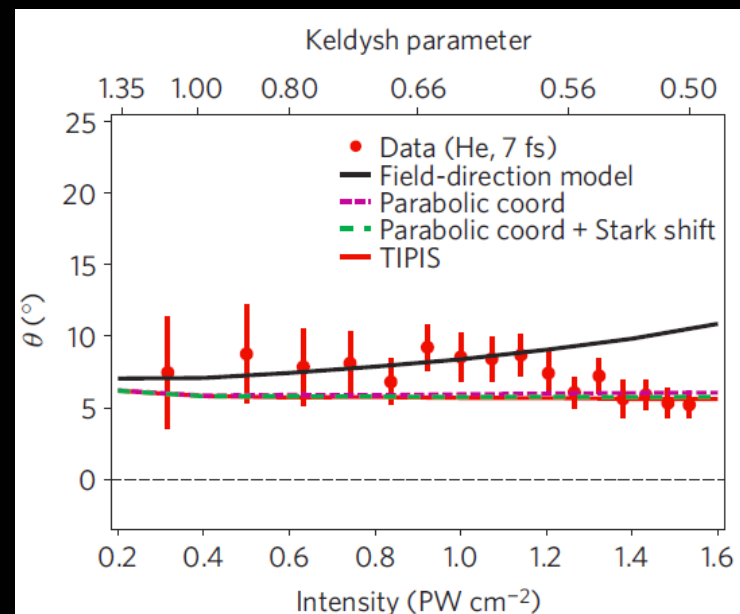
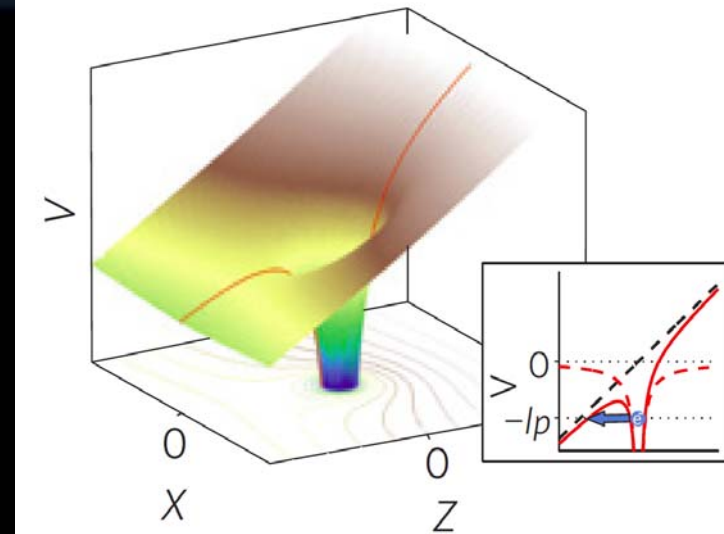
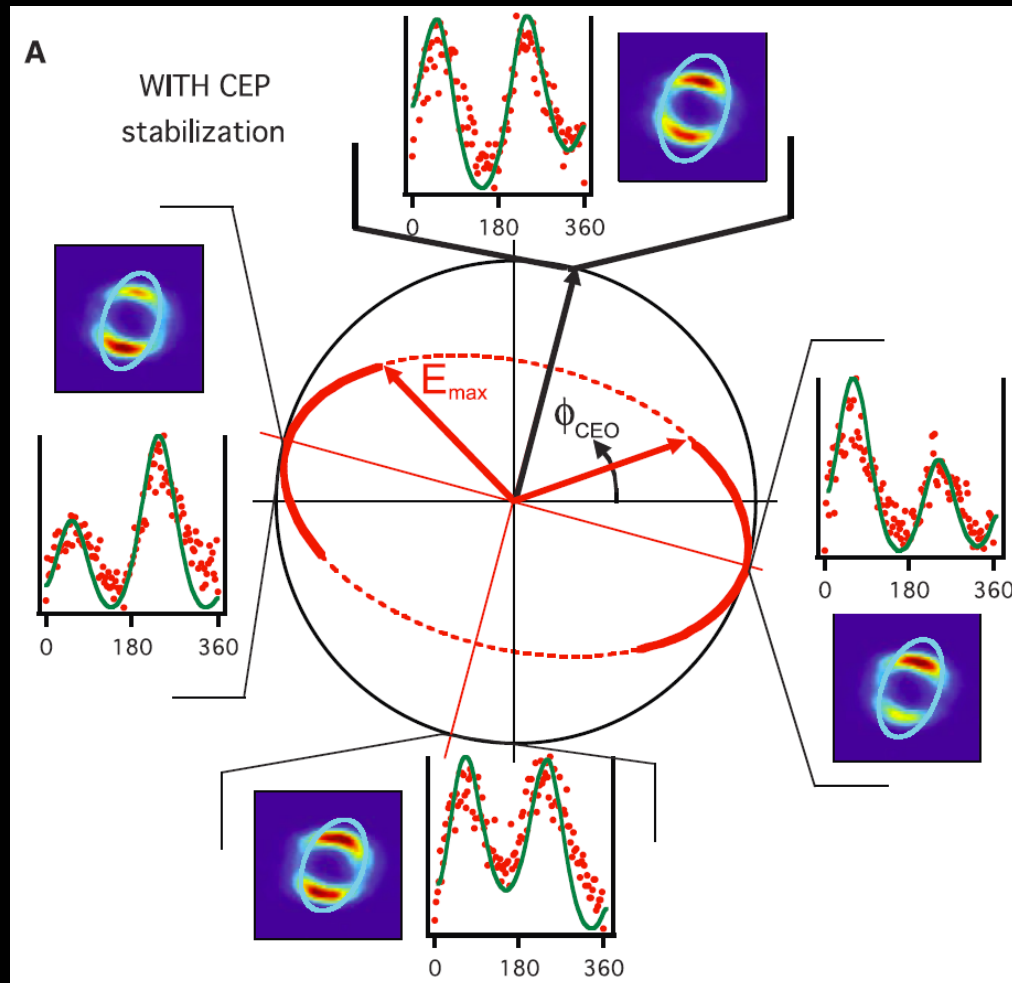


Pierre Agostini *et al.*
PRL **42**, 1127 (1979)



Strong-field ionization of Atoms

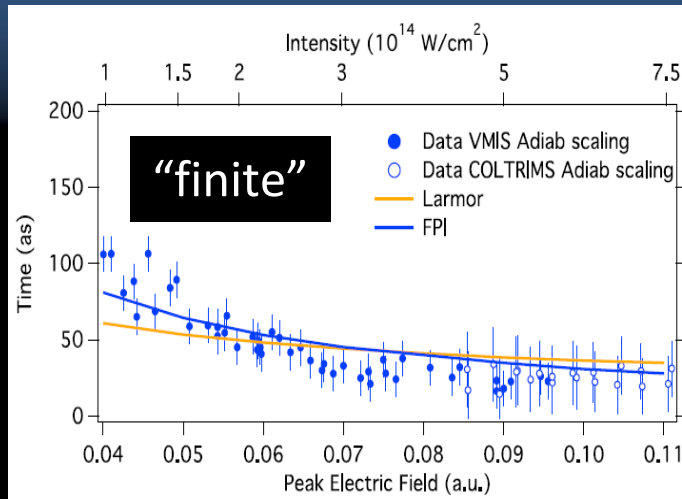
Keller group @ ETH
Attoclock, 10^{-18} sec precision



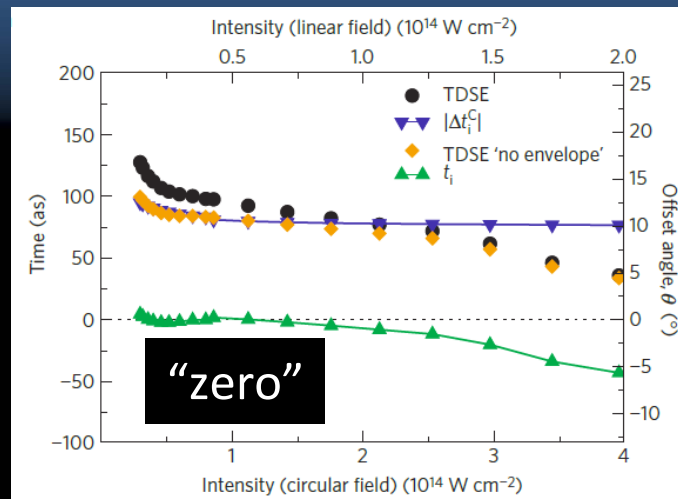
P. Eckle *et al.* Science **322**, 1525 (2008)

A. Pfeiffer *et al.* Nat. Physics (2012)

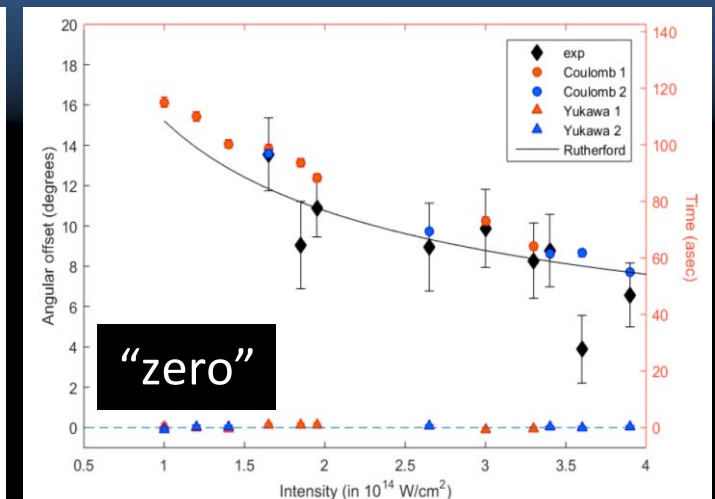
A never-ending controversy: “tunneling-time measurement”



Landsman *et al.* (Keller group)
Optica (2014)



Torlina *et al.* (Smirnova group)
Nat. Physics (2015)

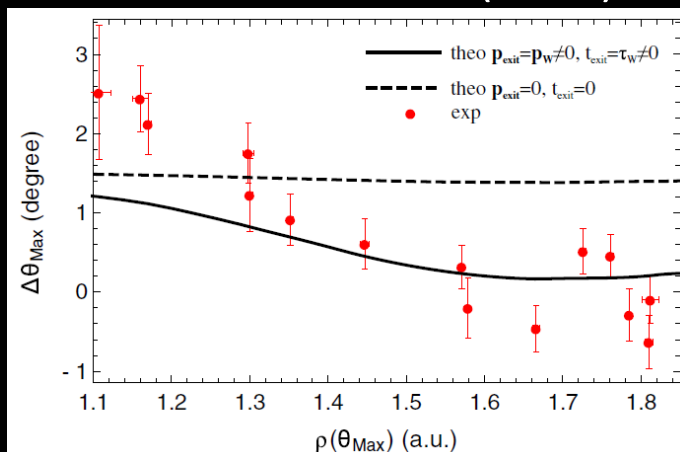


Sainadh *et al.* (Litvinyuk group)
arXiv:1707.05445 (2017)

Where is the problem?

- No time operator, hence: time not directly observable
- need to define a “time standard” (“controversy”)
- here in particular: mapping quantum dynamics (tunneling) into classical dynamics (trajectory)

Camus *et al.* PRL (2017)



Definition: Wigner time (Keitel group)

->Key: allow both **tunneling time** and **momentum**

Experimental test (Moshhammer group)

->Key: **Compare two species** (Ar, Kr)

by coincidence measurements
under otherwise **identical conditions**
with high angular precision

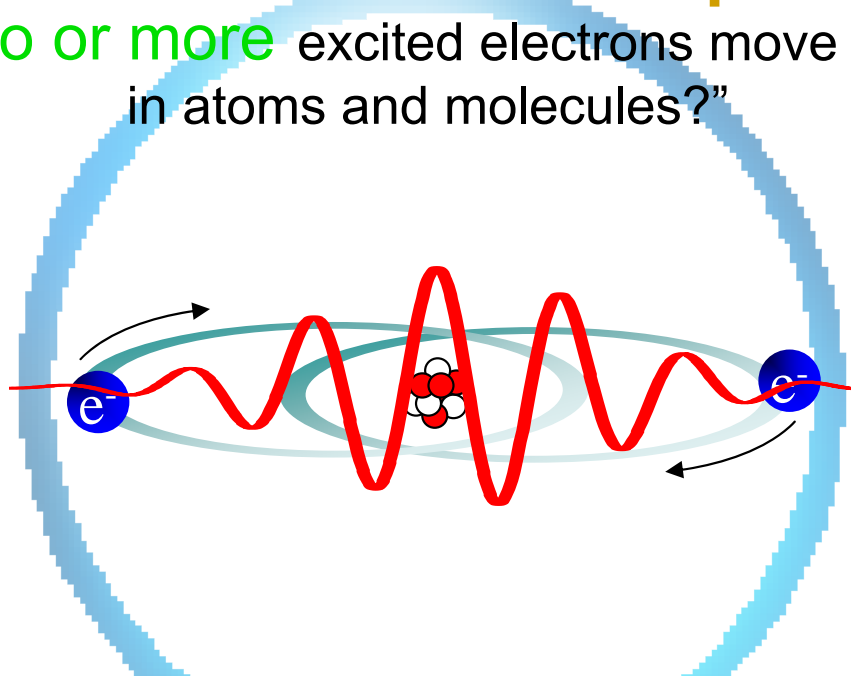
few-body quantum dynamics

a fundamental scientific question:

“how do **two or more** excited electrons move and interact in atoms and molecules?”

spatial scale
 $R \sim \text{sub/few } \text{\AA}$

temporal scale
 $T \sim \text{sub/few fs}$



The
“**quantum
few-body
problem**”
in strong fields

Scientific goal:

*measure / understand / control
the **quantum dynamics** of
few-body systems
*in strong fields**

0 as

(x-ray) movies of
single molecules

Laser control of
chemical reactions

Petahertz-clocked
computing

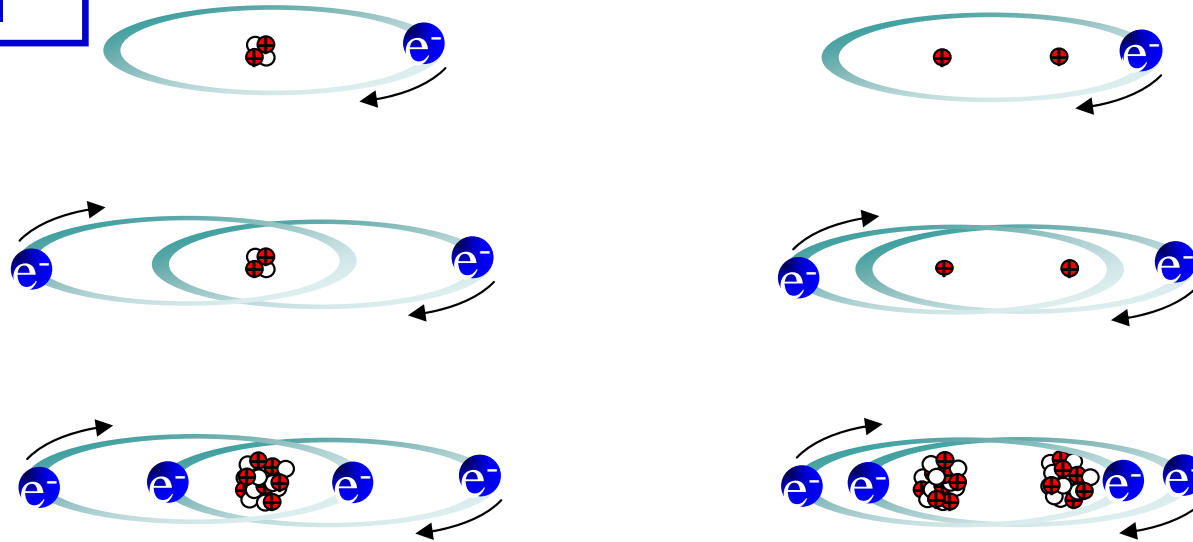
x-ray precision
spectroscopy

Fundamental Quantum Dynamics

Fermionic condensates in a nanotrap!

in atoms and small molecules:

more fundamental

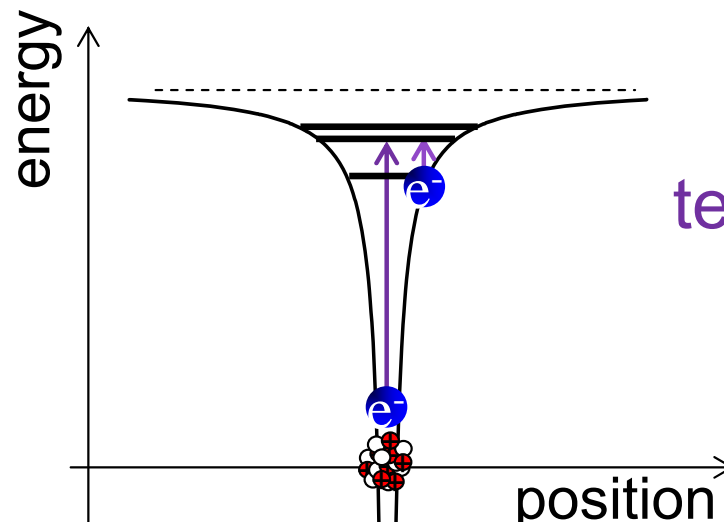


The “quantum few-body problem”

more complex

naturally well-defined, isolated, Å-sized 'labs'

What are the tools?



typical ΔE for transitions:
ten(s) of **eV** to **few keV**

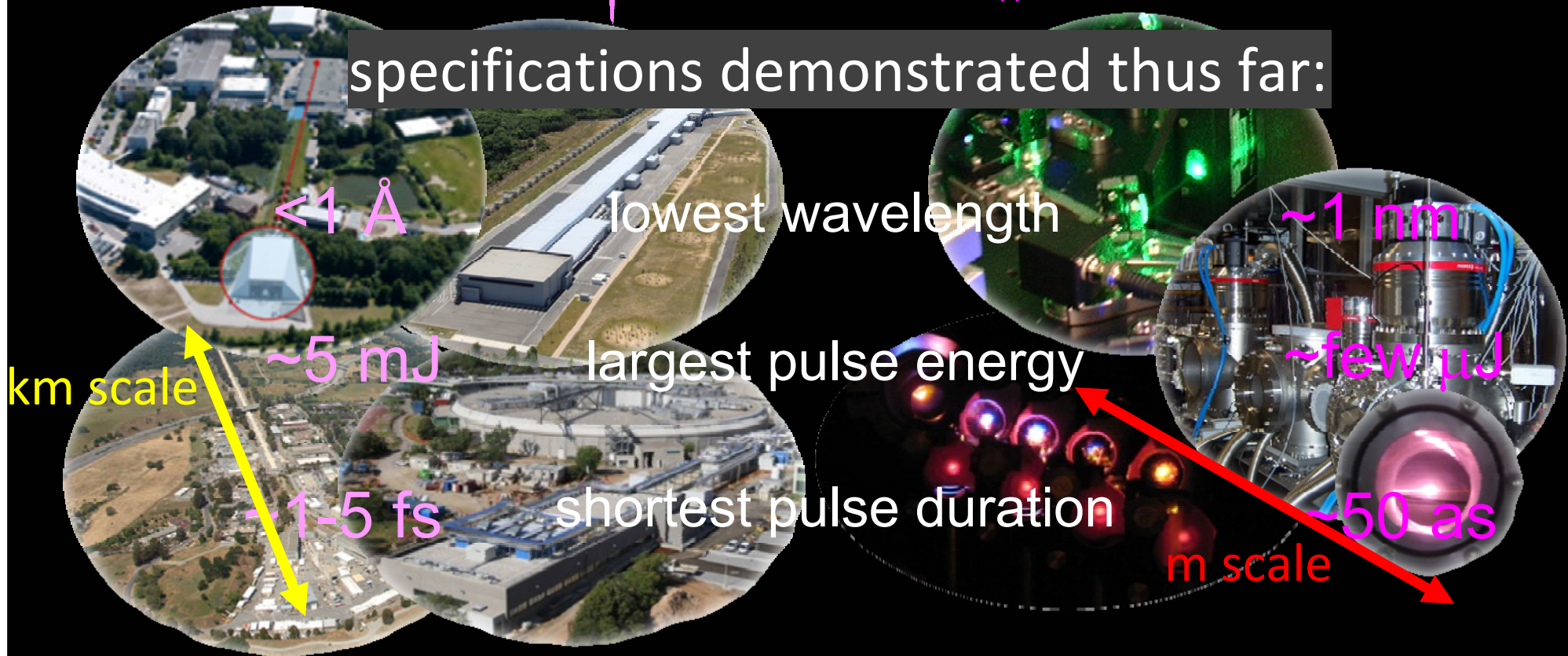
XUV to **x-ray**
light required for measurement and control

Two parallel revolutions in ultrashort **x-ray/XUV** laser science

Free **E**lectron **L**asers

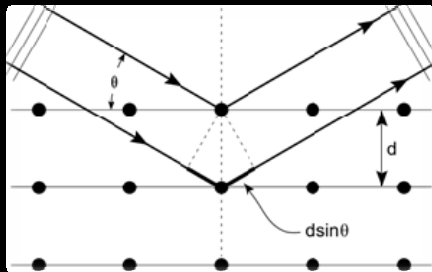
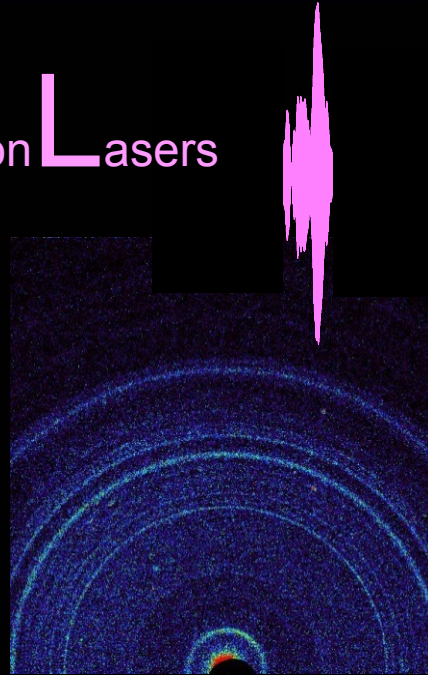
High **H**armonic **G**eneration

specifications demonstrated thus far:



Time-resolved Science with novel x-ray/XUV laser sources

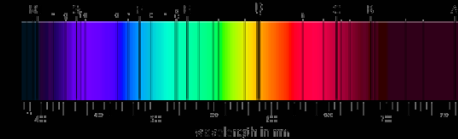
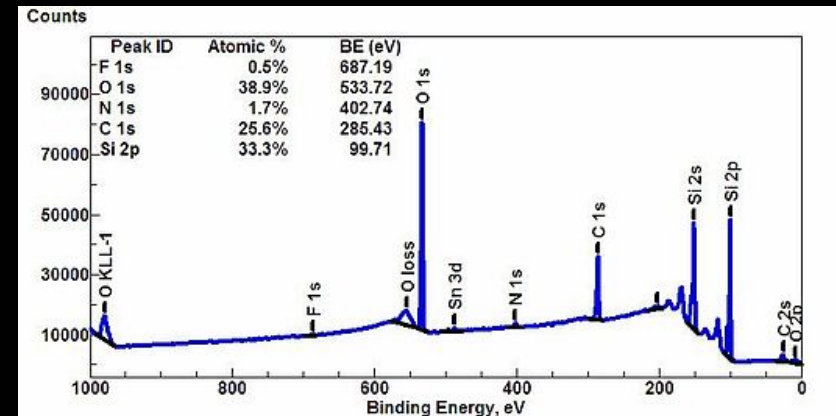
Free **E**lectron **L**asers



X-ray diffraction

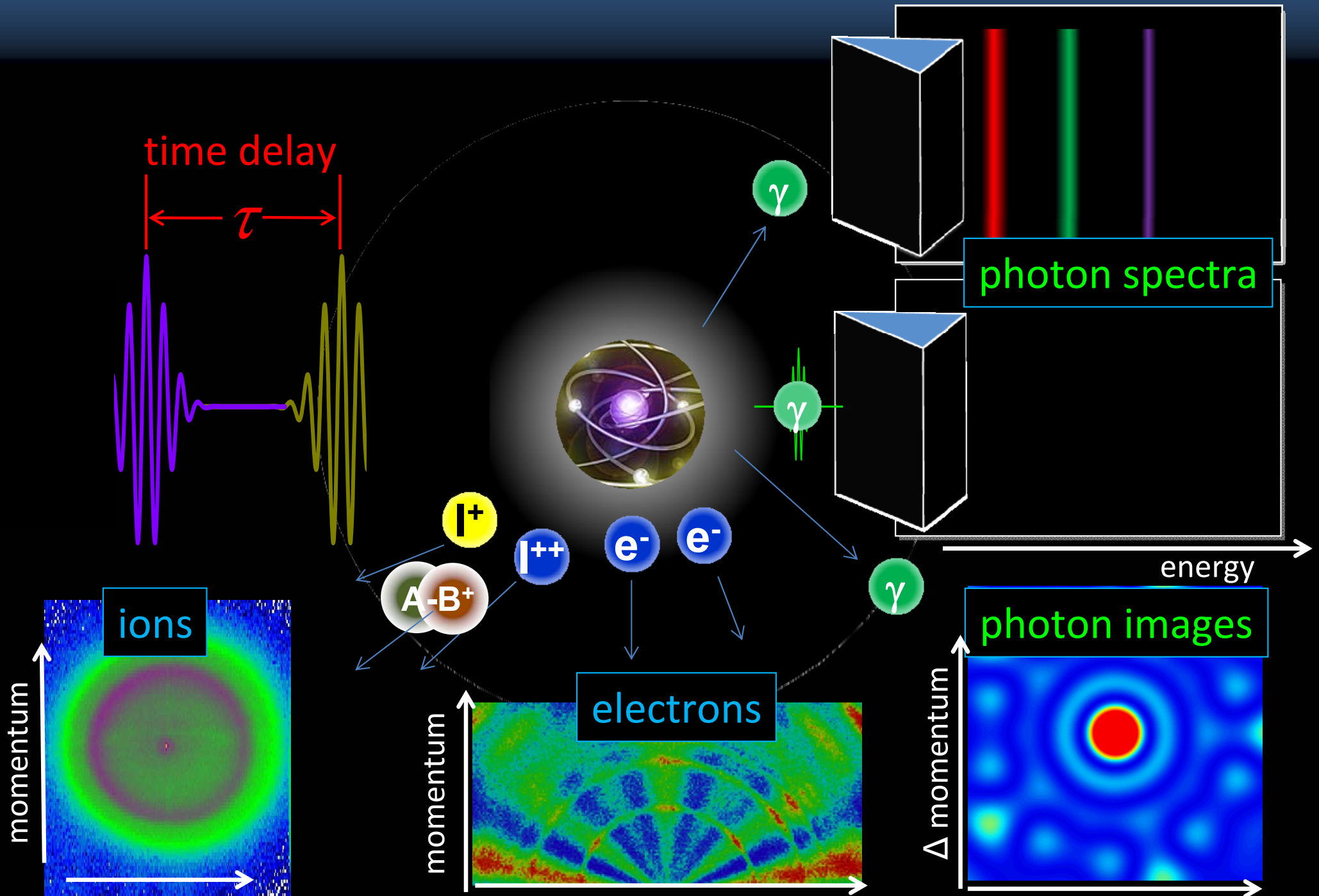


High **H**armonic **G**eneration



XUV spectroscopy

Time-resolved imaging and spectroscopy

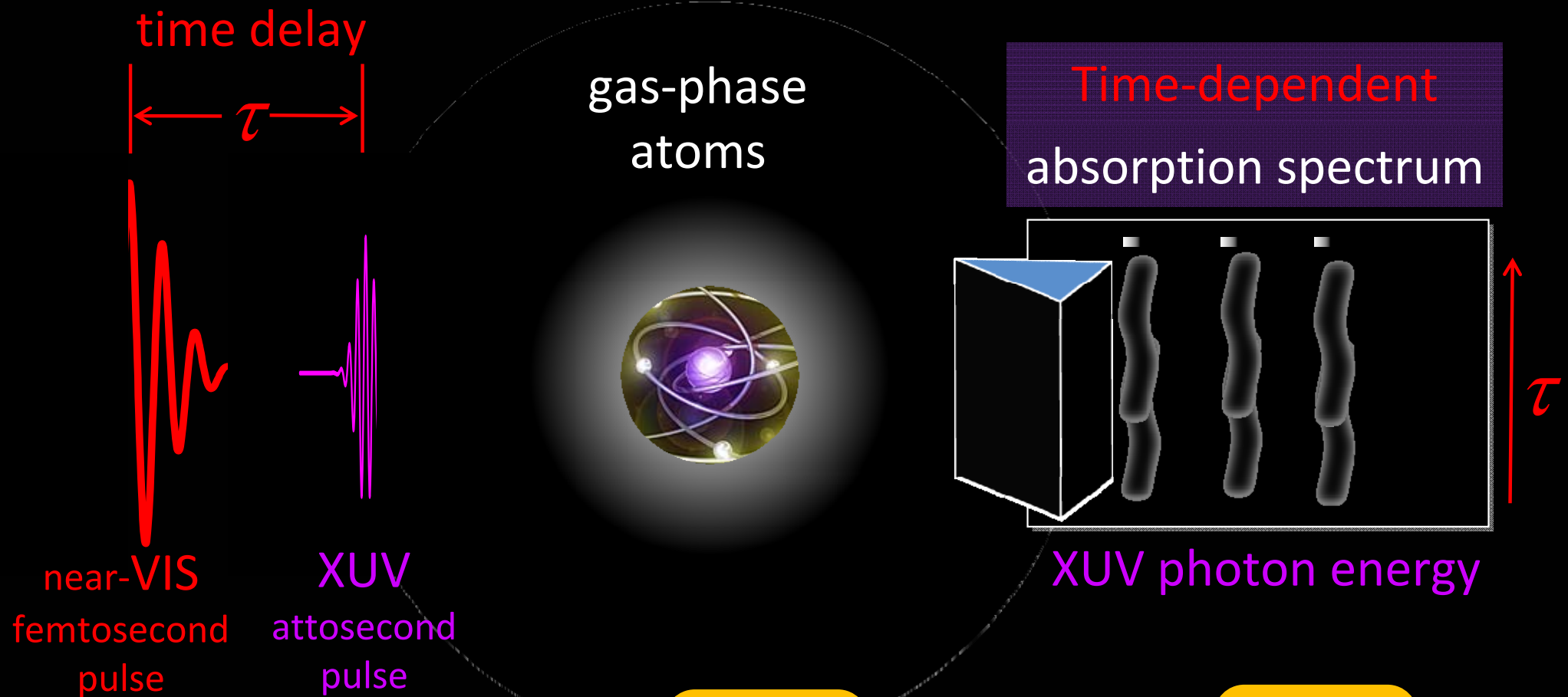


time-dependent XUV absorption spectroscopy

HHG-based applications, active groups:

Exp.: Leone, Neumark, Keller, Gallmann, Chang, Krausz, Sansone, Kim, Sandhu, Vrakking, Wörner ...

Theory: Schafer, Gaarde, Santra, Martín, Argenti, Rost, Greene, Keitel, Stockman, ...



$$\psi(x, t) \propto (a_1) \psi_1(x) e^{-\frac{i}{\hbar} E_1 t} + (a_2) \psi_2(x) e^{-\frac{i}{\hbar} E_2 t}$$

Phase φ_1

Phase φ_2

time-dependent XUV absorption spectroscopy

What happens to bound states and resonances in *short* and *strong* fields ?



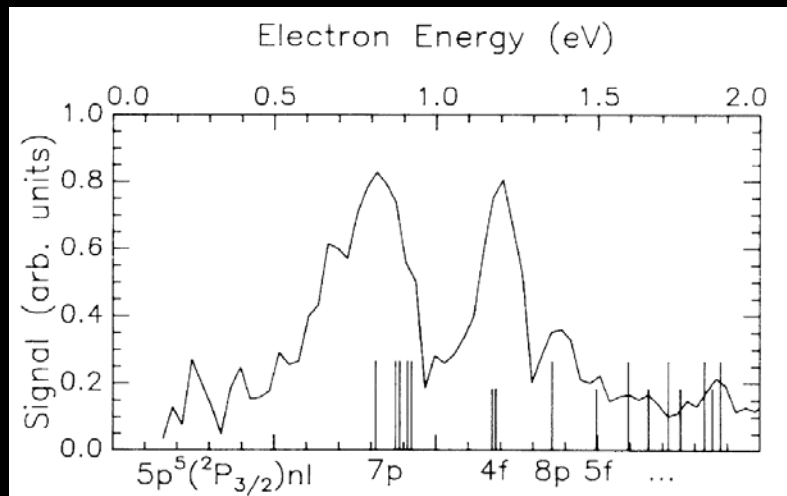
$$\psi(x, t) \propto (a_1) \psi_1(x) e^{-\frac{i}{\hbar} E_1 t} + (a_2) \psi_2(x) e^{-\frac{i}{\hbar} E_2 t}$$

Phase φ_1

Phase φ_2

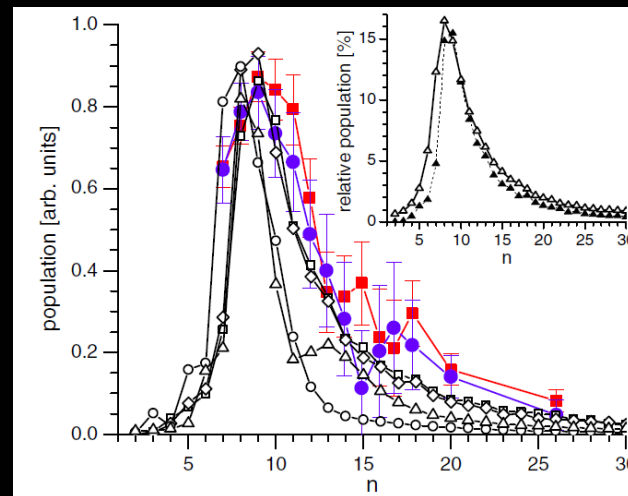
Role of bound (excited) states and resonances in (short) strong fields

Freeman resonances in strong-field ionization



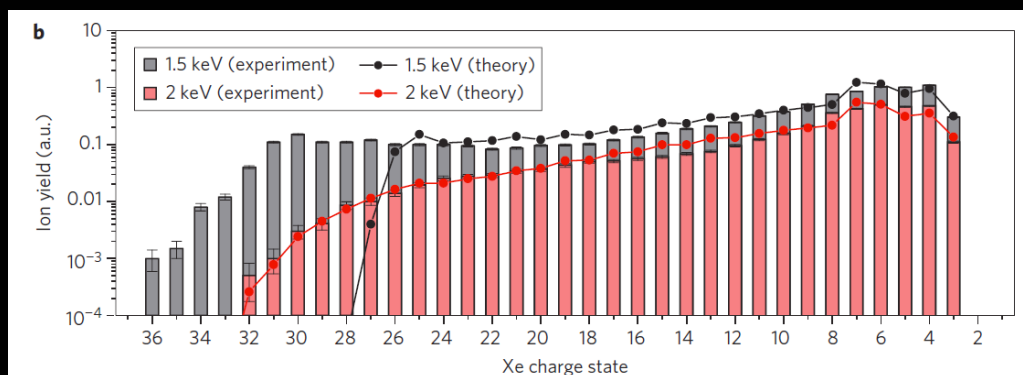
Freeman, Bucksbaum *et al.* PRL **59**, 1092 (1987)

frustrated tunnel ionization



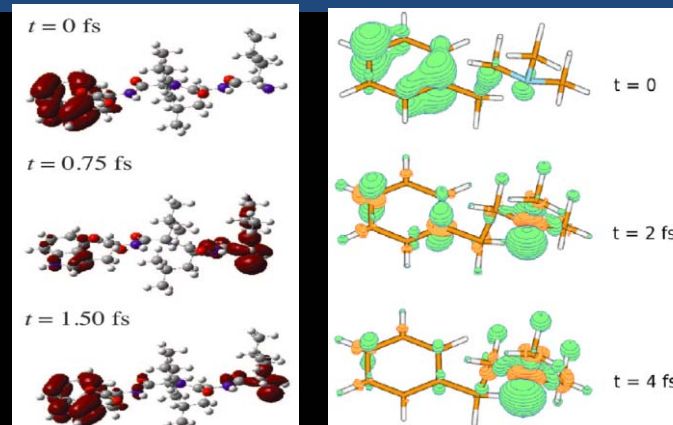
Zimmermann, Buller, Eilzer, Eichmann PRL **114**, 123003 (2015)

x-ray FEL multiple ionization



Rudek *et al.* Nat. Photonics **6**, 858 (2012)

molecular charge migration



Remacle, Levine PNAS **103**, 6793 (2006)

Lünnemann, Kuleff, Cederbaum CPL **450**, 232 (2008)

General Physics: coupling of states

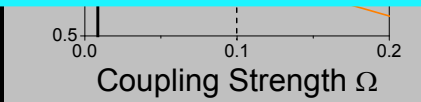
coupling of one to one other state

coupling of one to multiple states

coupling of one to a continuum of states

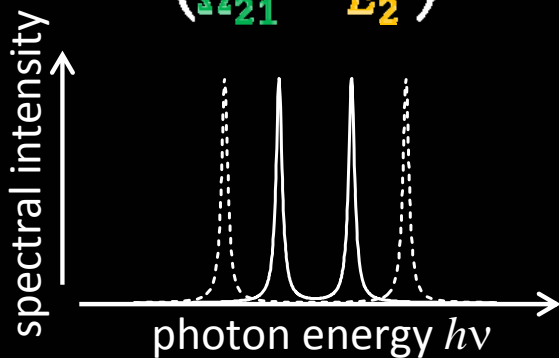
Quite well understood for cases of time-independent (or adiabatic) couplings

What happens when the coupling is **ultrashort** ?



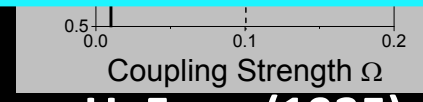
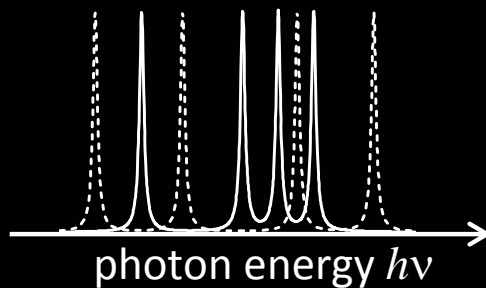
Rabi oscillation
in strong coupling

$$\begin{pmatrix} E_1 & \Omega_{12} \\ \Omega_{21} & E_2 \end{pmatrix}$$



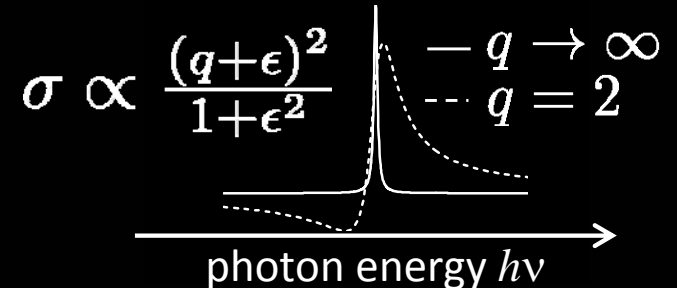
Breit-Rabi
e.g. Paschen-Back regime

$$\begin{pmatrix} E_1 & \Omega_{12} & \Omega_{13} & \Omega_{14} \\ \Omega_{21} & E_2 & \Omega_{23} & \Omega_{24} \\ \Omega_{31} & \Omega_{32} & E_3 & \Omega_{34} \\ \Omega_{41} & \Omega_{42} & \Omega_{43} & E_4 \end{pmatrix}$$



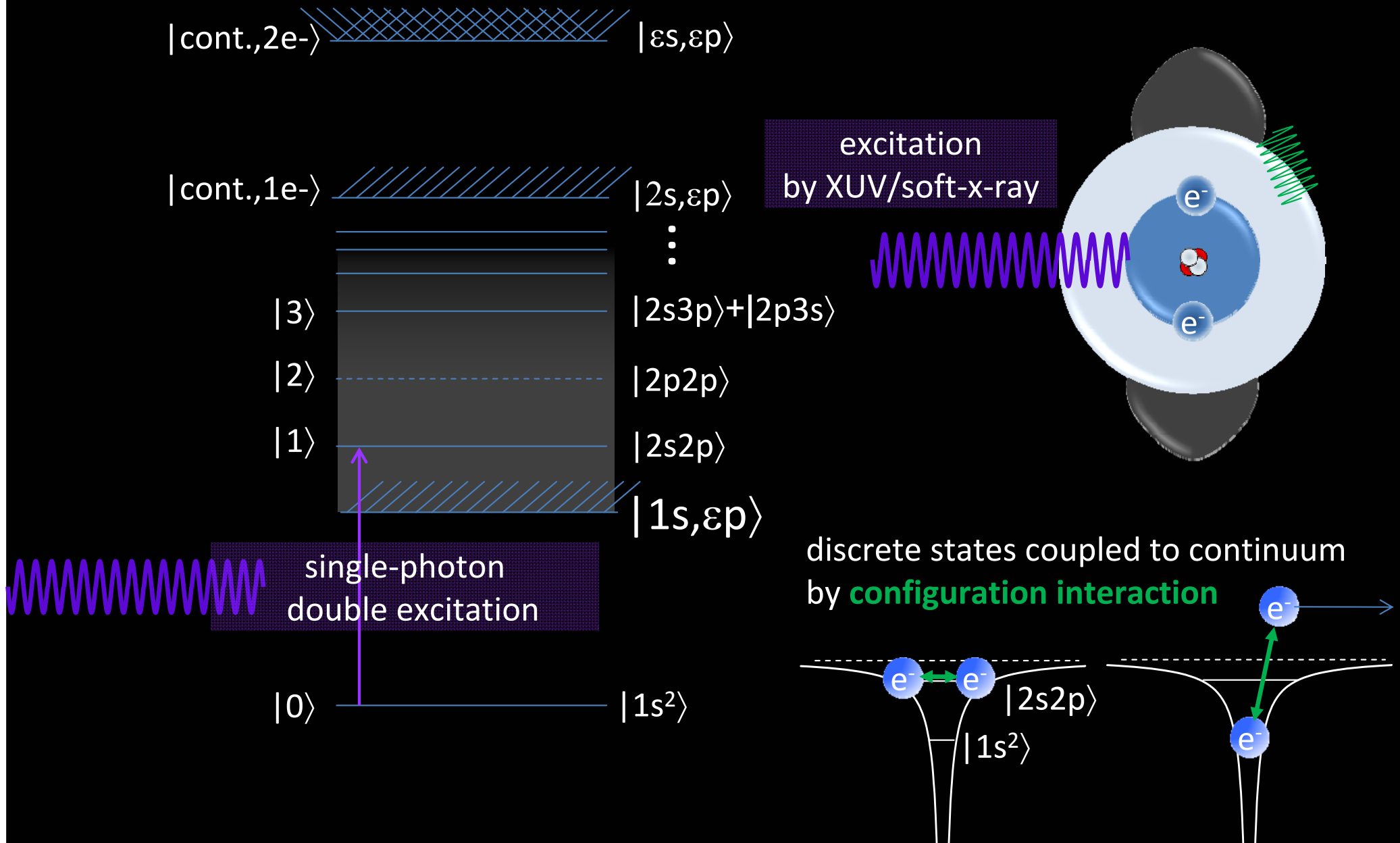
U. Fano (1935)
Phys. Rev. **124**, 1866 (1961)
Nuovo Cim. **12**, 154 (1935)

$$\begin{pmatrix} E & \Omega(\epsilon) \\ \Omega(\epsilon) & \epsilon \end{pmatrix}$$



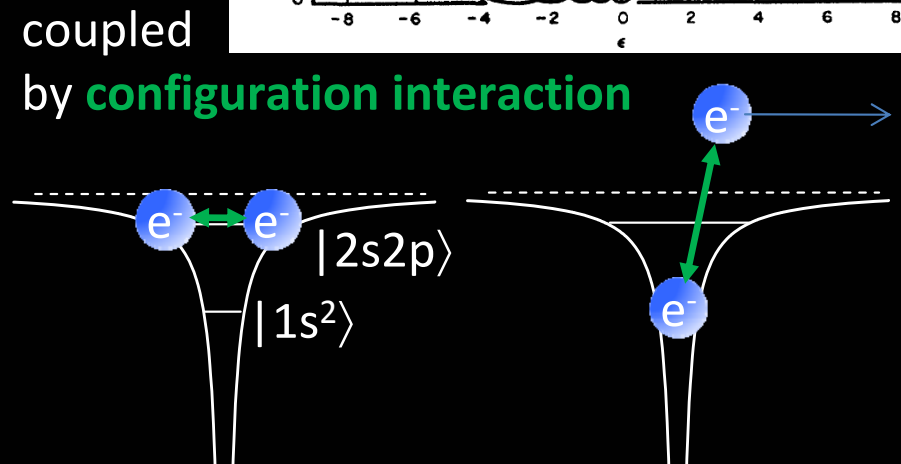
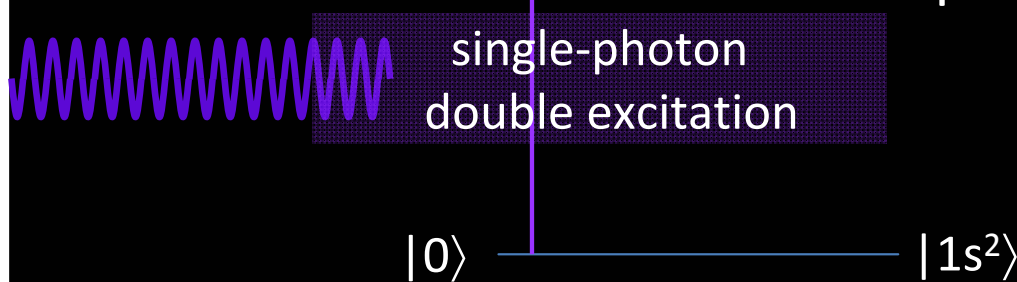
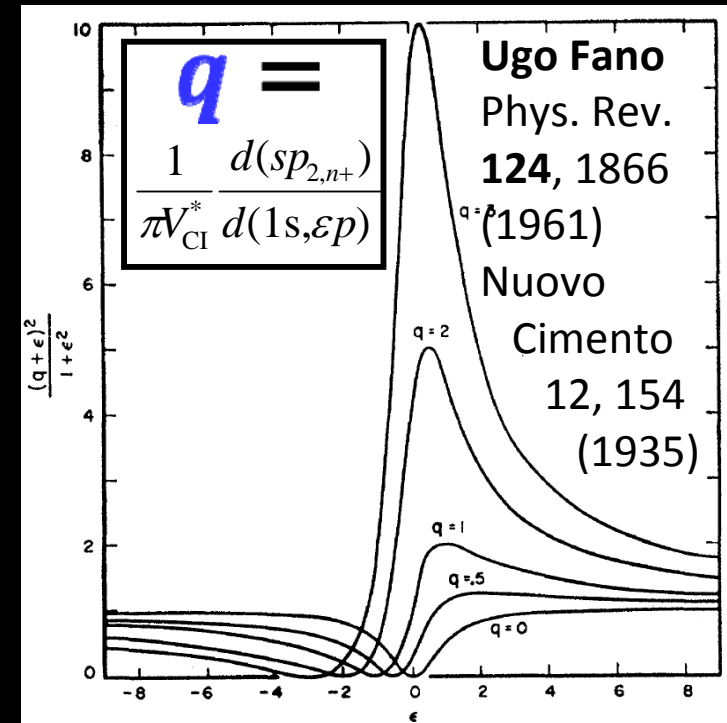
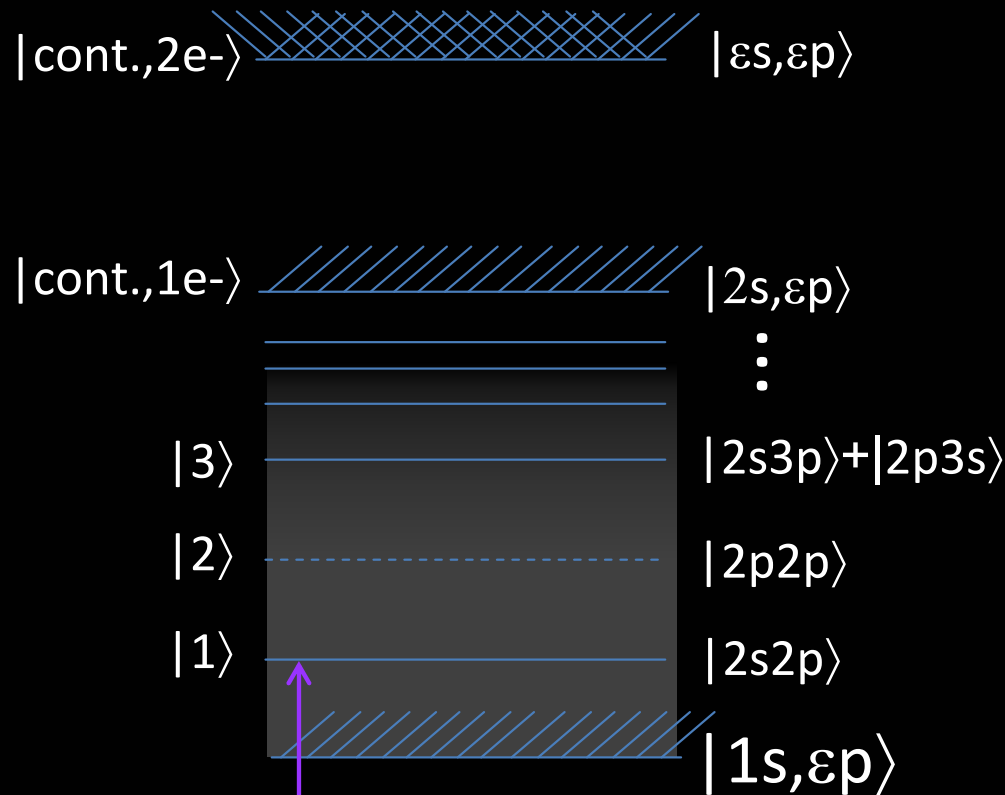
correlated bound-state dynamics: Doubly excited helium

a prototype system
for electron correlation



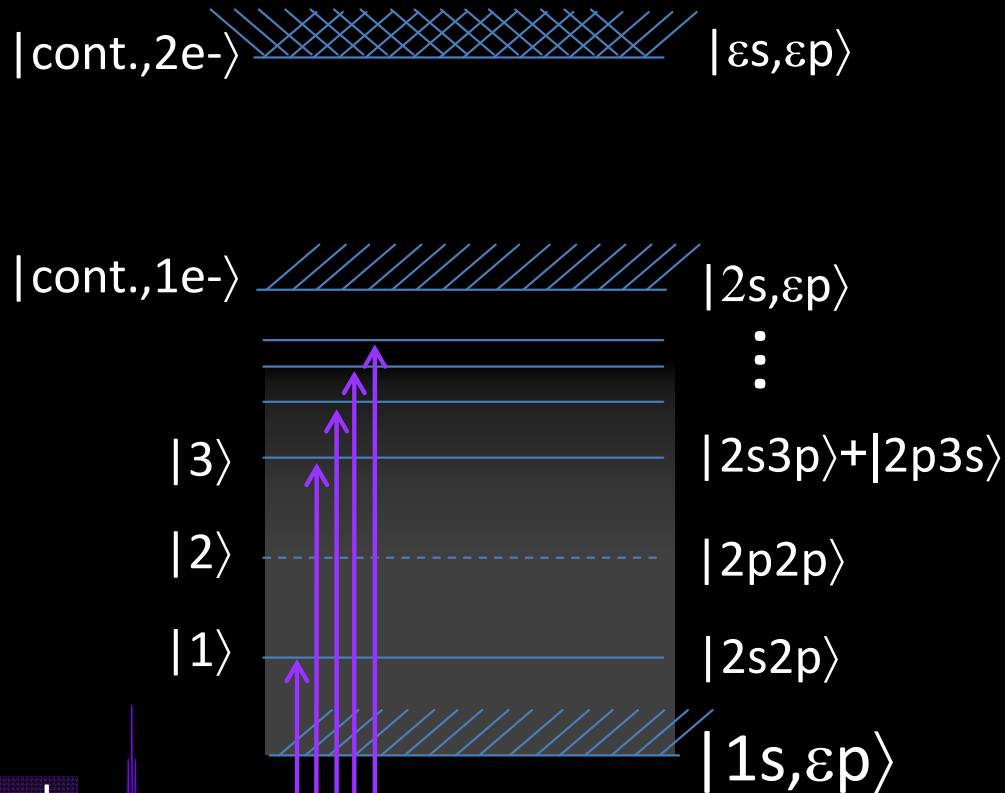
doubly excited helium: Fano resonance

$$\sigma_{Fano} \sim \frac{(q + \epsilon)^2}{1 + \epsilon^2}$$



doubly excited helium: Fano resonance

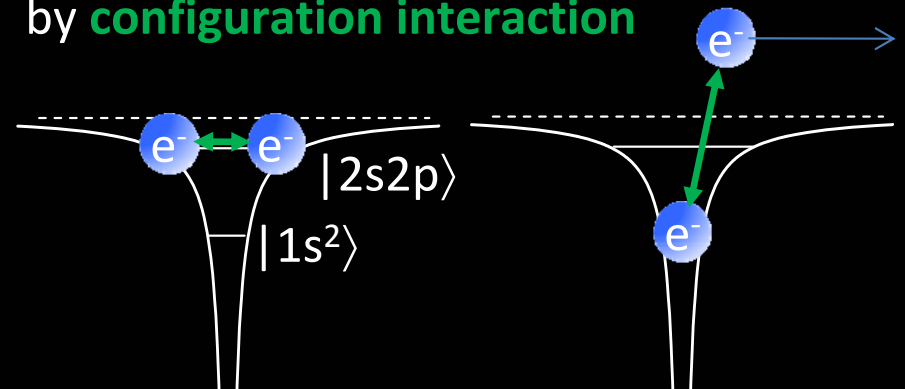
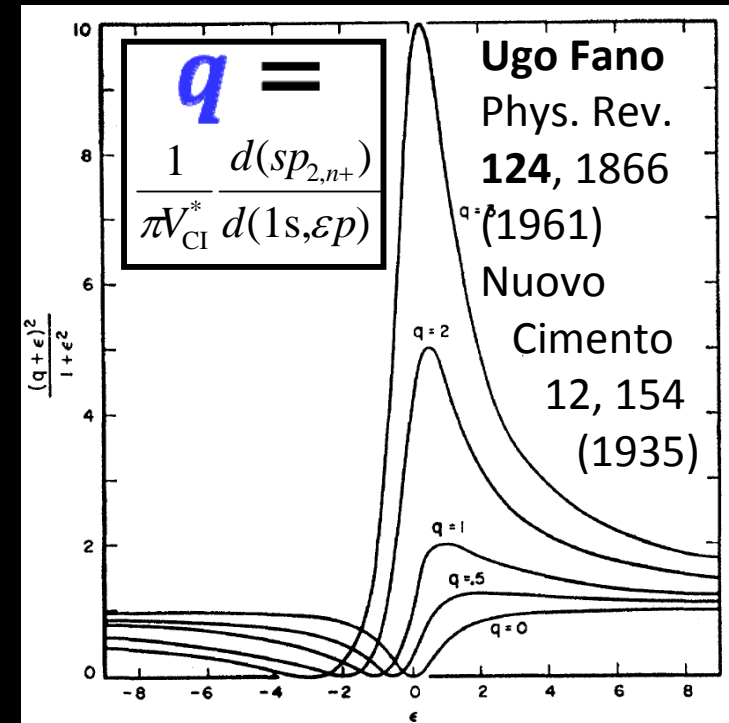
$$\sigma_{Fano} \sim \frac{(q + \epsilon)^2}{1 + \epsilon^2}$$



attosecond pulse
60 eV

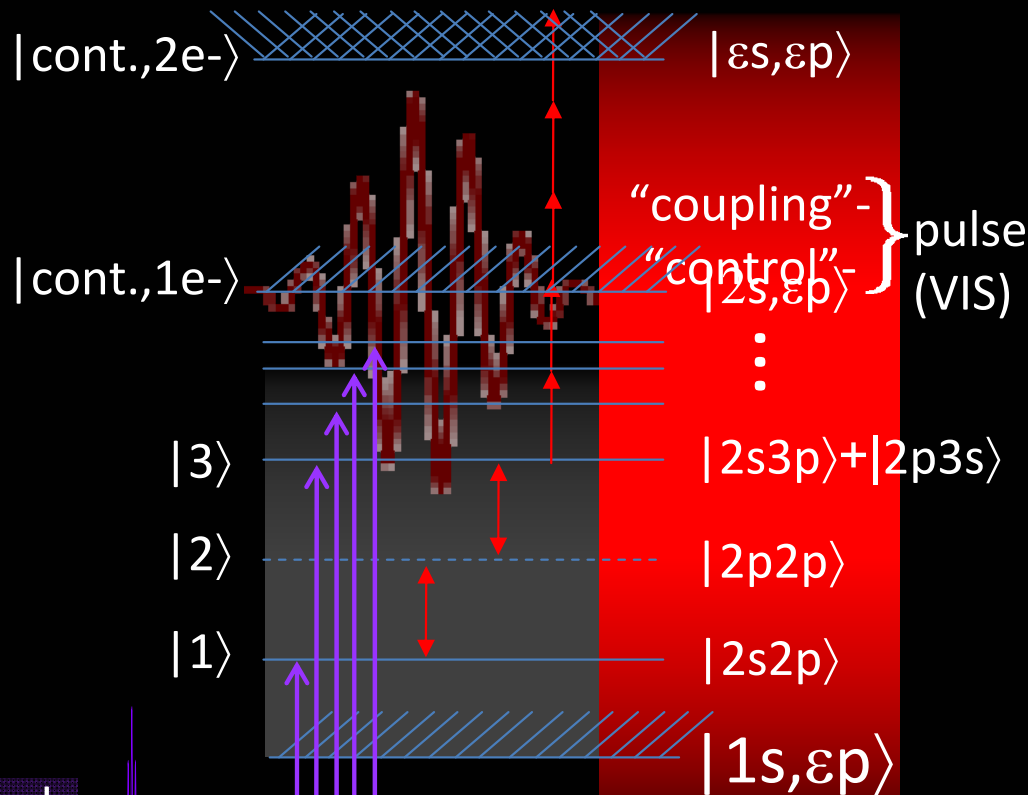
single-photon
double excitation

coupled
by **configuration interaction**



doubly-excited helium, in a strong laser field

$$\sigma_{Fano} \sim \frac{(q + \epsilon)^2}{1 + \epsilon^2}$$

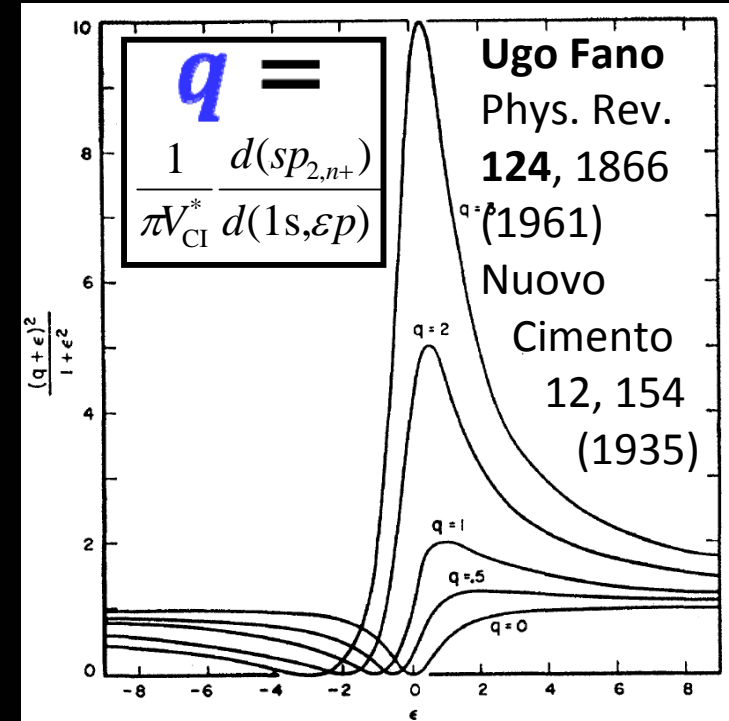
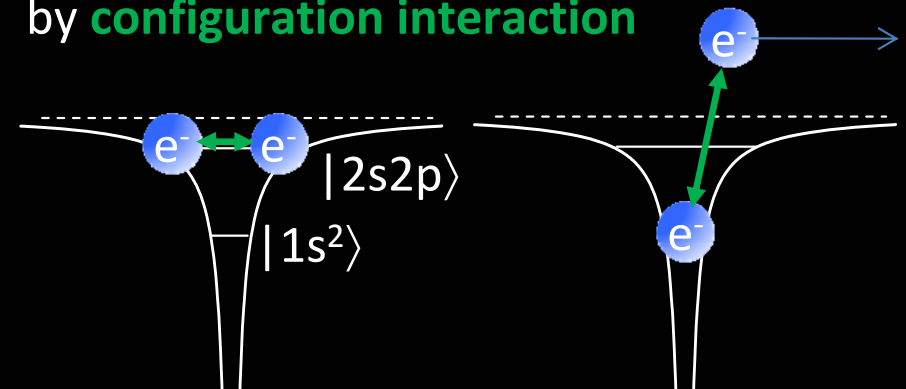


attosecond pulse
60-70 eV

single-photon
double excitation

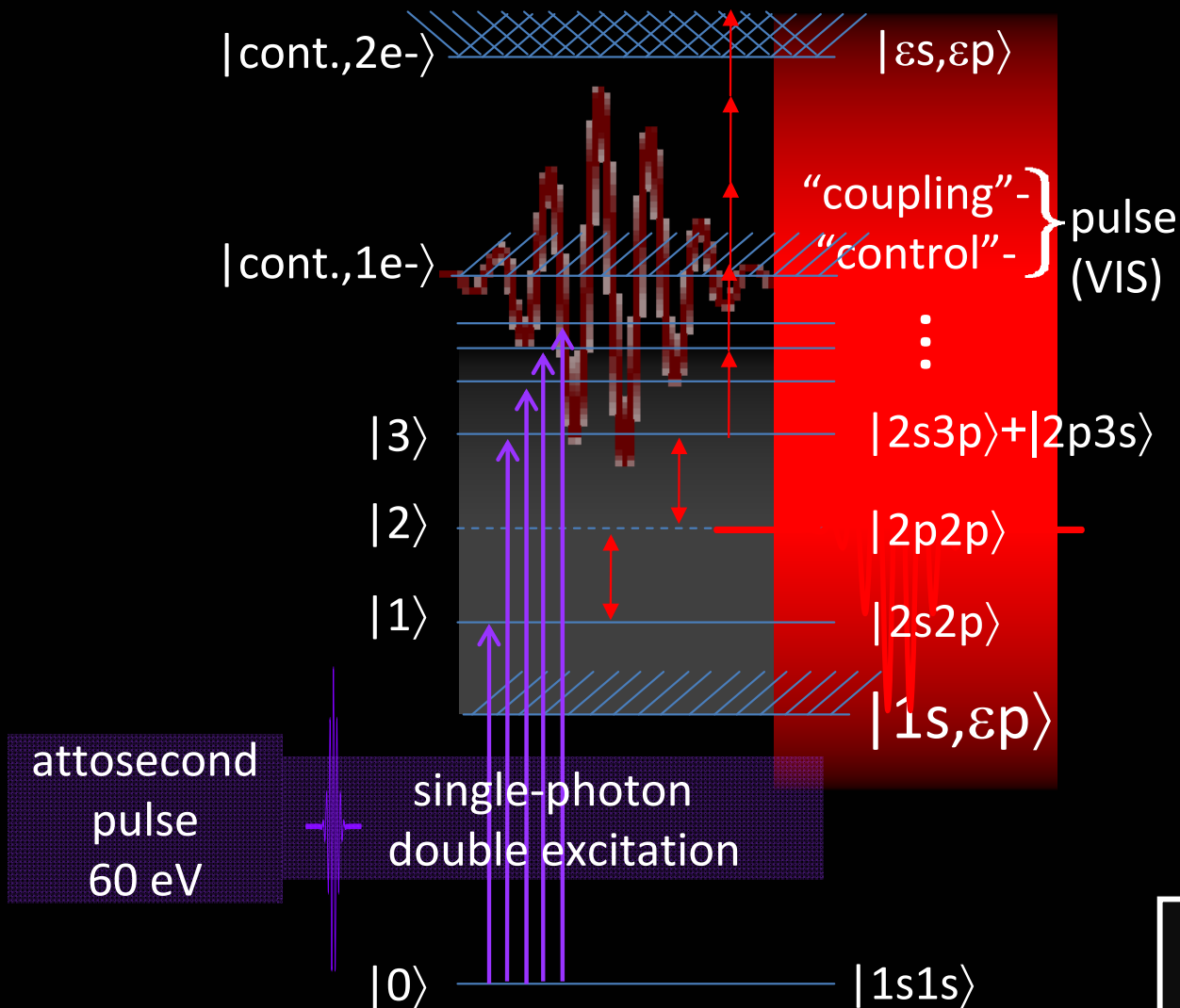
coupled

by **configuration interaction**



Ugo Fano
Phys. Rev.
124, 1866
(1961)
Nuovo
Cimento
12, 154
(1935)

doubly-excited helium, coupled to a laser field



Previous work

(on laser coupling of doubly-excited helium):

Theory:

- Madsen, Themelis, Lambropoulos
- Zhao, Chu, Lin et al.

...

Experiment:

- Loh, Greene, Leone, et al.
- Gilbertson, Chang et al.

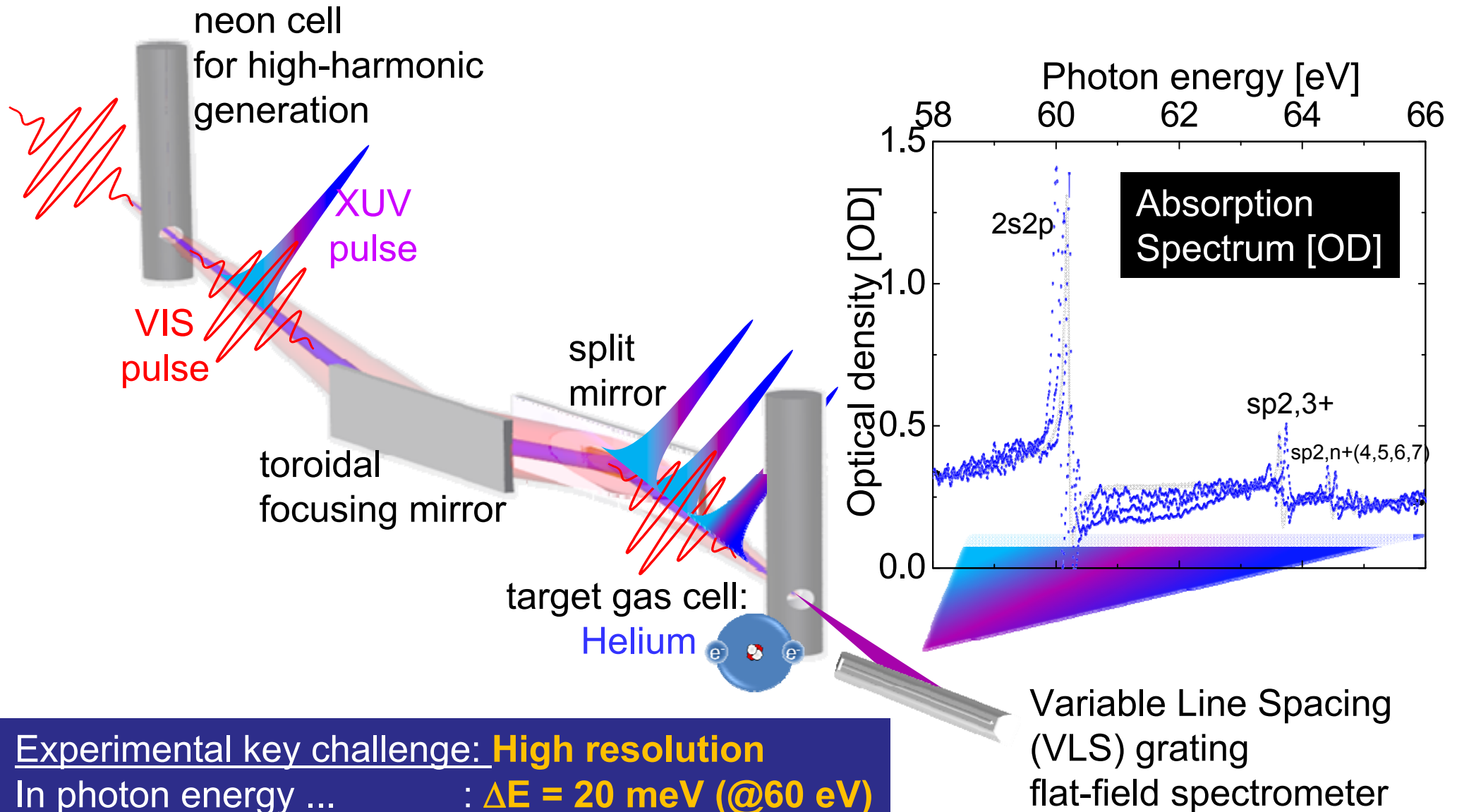
...

Experimental challenge:

- high (asec) temporal and
- high (meV) **spectral resolution** required at the same time

Experimental setup

for **time-resolved** XUV absorption spectroscopy



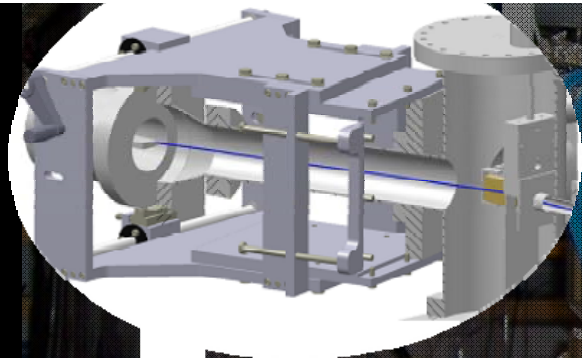
Experimental key challenge: **High resolution**

In photon energy ... : $\Delta E = 20 \text{ meV} (@60 \text{ eV})$

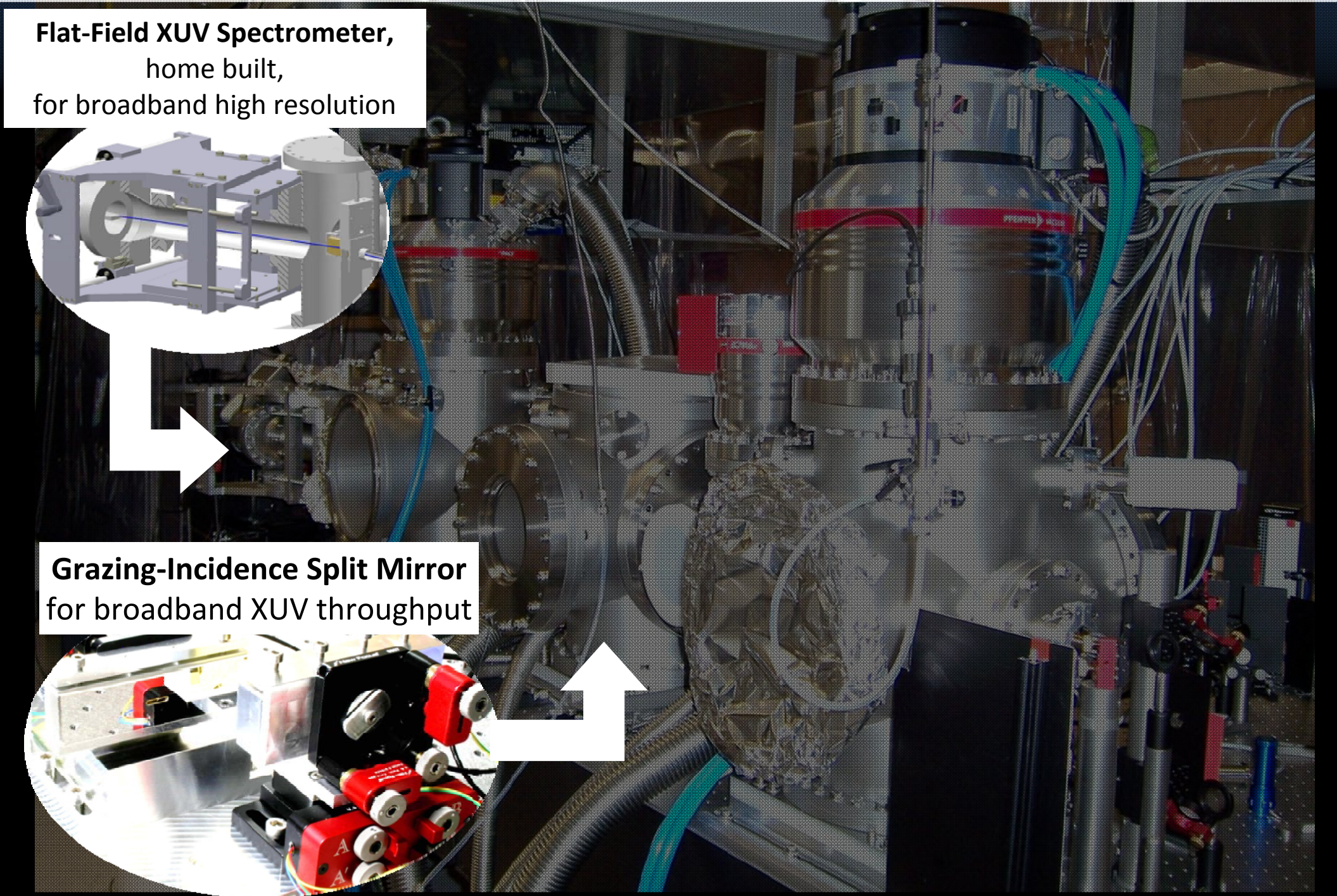
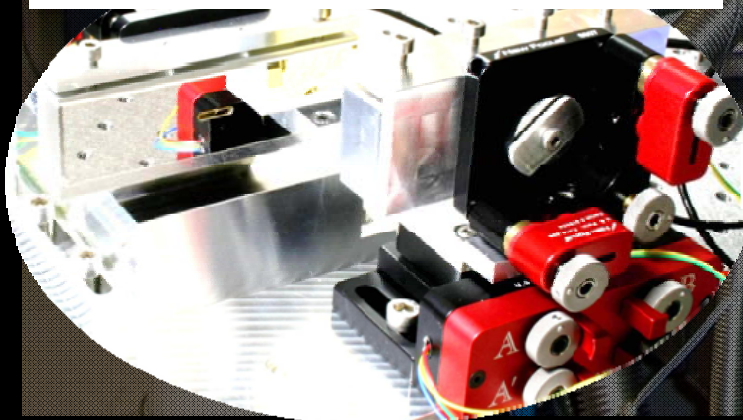
... and time delay : $\Delta \tau = 10 \text{ as}$

Experimental Setup in the Lab

Flat-Field XUV Spectrometer,
home built,
for broadband high resolution

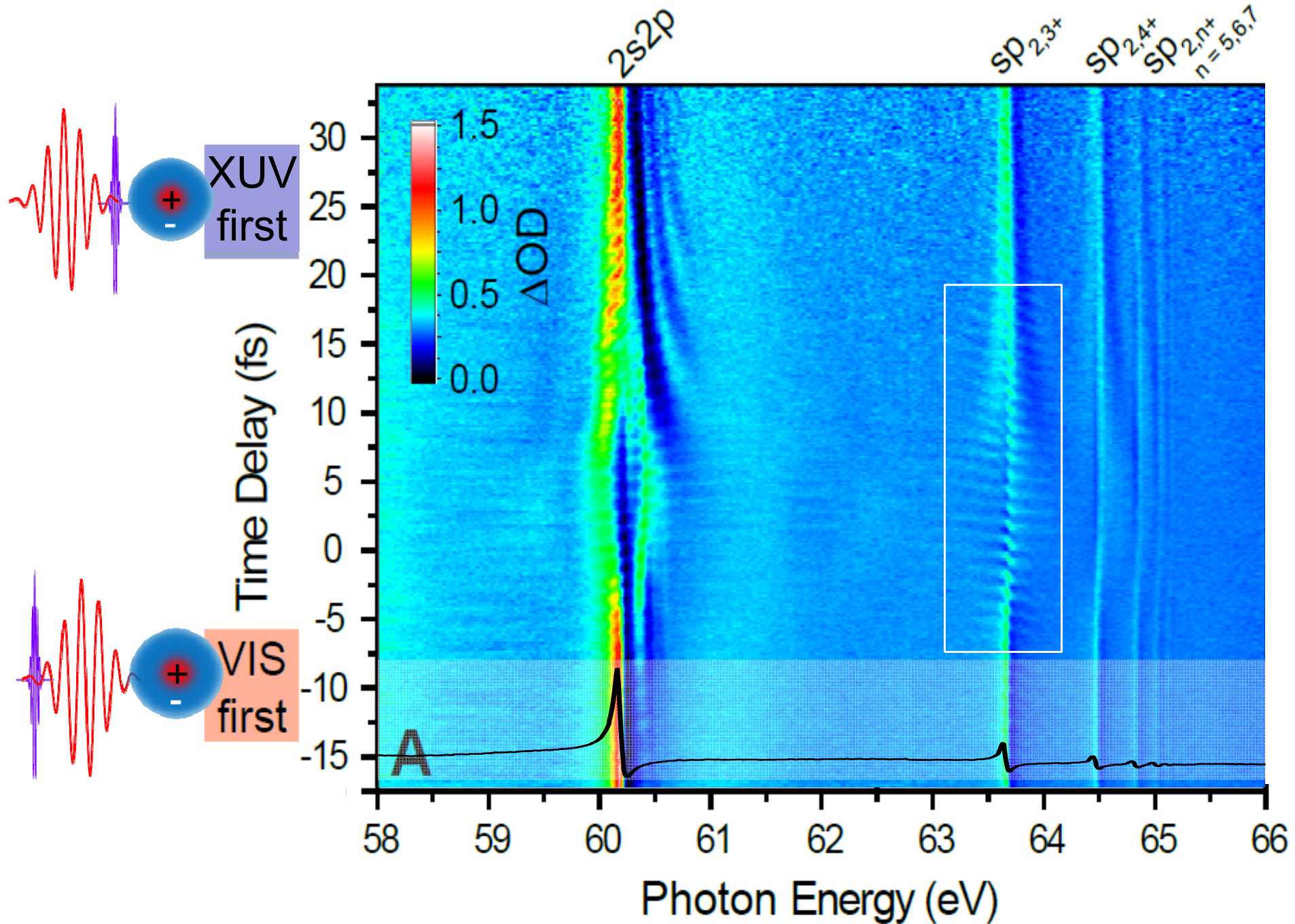


Grazing-Incidence Split Mirror
for broadband XUV throughput

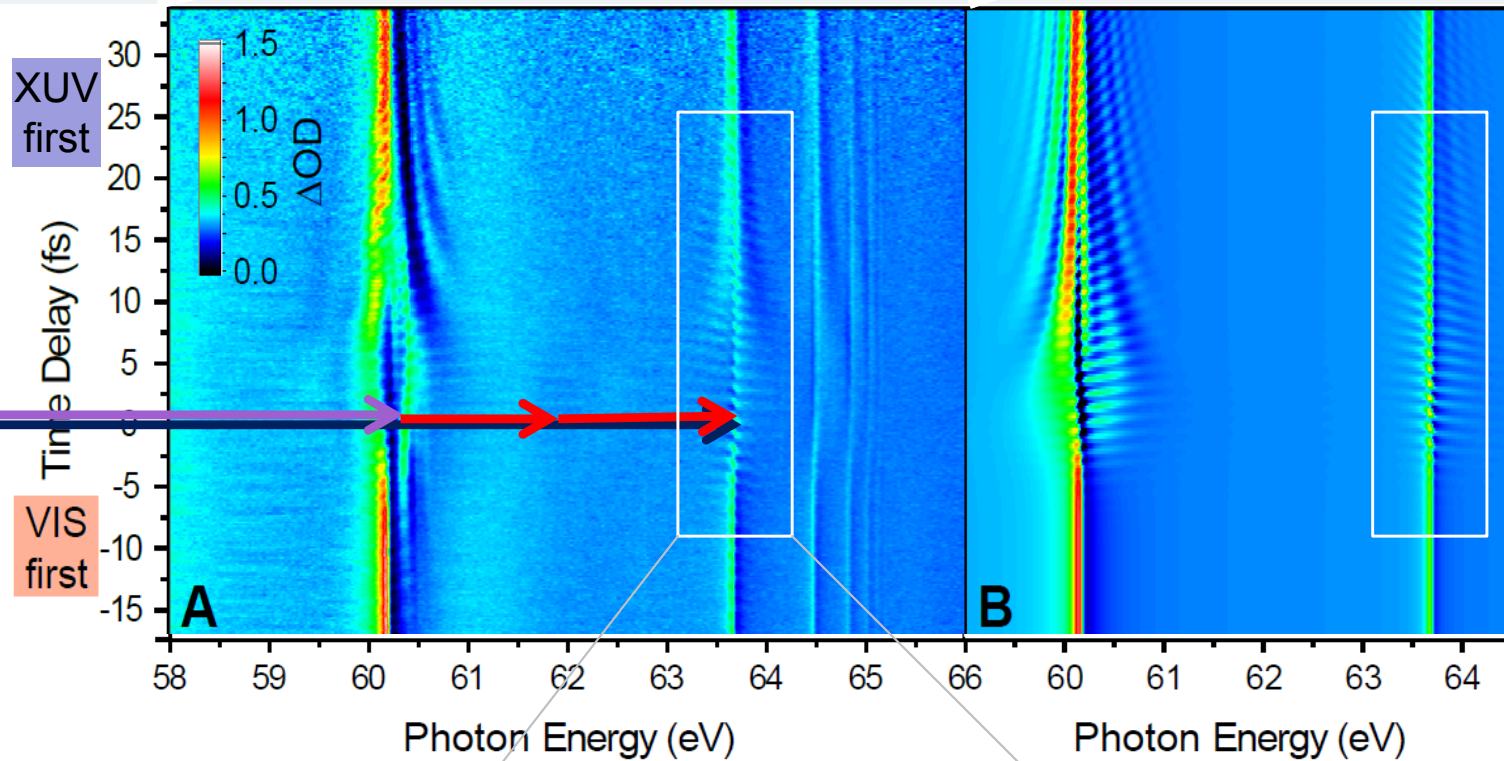


Time-resolved doubly-excited $2e^-$ dynamics in He

Experimental data



comparison experiment and theory

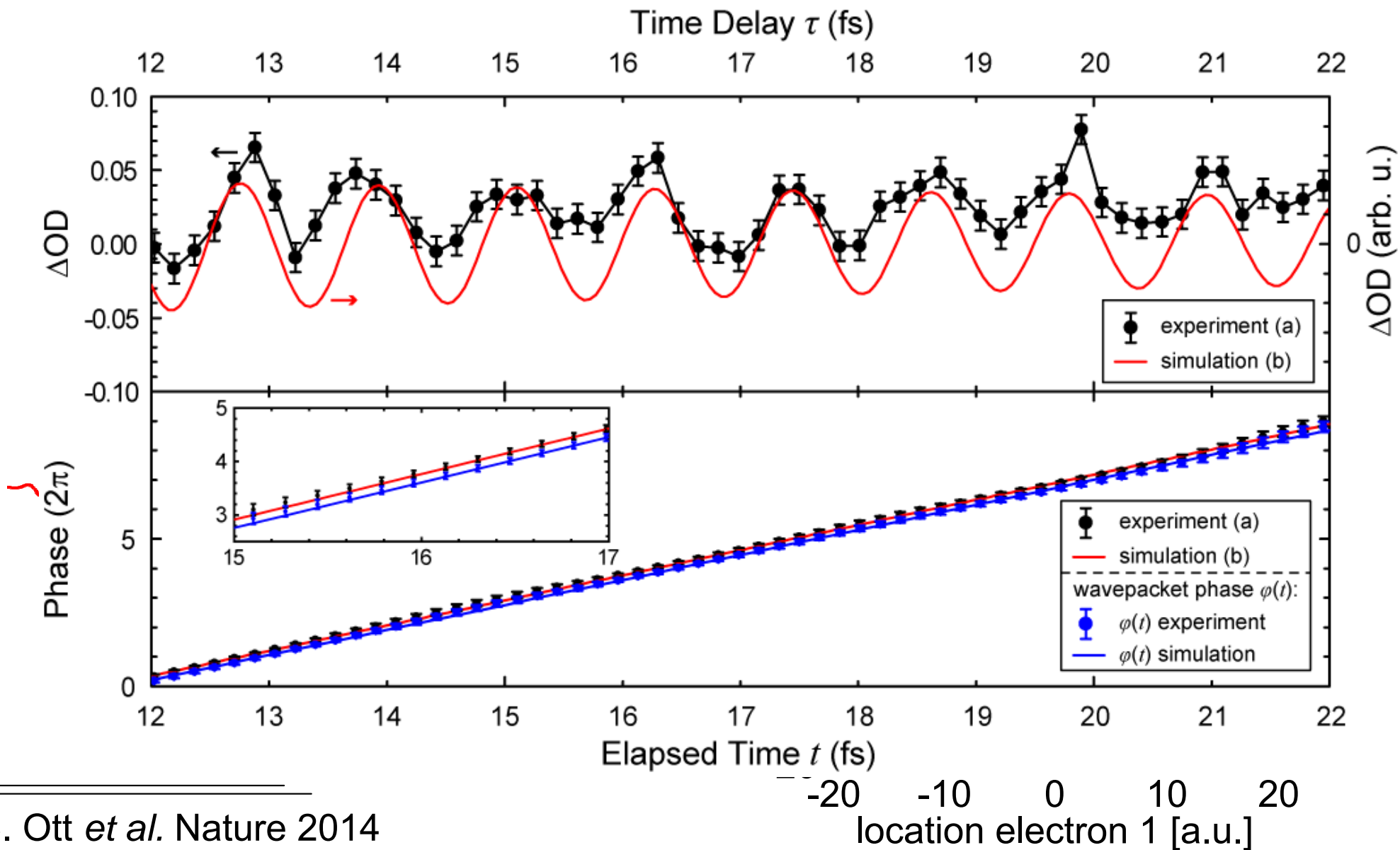


quantum-state interference
 \Rightarrow phase $\varphi(t)$

Probing a time-dependent
superposition state:
 \Rightarrow Wavepacket

Measuring the time-dependent phase difference of 2s2p and sp_{23+} autoionization states

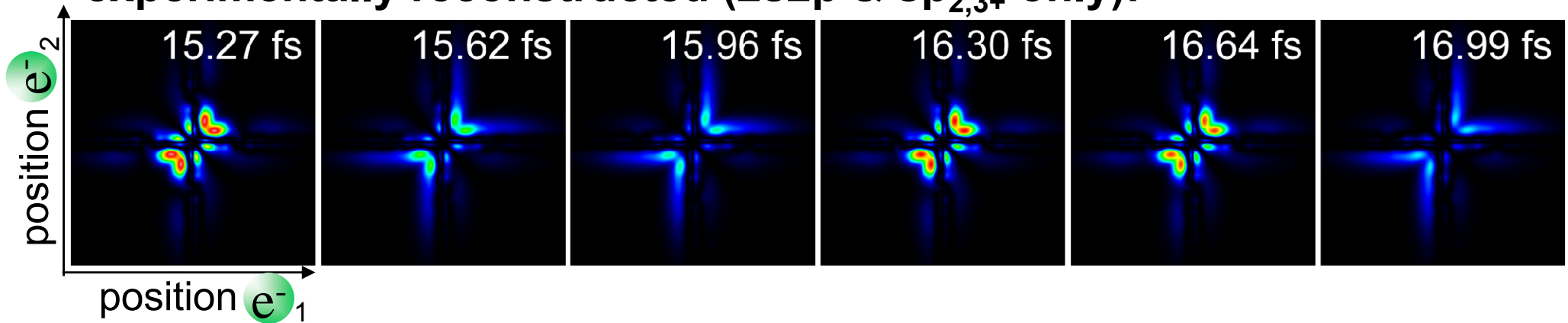
cooperation with Javier Madroño (Theory, TU München)
Luca Argenti, Fernando Martín (Theory, UAM Madrid)



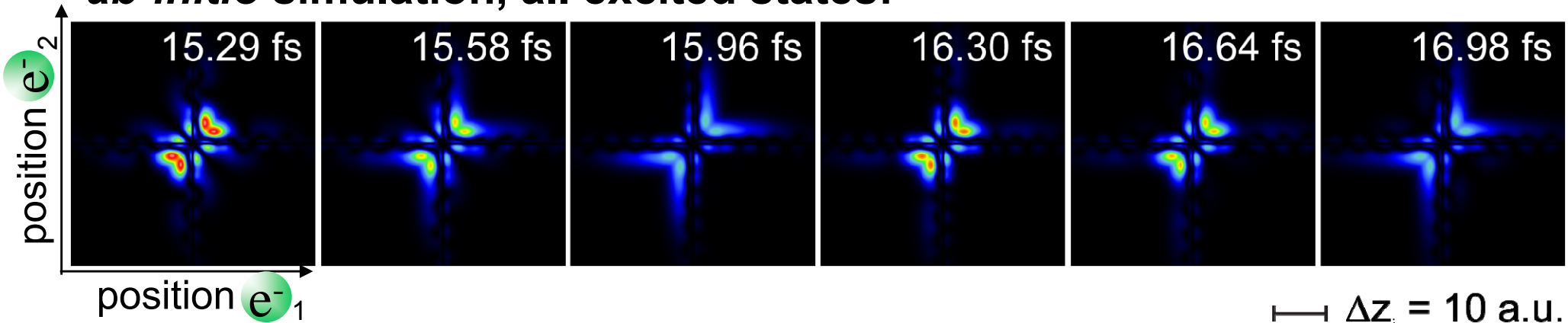
Testing *ab-initio* theory of e^- correlation dynamics

cooperation: Luca Argenti & Fernando Martín (UAM Madrid, Spain)
Javier Madroñero (TU Munich)

experimentally reconstructed ($2s2p$ & $sp_{2,3+}$ only):

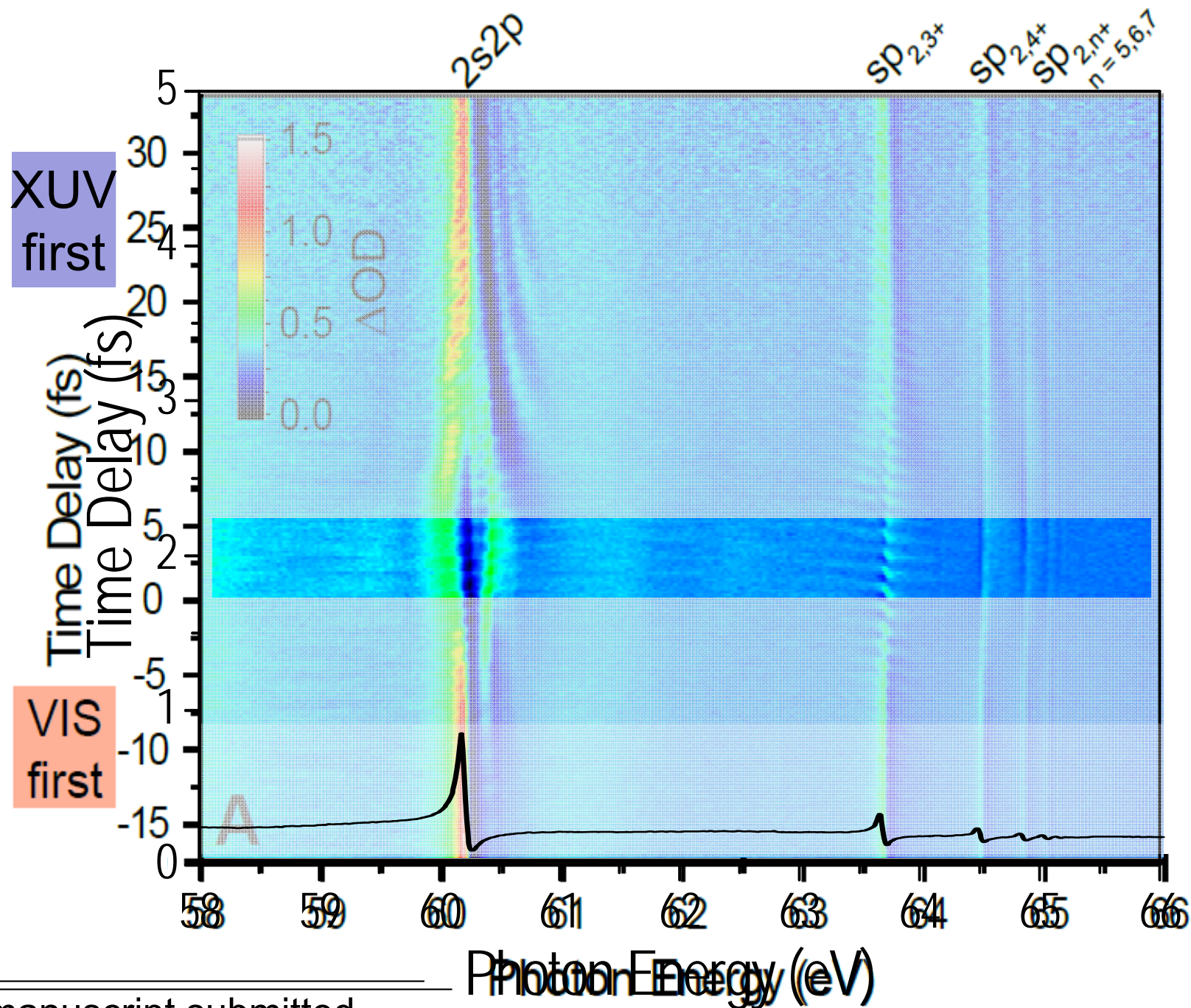


ab-initio simulation, all excited states:

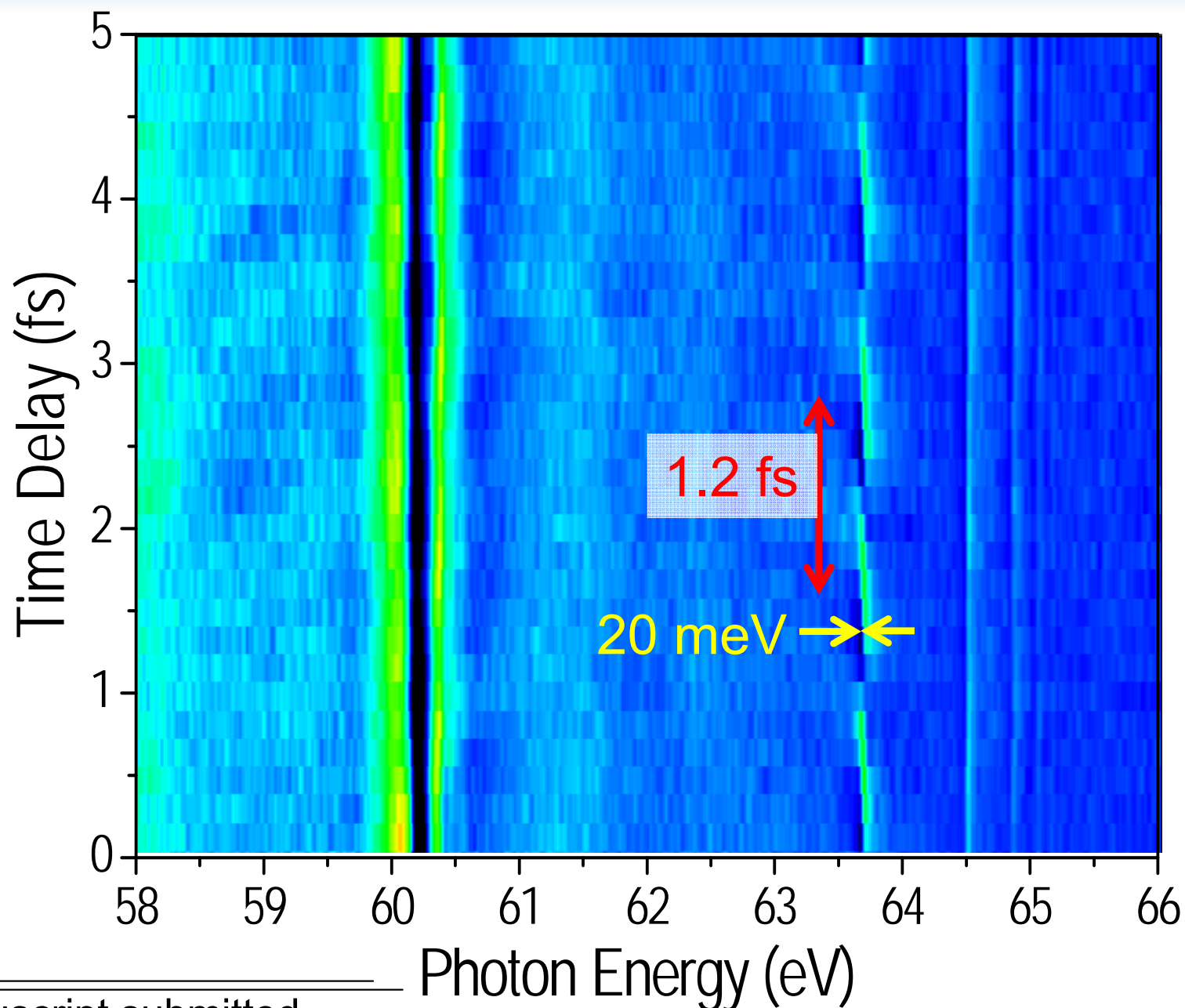


first experimental observation
of **two-electron wavepacket** motion

Time-resolved doubly-excited $2e^-$ dynamics in He



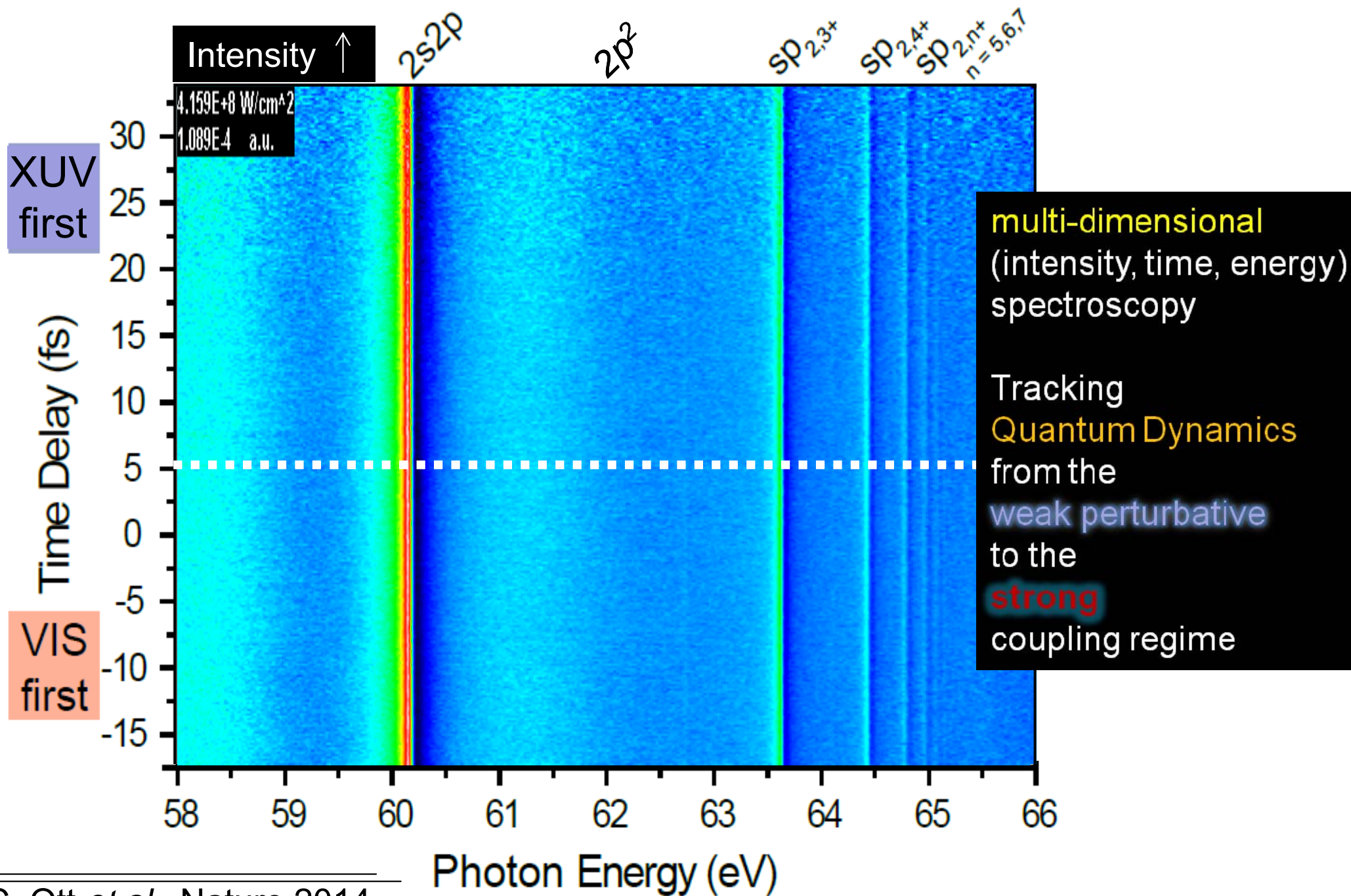
Time-resolved doubly-excited $2e^-$ dynamics in He



high **temporal**
and
high **spectral**
resolution
are
required
simultaneously

Intensity dependence of 2-e⁻ quantum dynamics

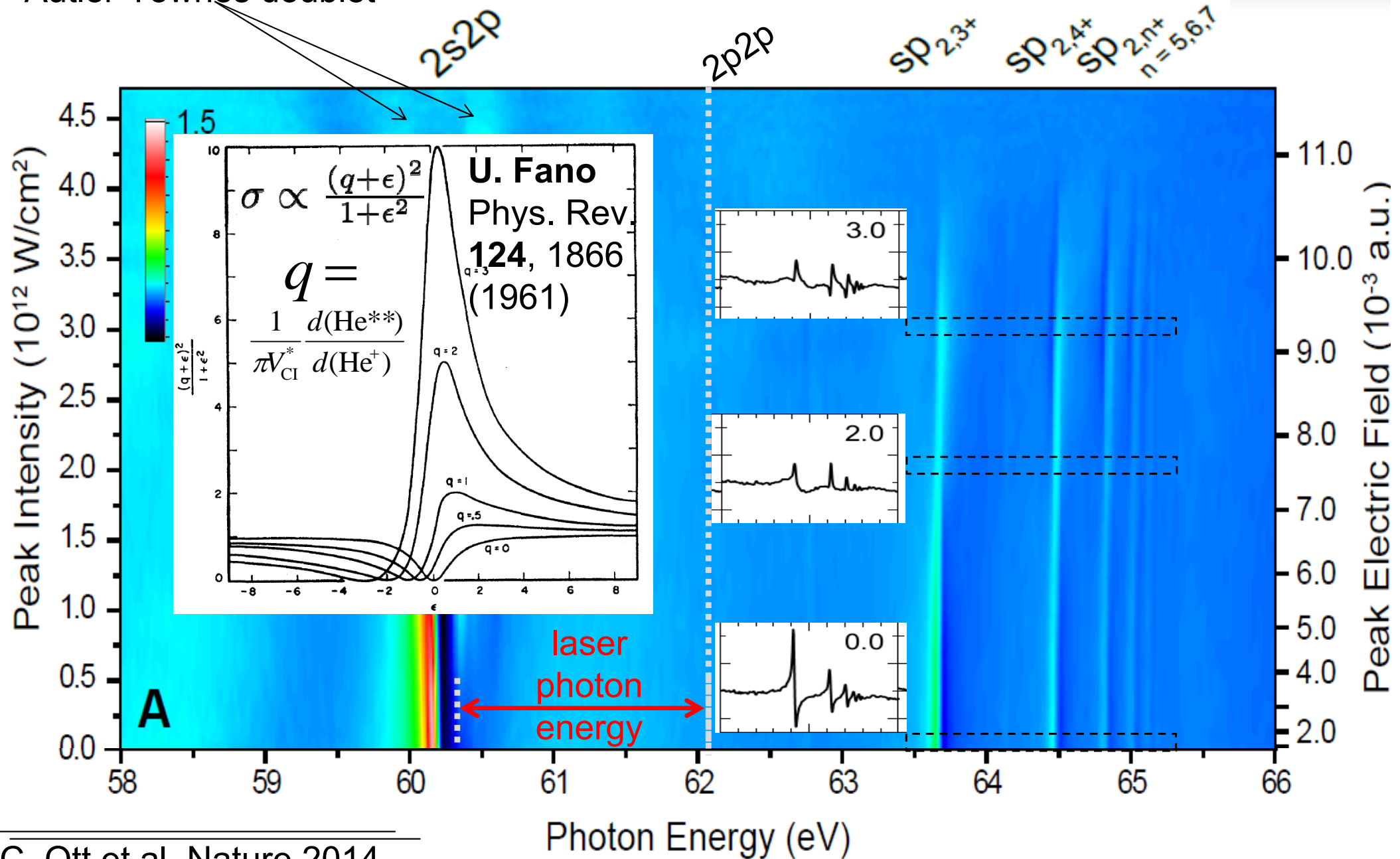
Experimental data



Intensity, a key parameter (the coupling strength)

Rabi cycling of autoionizing states)
Autler-Townes doublet

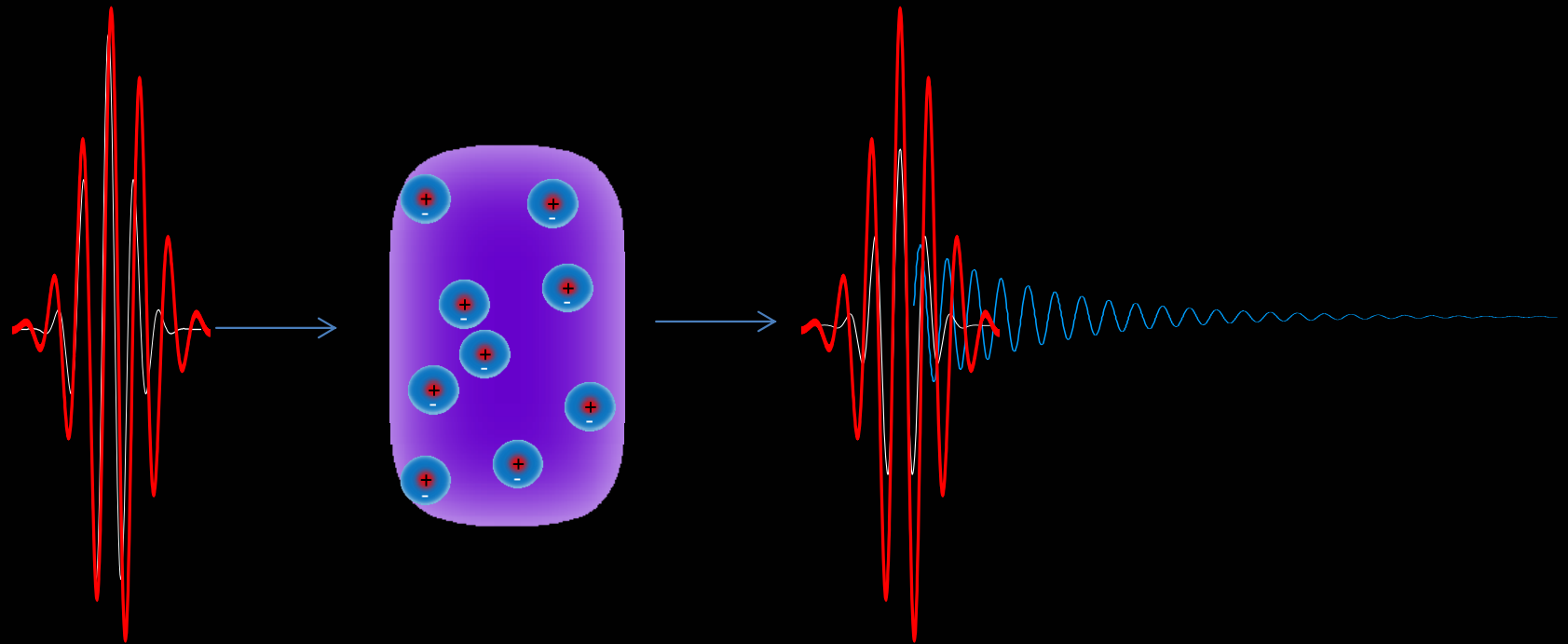
tuning from weak perturbative to strong fields



asking two e^- electrons in He
some quick questions ...

What is absorption?

And how does it respond to **intense pulsed light**?



The refractive index n

$$E(x,t) = E_0 e^{i(k_{\text{med}}x - \omega t)} \quad k_{\text{med}} = \frac{\omega}{c_{\text{med}}} = n \cdot \frac{\omega}{c_{\text{vac}}} = n k_{\text{vac}}$$

$$\hookrightarrow E = E_0 e^{i(n k_{\text{vac}}x - \omega t)} \quad n = \text{Re}(n) + i \text{Im}(n)$$

$$\hookrightarrow E = E_0 e^{i(\text{Re}(n)k_{\text{vac}}x - \omega t)} e^{-\text{Im}(n)k_{\text{vac}}x}$$

\swarrow macro. polarization rel. permittivity E_0
 $(\text{Im}(n)k_{\text{vac}})^{-1}$ x

$$P = \chi \cdot E$$

\swarrow susceptibility (linear) $\chi = \epsilon_r - 1$

$$[D = \epsilon_0 \epsilon_r E = P + \epsilon_0 E]$$

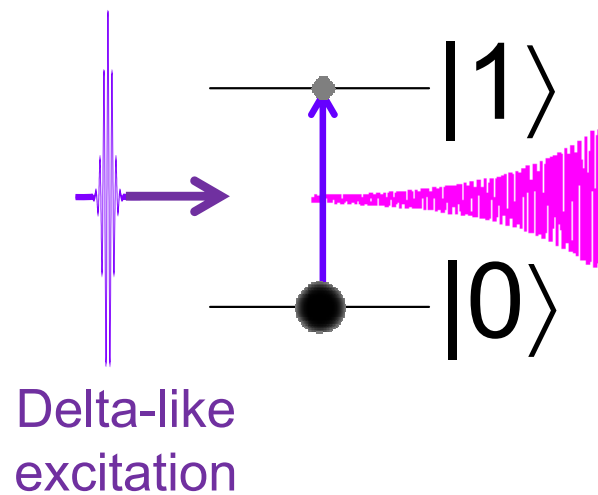
$$P = \int \rho d$$

\swarrow dipole moment
 \swarrow density of dipoles (e.g. atoms)

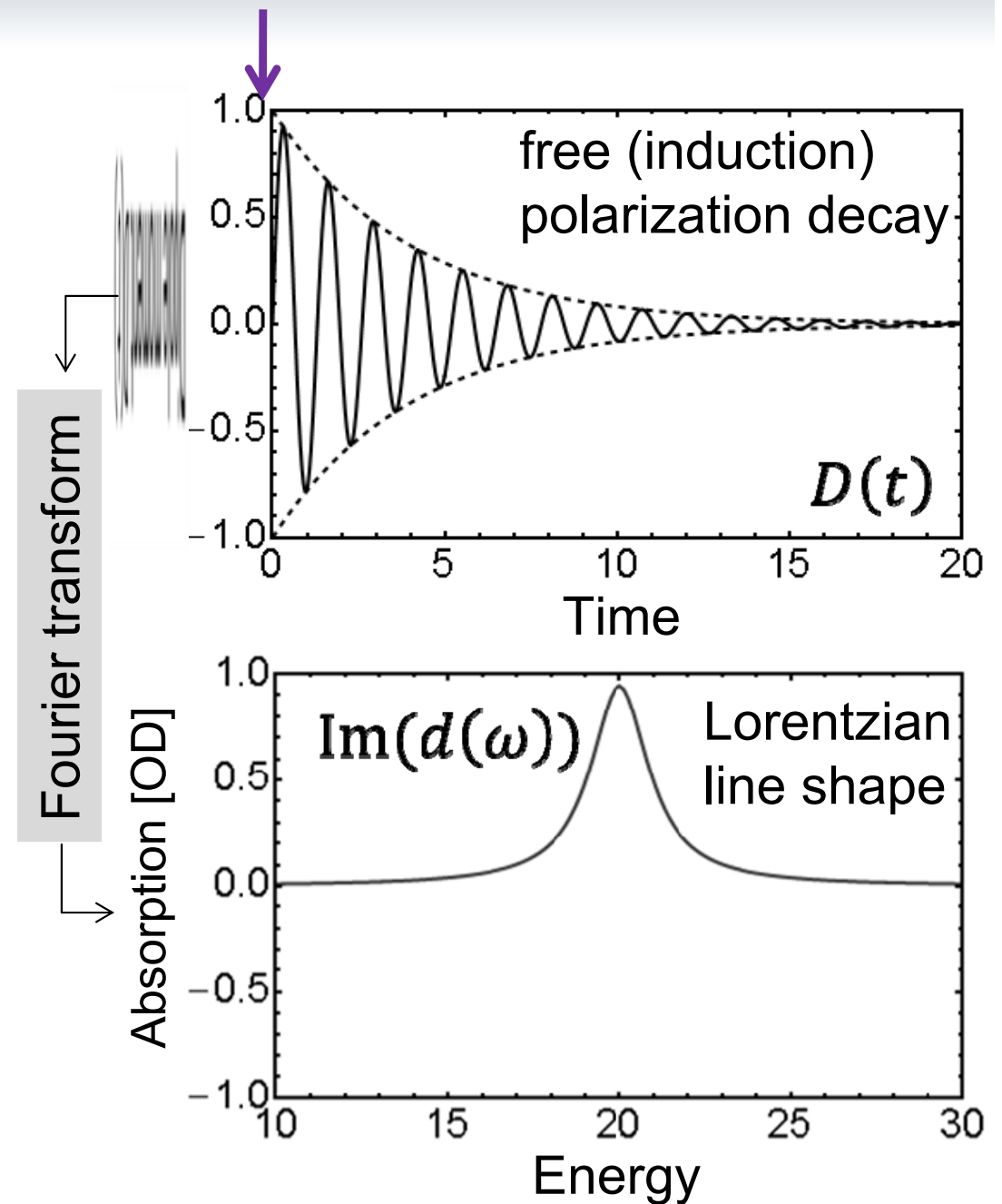
$$n^2 = \epsilon_r = 1 + \chi$$

$$n = \sqrt{1 + \chi} \sim 1 + \frac{\chi}{2} = 1 + \frac{\int \rho d}{E} \Rightarrow \Delta n(\omega) = \int \frac{d(\omega)}{E(\omega)}$$

Optical response and absorption

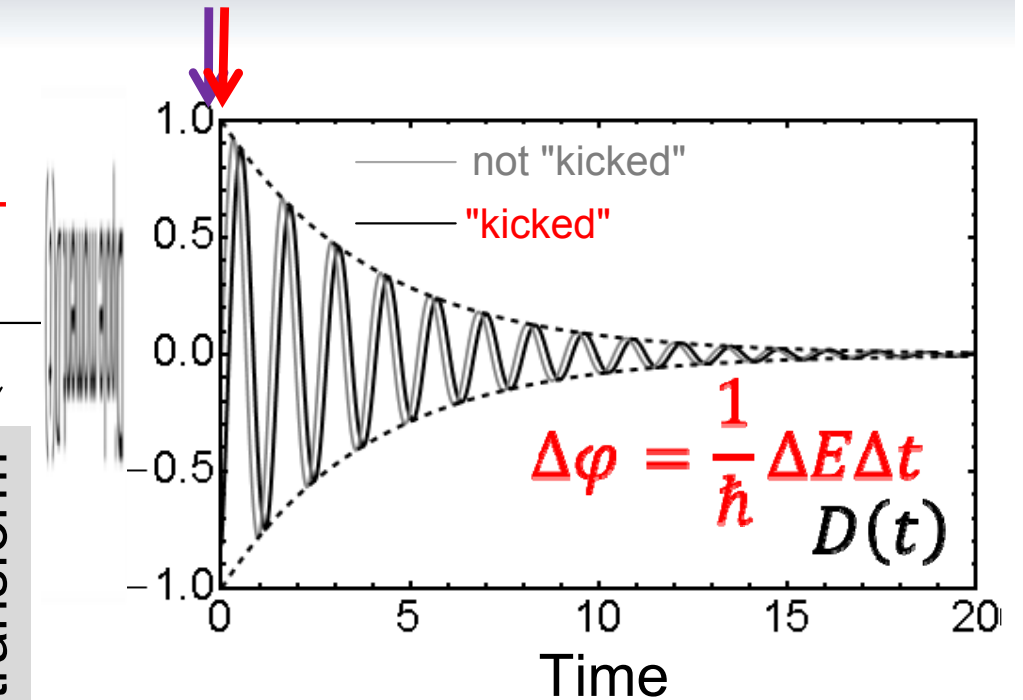
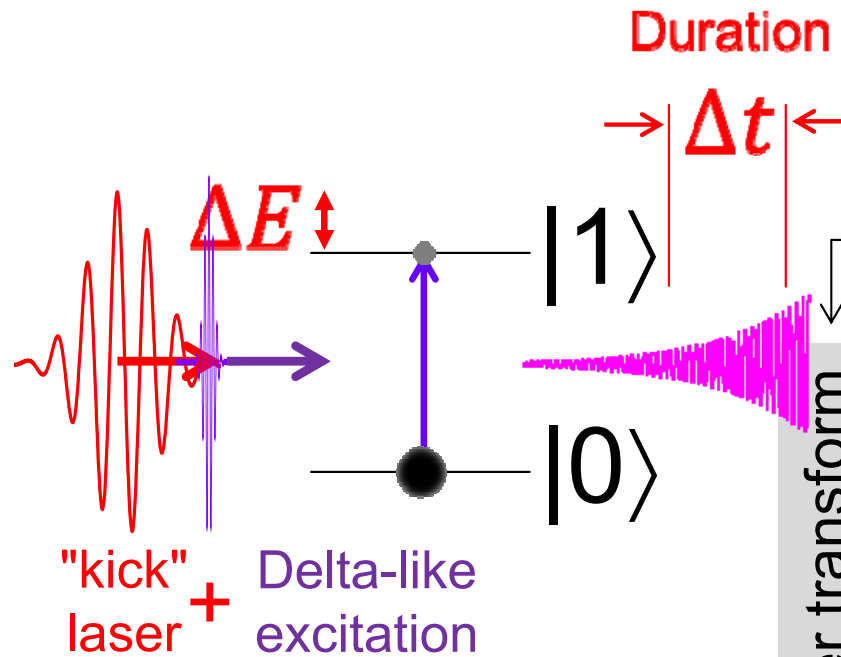


$$\Rightarrow \Delta n(\omega) = \int \frac{d(\omega')}{E(\omega')}$$



Resonance absorption in the Time Domain

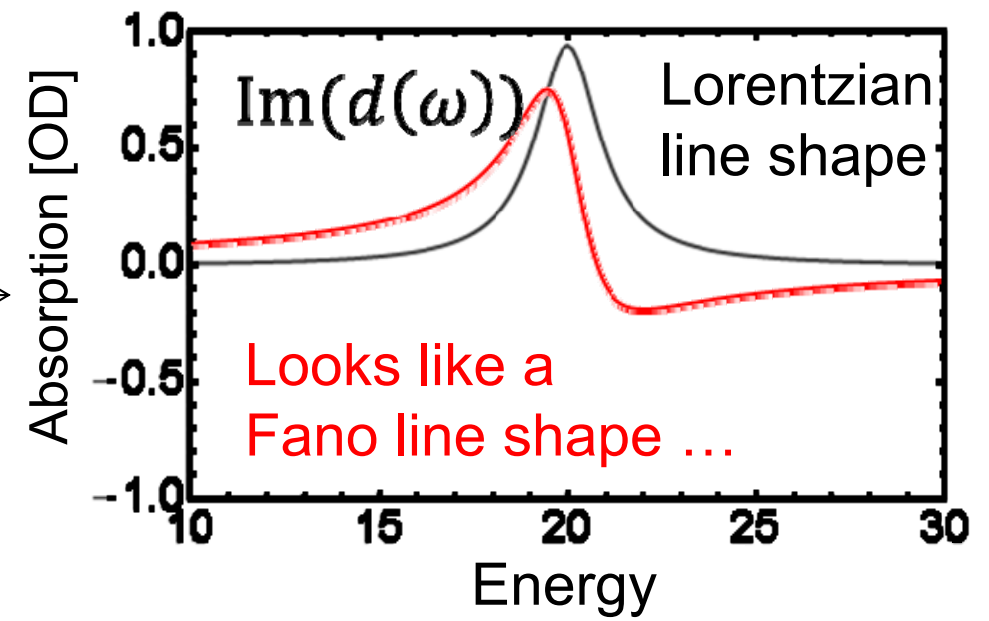
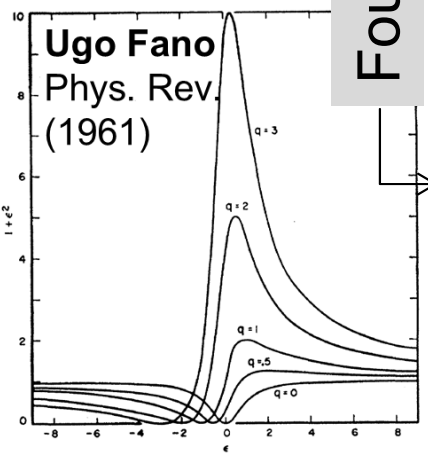
Science **340**, 716 (2013)



...it IS a Fano line shape!

$$\varphi(q) \Rightarrow \text{Im}(d(\omega)) = \frac{q^2}{1+q^2}$$

$$q(\varphi) = \cot\left(\frac{\varphi}{2}\right)$$



The Fano dipole phase

Exact mapping from Fano q parameter to temporal phase shift φ

cooperation with C. Greene (Purdue), J. Evers and C. Keitel

$$\sigma_{Fano} \sim \frac{(\varepsilon + q)^2}{\varepsilon^2 + 1}$$

=

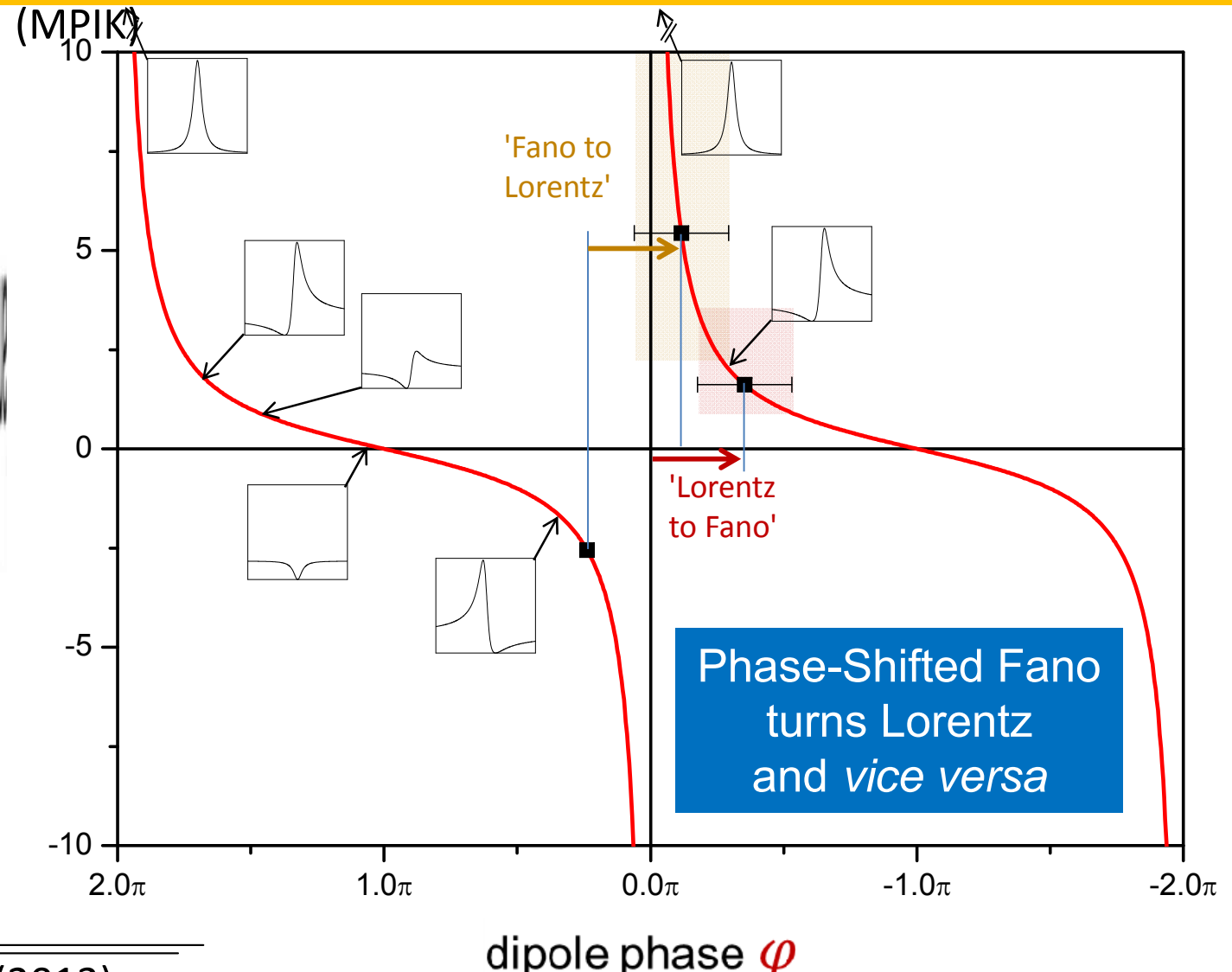
$$\text{Im}\left(\frac{-1}{i+\varepsilon} \exp(i\varphi)\right) + \text{const.}$$

Phase-shifted Lorentzian

...it IS a
Fano line shape!

$$\varphi(q) = 2 \arg(q - i)$$

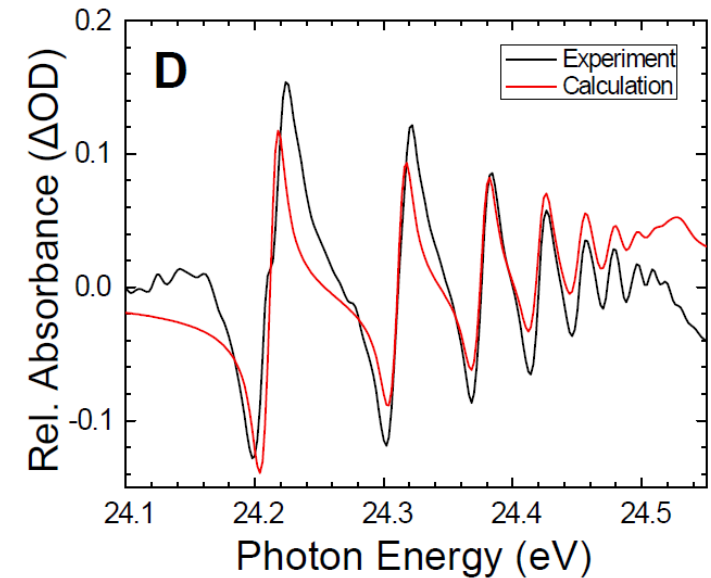
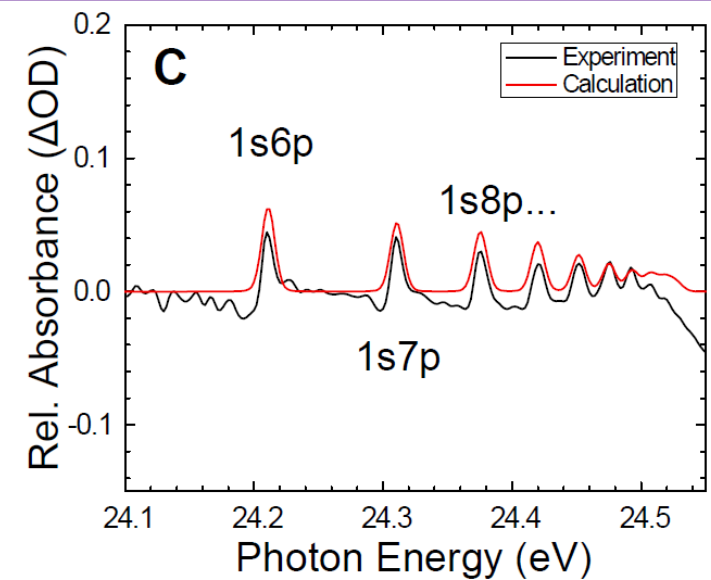
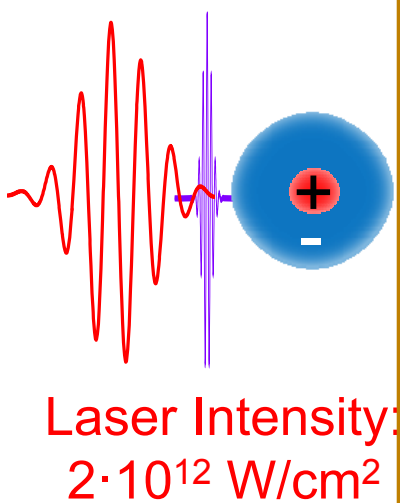
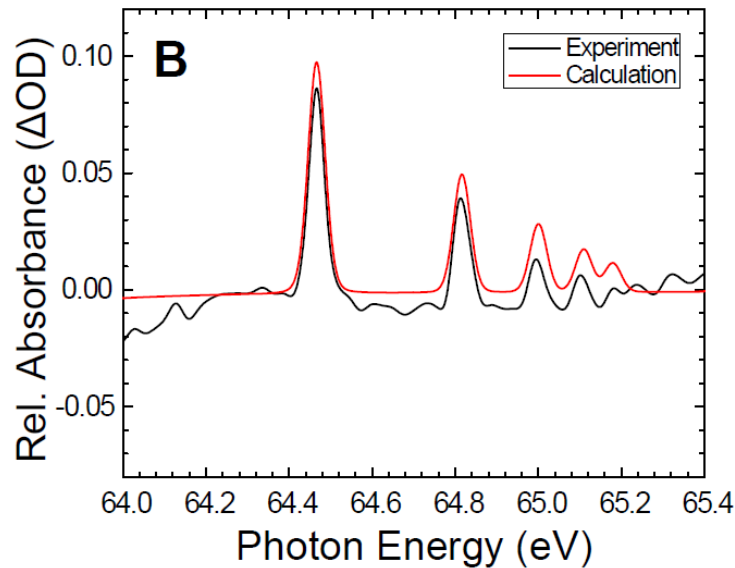
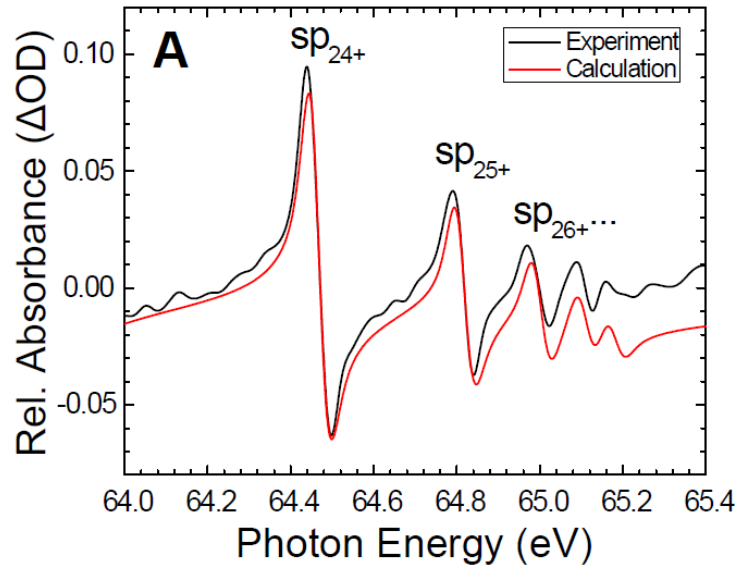
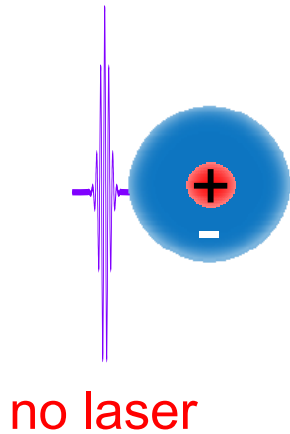
$$q(\varphi) = -\cot\left(\frac{\varphi}{2}\right)$$



Fano to Lorentz, and Lorentz to Fano

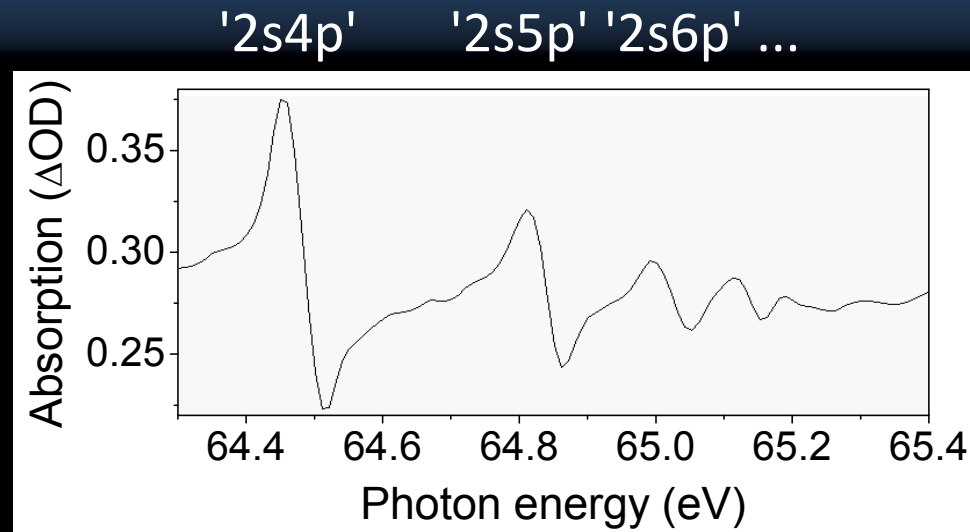
doubly-excited Helium
originally Fano lineshape

singly-excited Helium
originally Lorentzian



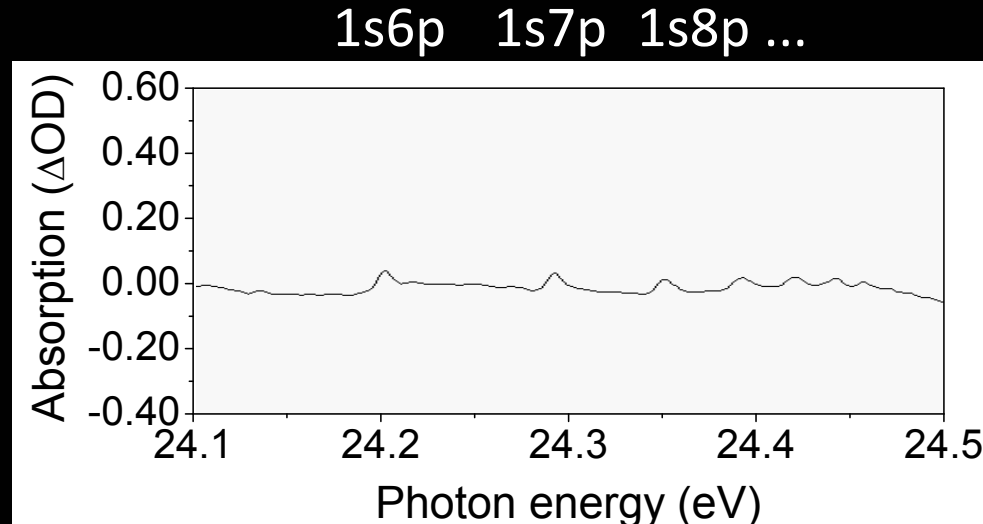
Changing the (spectral) shape of atoms

Helium
doubly excited
 (above the
 continuum
 threshold)



turning
 original 'Fano'
 into 'Lorentz'

Helium
singly excited
 (below the
 continuum
 threshold)



turning
 original 'Lorentz'
 into 'Fano'

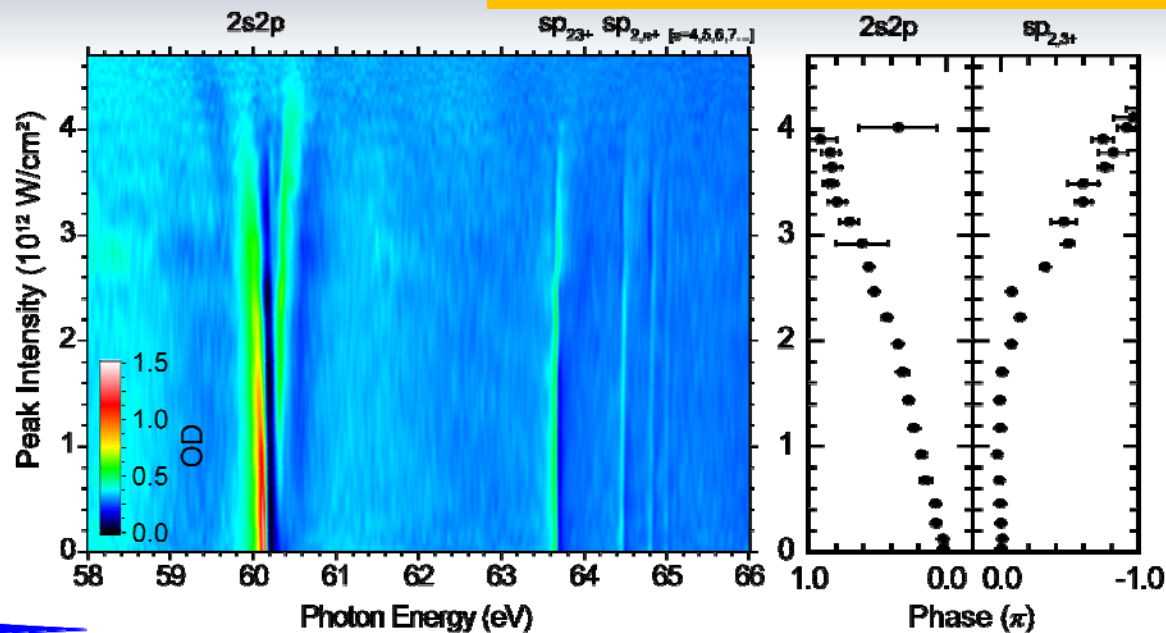
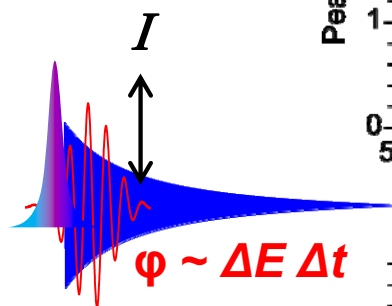
$$q = -\cot\left(\frac{\varphi}{2}\right)$$

$$q = -2.55 \Rightarrow \varphi = 0.24\pi$$

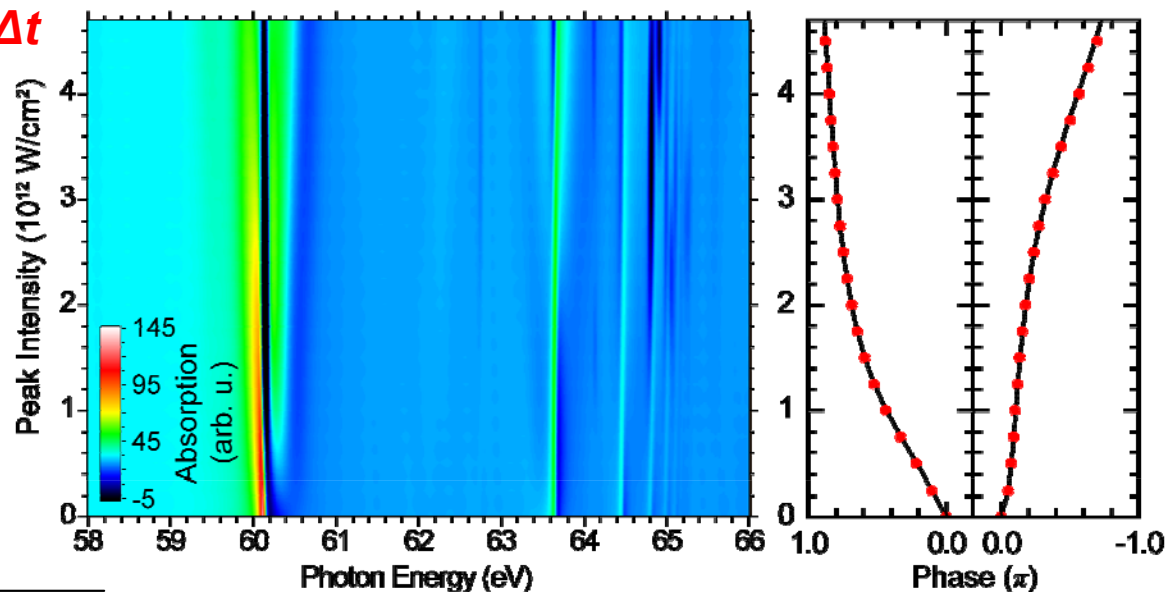
Extracting the laser-induced phase shift

Cooperation with J. Madronero, L. Argenti, F. Martín

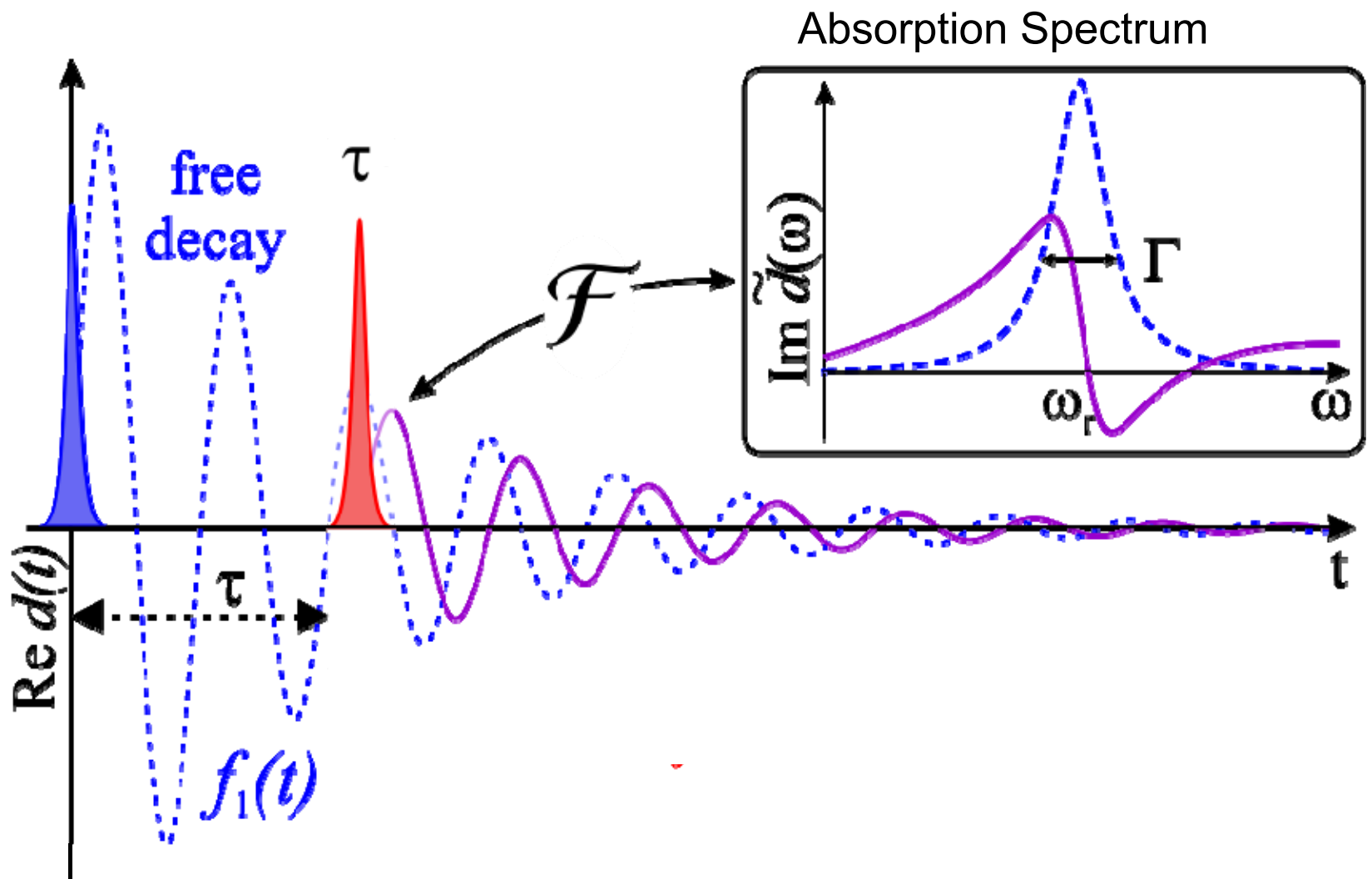
experiment



ab-initio theory
(Argenti/Martin)

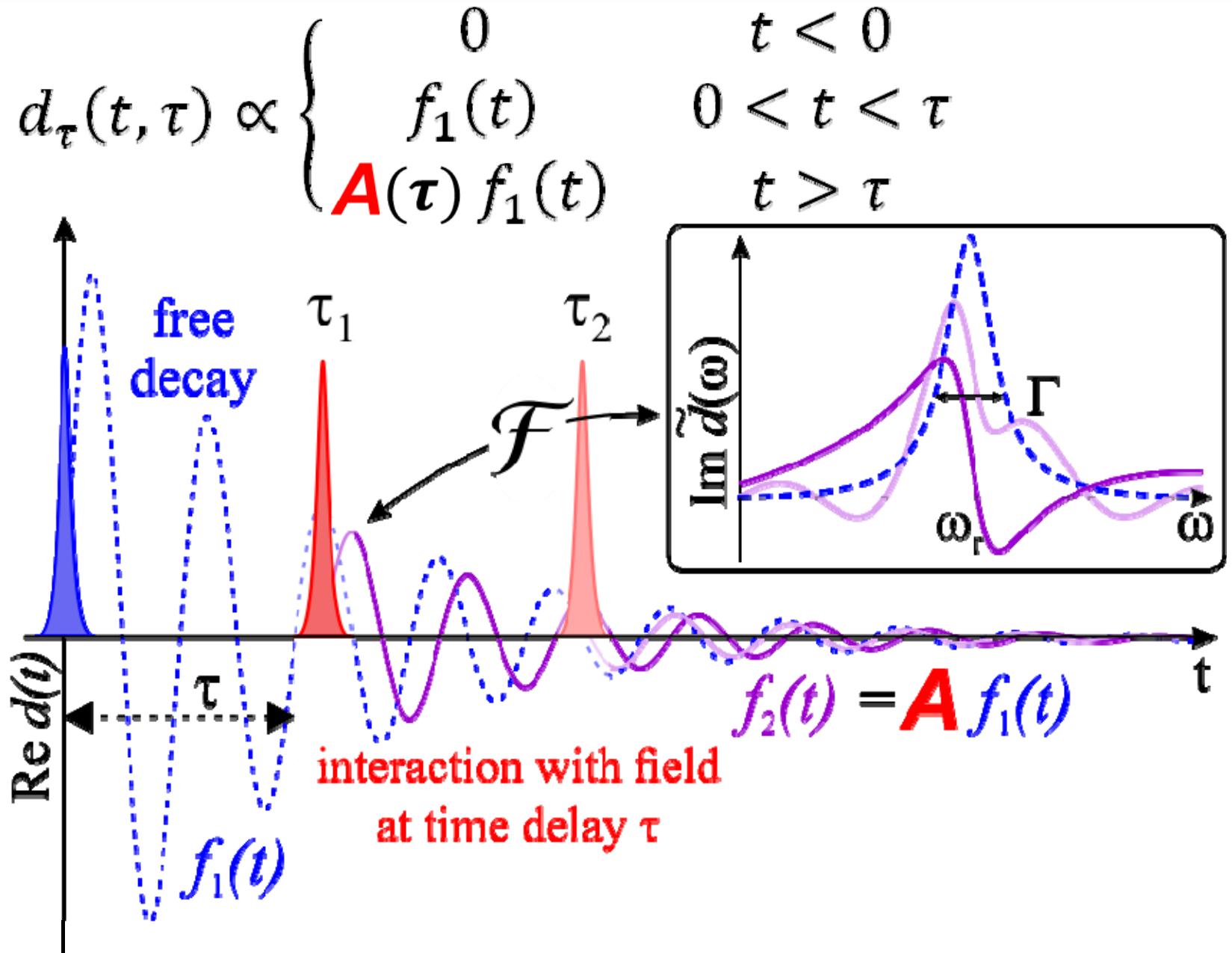


The Fano Phase Shift



The dipole-control model

Blättermann *et al.* J. Phys. B: At. Mol. Opt. Phys. 47 124008 (2014)

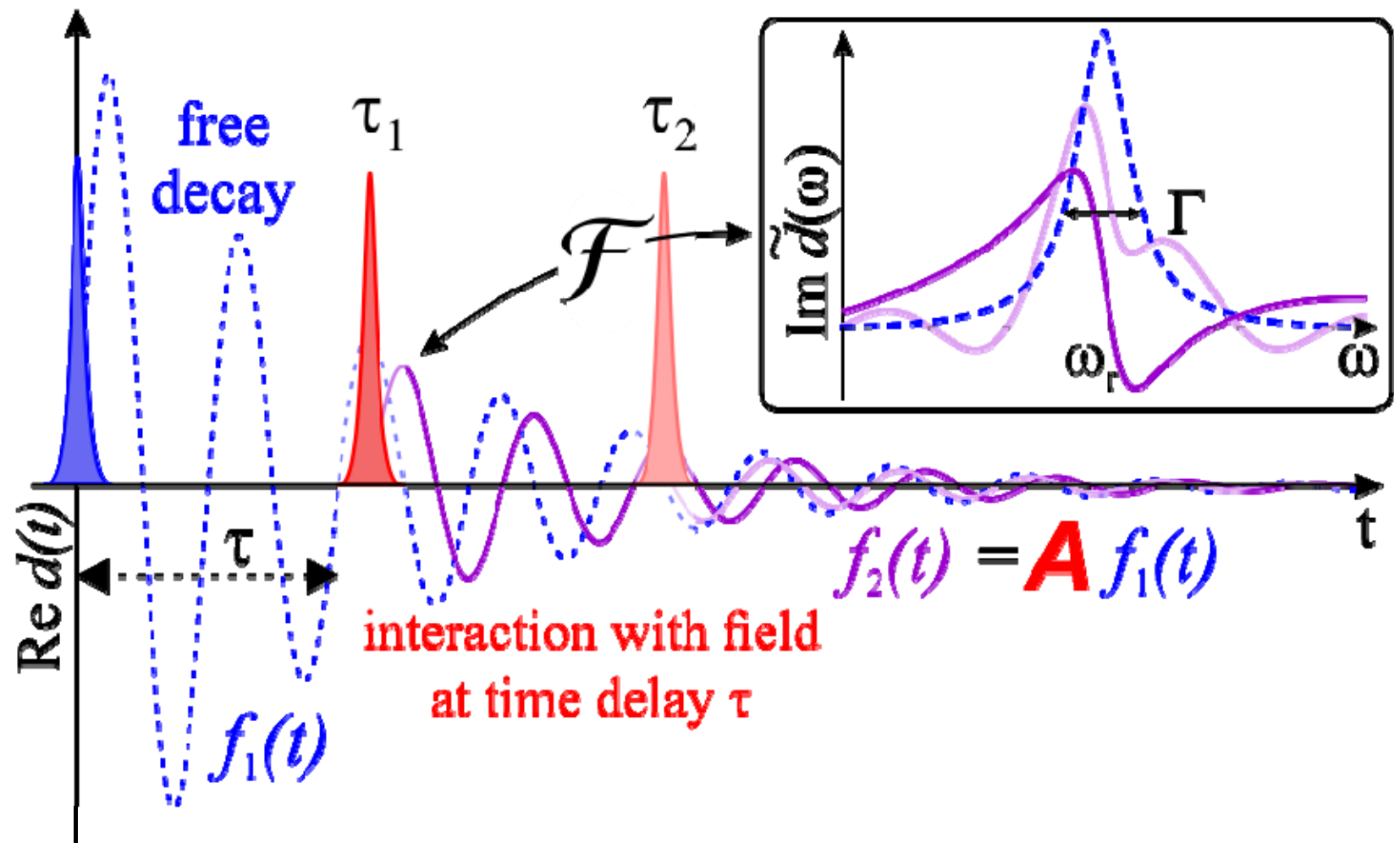


The dipole-control model

Analytical result:

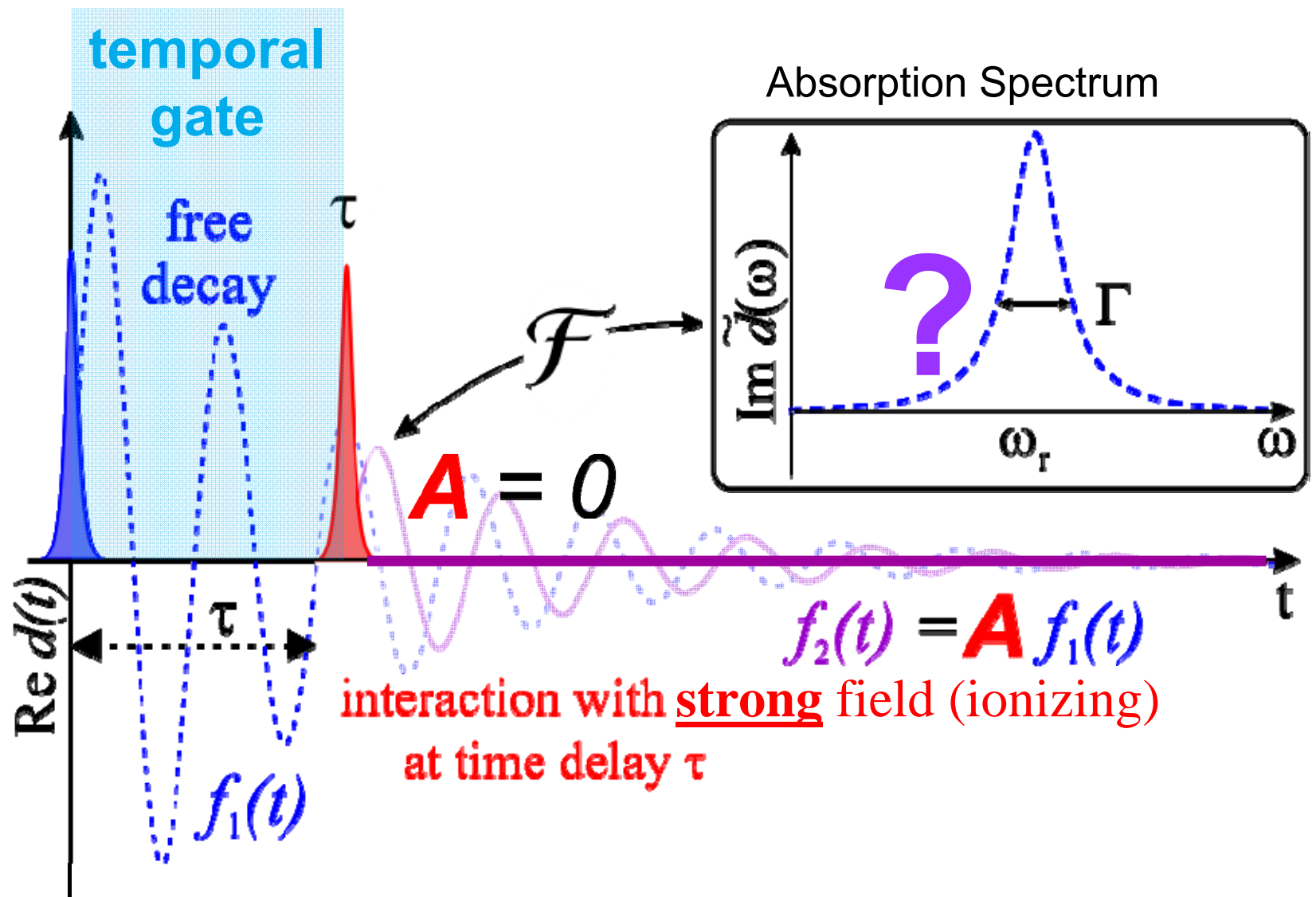
Blättermann *et al.* J. Phys. B: At. Mol. Opt. Phys. **47** 124008 (2014)

$$\tilde{d}_\tau(\omega, \tau) \propto -i \frac{1 - e^{i(\omega_r - \omega)\tau - \frac{\Gamma}{2}\tau} (1 - \mathbf{A}(\tau))}{i(\omega_r - \omega) - \Gamma/2}$$



The Dipole-Amplitude Gate

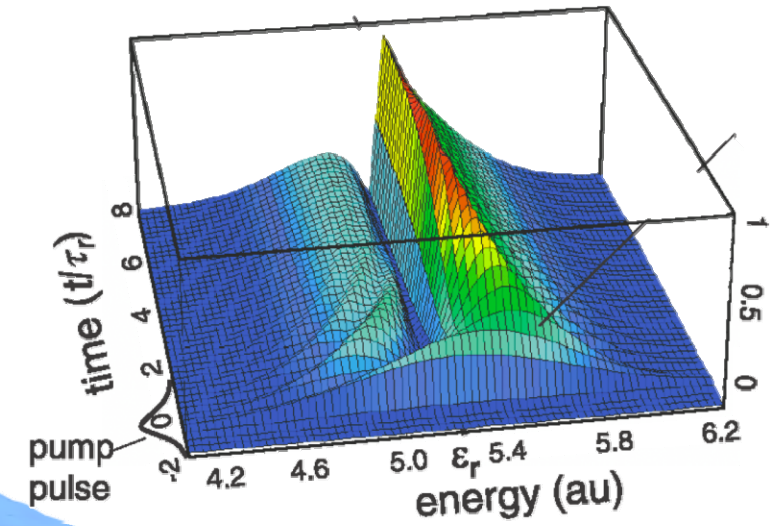
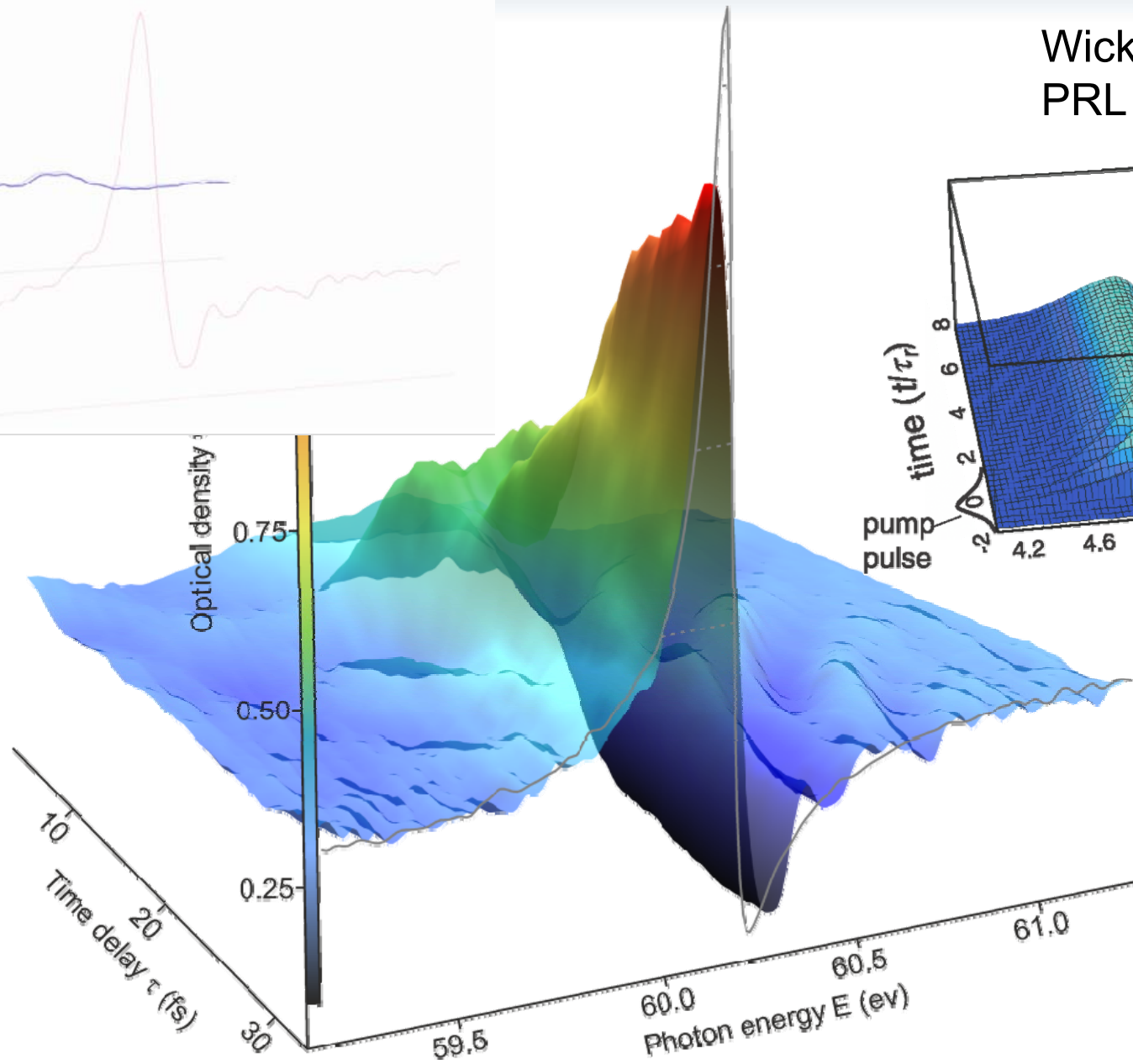
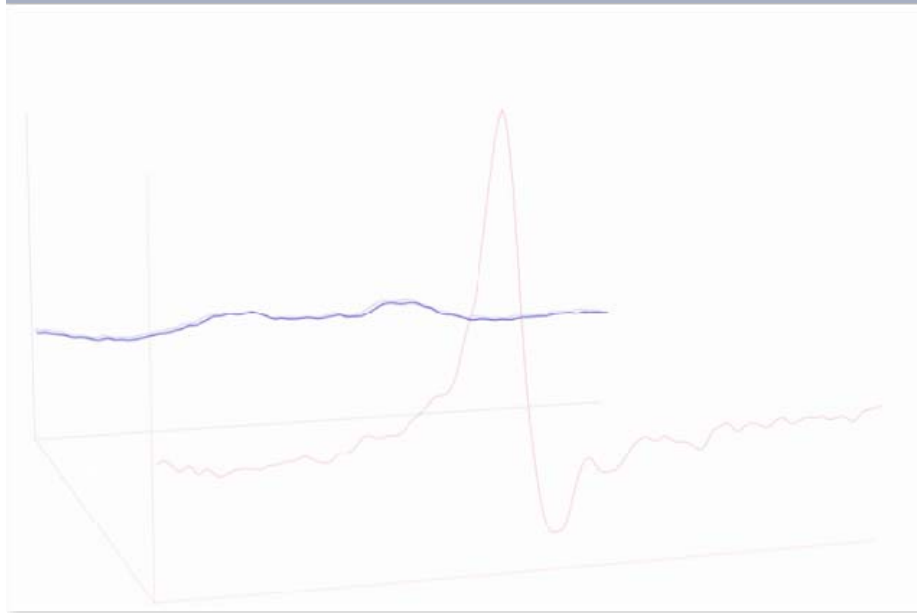
Kaldun, Blättermann *et al.* Science 2016



The birth of a Fano resonance

Kaldun, Blättermann *et al.* Science 2016

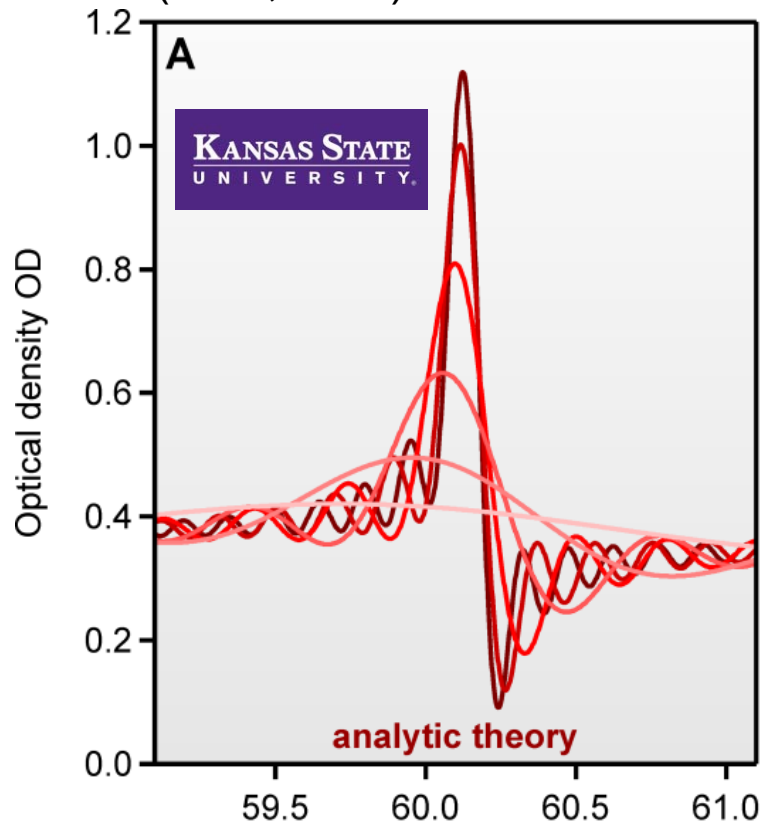
Wickenhauser *et al.*
PRL **94**, 023002 (2005)



From theory to experiment

Kaldun, Blättermann *et al.* Science 2016

by Hui Wei and
Chii-Dong Lin
(KSU, USA)

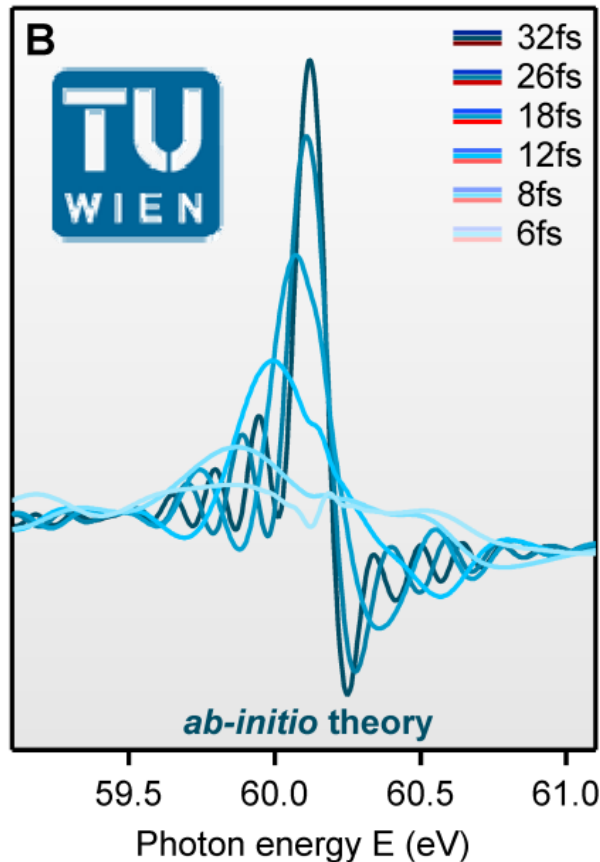


analytical result:

$$\sigma(\epsilon, \tau) \propto \text{Re} \left\{ 1 + \frac{(q - i)^2}{1 - i\epsilon} \left[1 - e^{-\frac{\Gamma}{2}(1 - i\epsilon)\tau} \right] \right\}$$

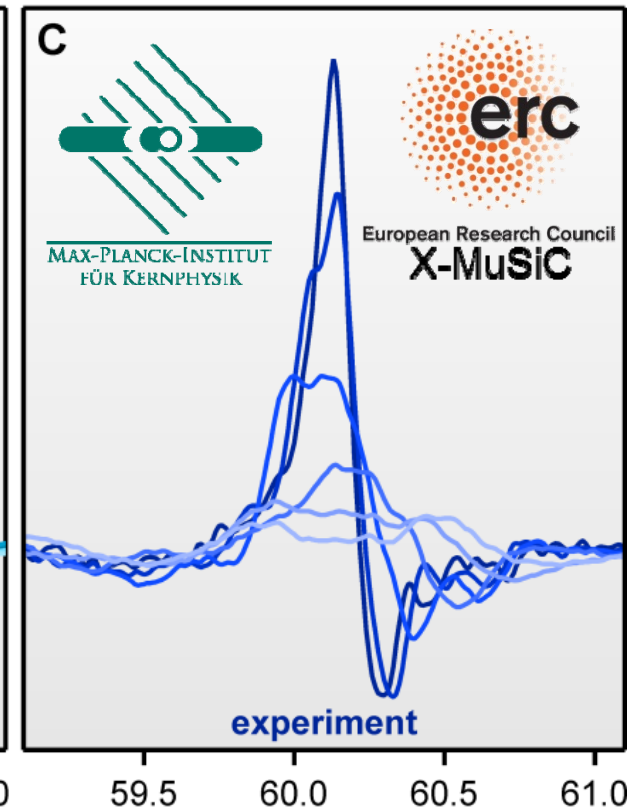
ideal case approximation

by R. Pazourek, S. Donsa
J. Burgdörfer
(TU Vienna, Austria)



TDSE
(TDCC)

Experiment
(MPIK, Heidelberg)



time-gated
absorption
spectroscopy

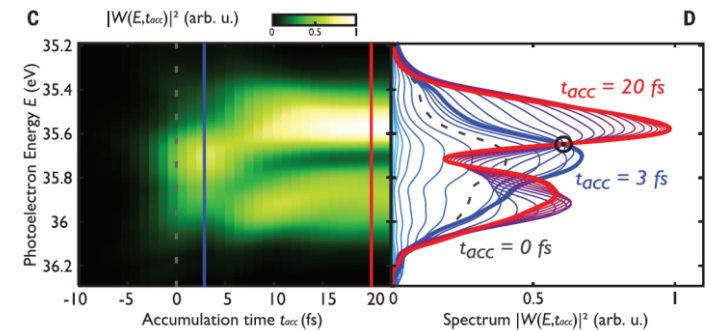
Two complementary studies

CHEMICAL PHYSICS

Attosecond dynamics through a Fano resonance: Monitoring the birth of a photoelectron

V. Gruson,^{1*} L. Barreau,^{1*} Á. Jiménez-Galan,² F. Risoud,³ J. Caillat,³ A. Maquet,³ B. Carré,¹ F. Lepetit,¹ J.-F. Hergott,¹ T. Ruchon,¹ L. Argenti,^{2†} R. Taïeb,³ F. Martín,^{2,4,5‡} P. Salières^{1‡}

Photoelectron spectroscopy



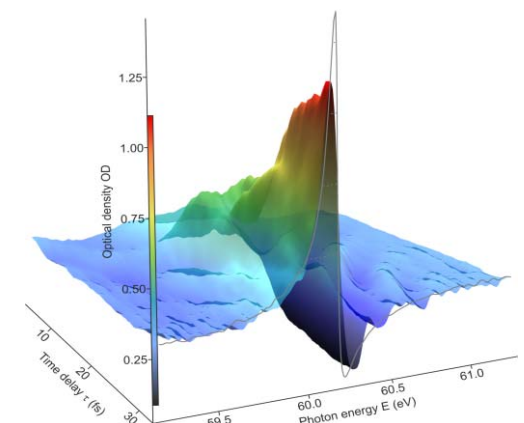
Science **354** 734 (2016)

CHEMICAL PHYSICS

Observing the ultrafast buildup of a Fano resonance in the time domain

A. Kaldun,^{1*†} A. Blättermann,^{1†} V. Stooß,¹ S. Donsa,² H. Wei,³ R. Pazourek,² S. Nagele,² C. Ott,¹ C. D. Lin,³ J. Burgdörfer,² T. Pfeifer^{1,4‡}

Absorption spectroscopy



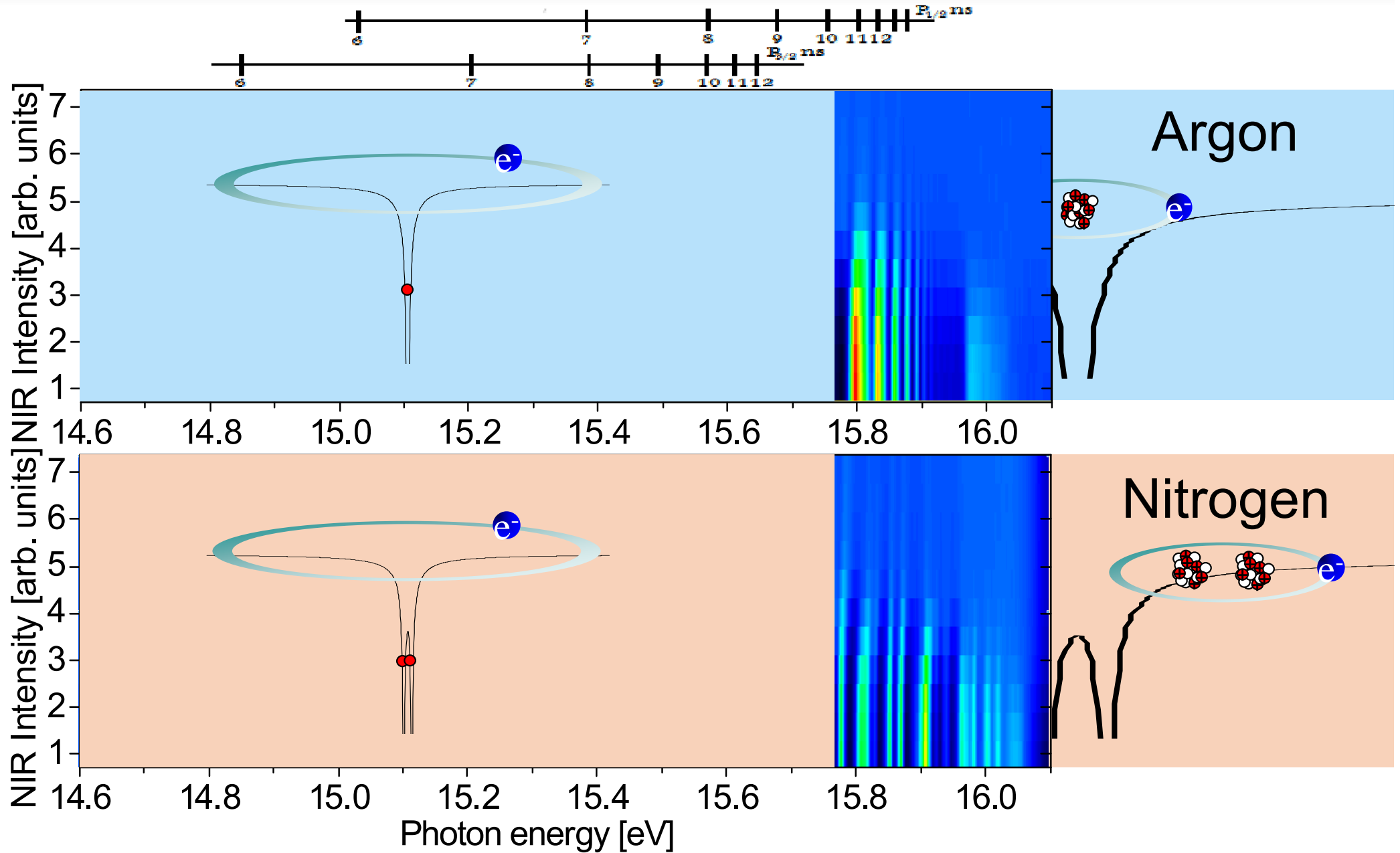
Science **354** 738 (2016)

From atoms to molecules

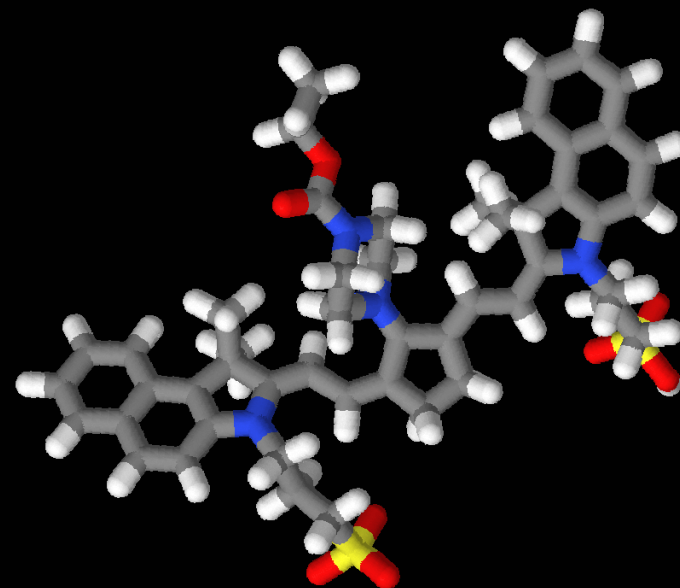
Can we change the shape
of molecular states?

If such control works for (two) bonding electrons in molecules ...
... this would open doors to laser-directed chemistry.

Atomic and molecular resonances interacting with weak to strong laser fields



Can we change the shape of complex molecules?

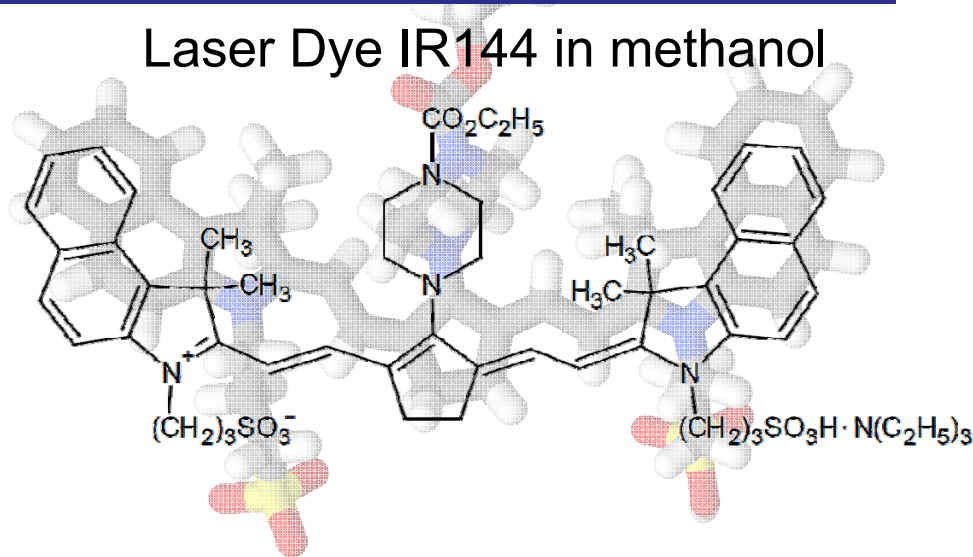


Fano control of molecules in the liquid phase

cooperation with J.-M. Mewes, A. Dreuw, Univ. Heidelberg

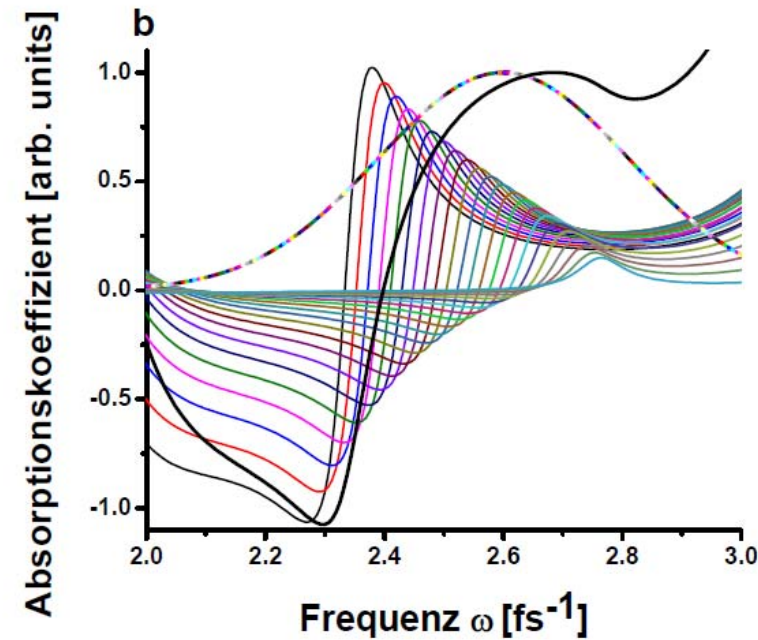
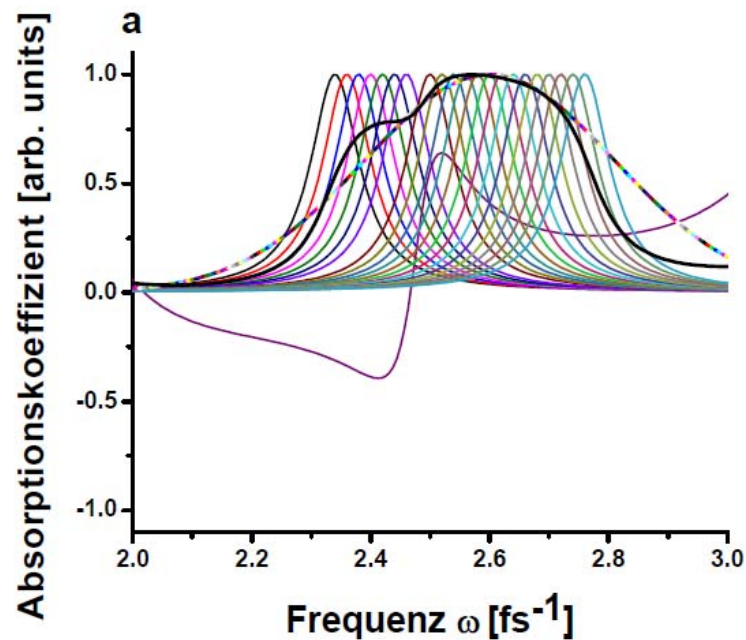
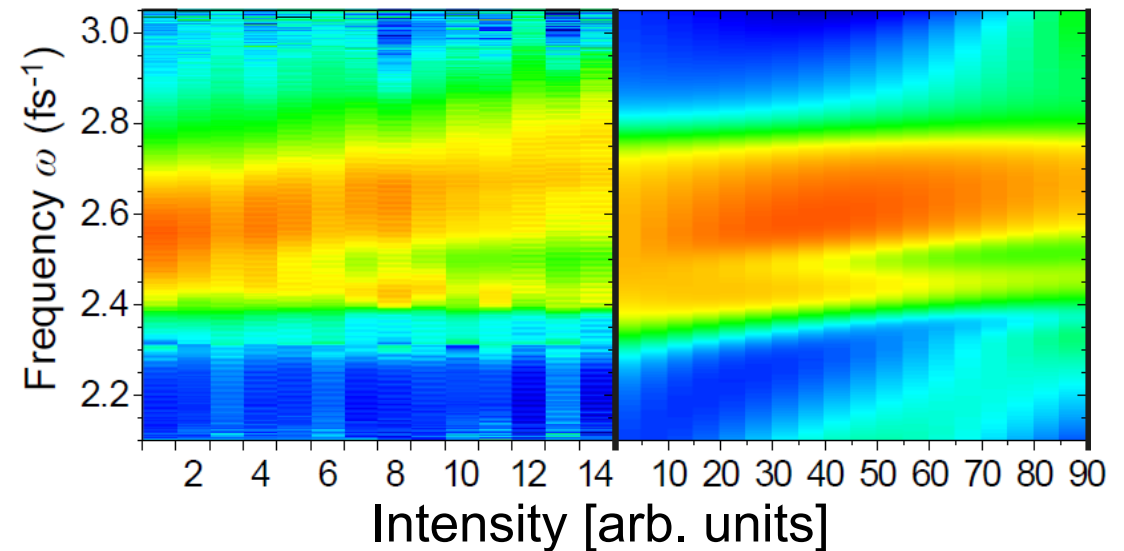
Kristina Meyer *et al.*, PNAS (2015)

Laser Dye IR144 in methanol



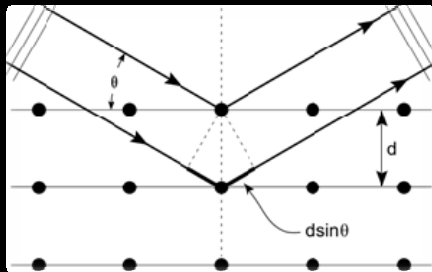
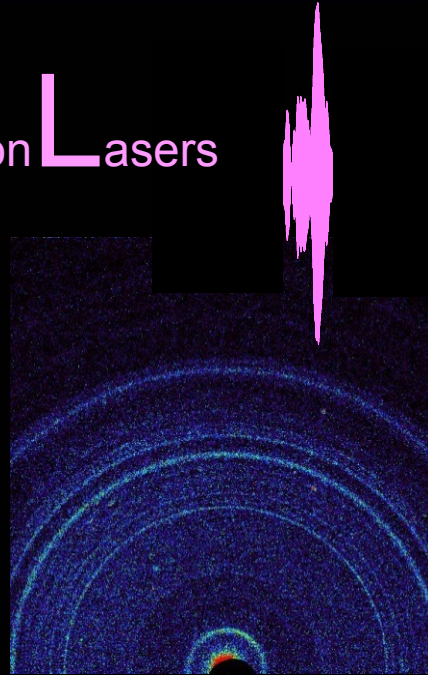
Measurement

Simulation



Time-resolved Science with novel x-ray/XUV laser sources

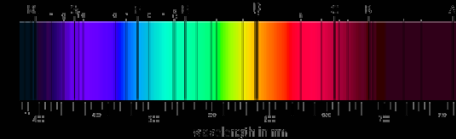
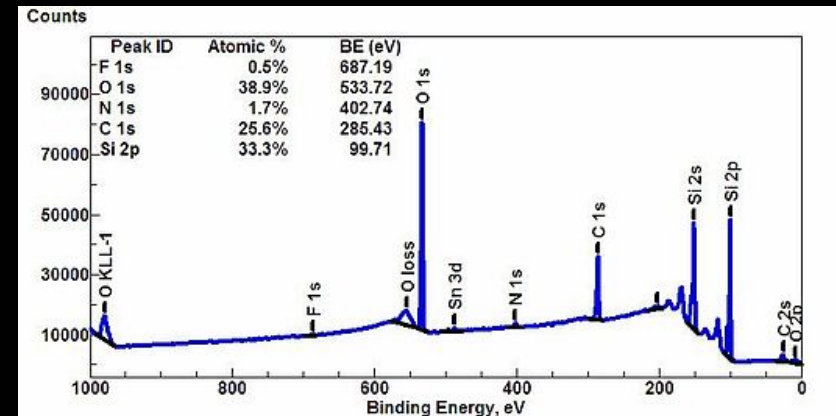
Free **E**lectron **L**asers



X-ray diffraction

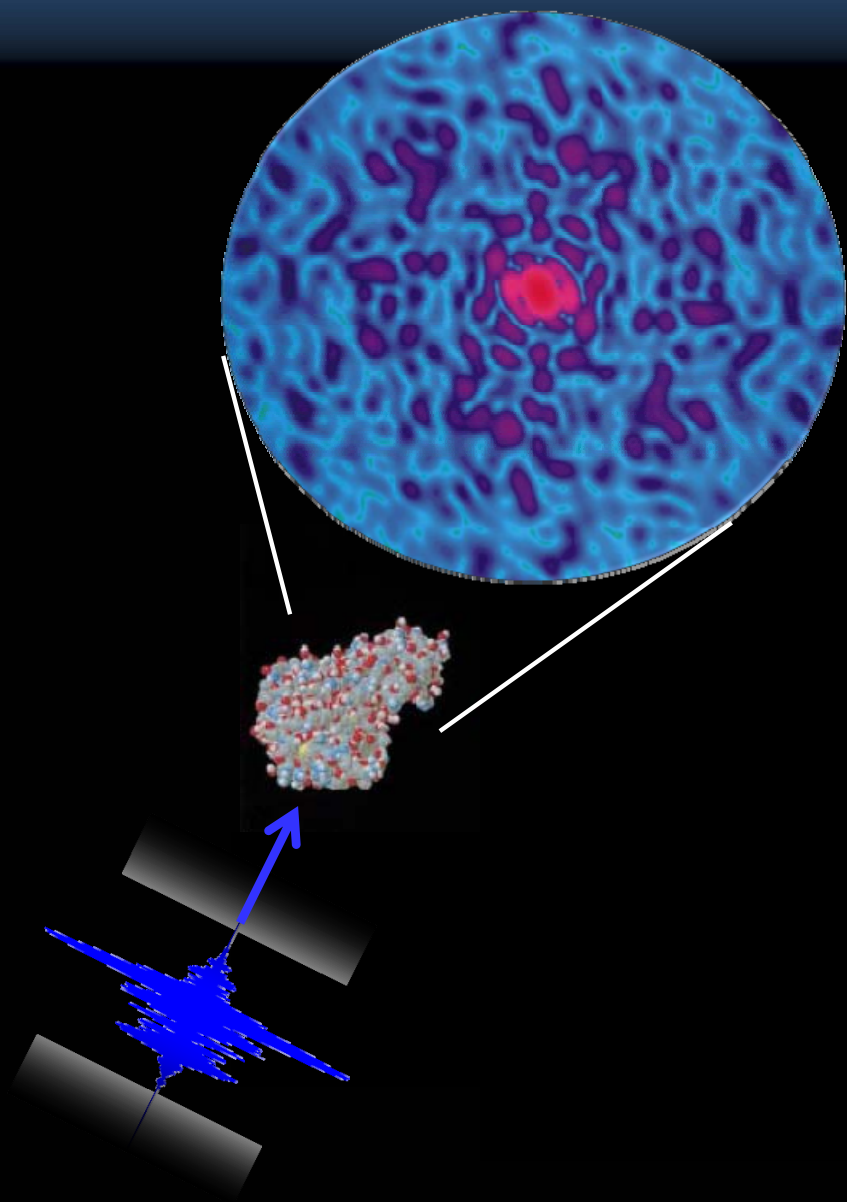


High **H**armonic **G**eneration

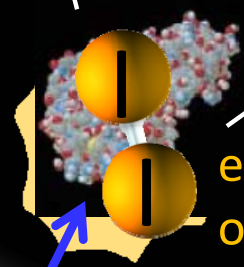
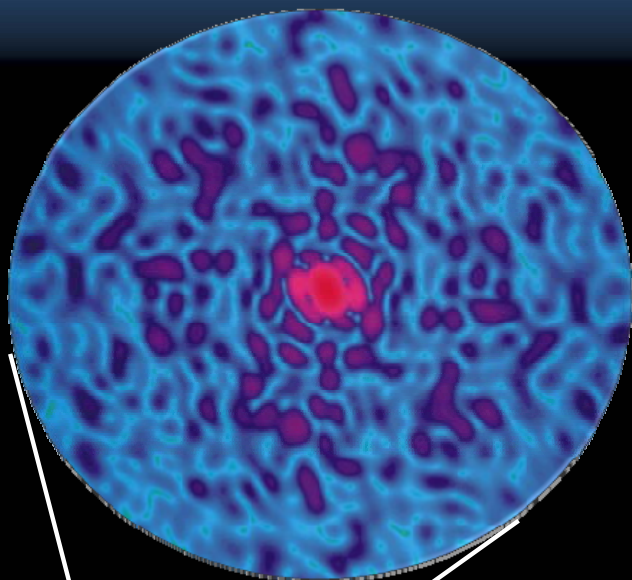


XUV spectroscopy

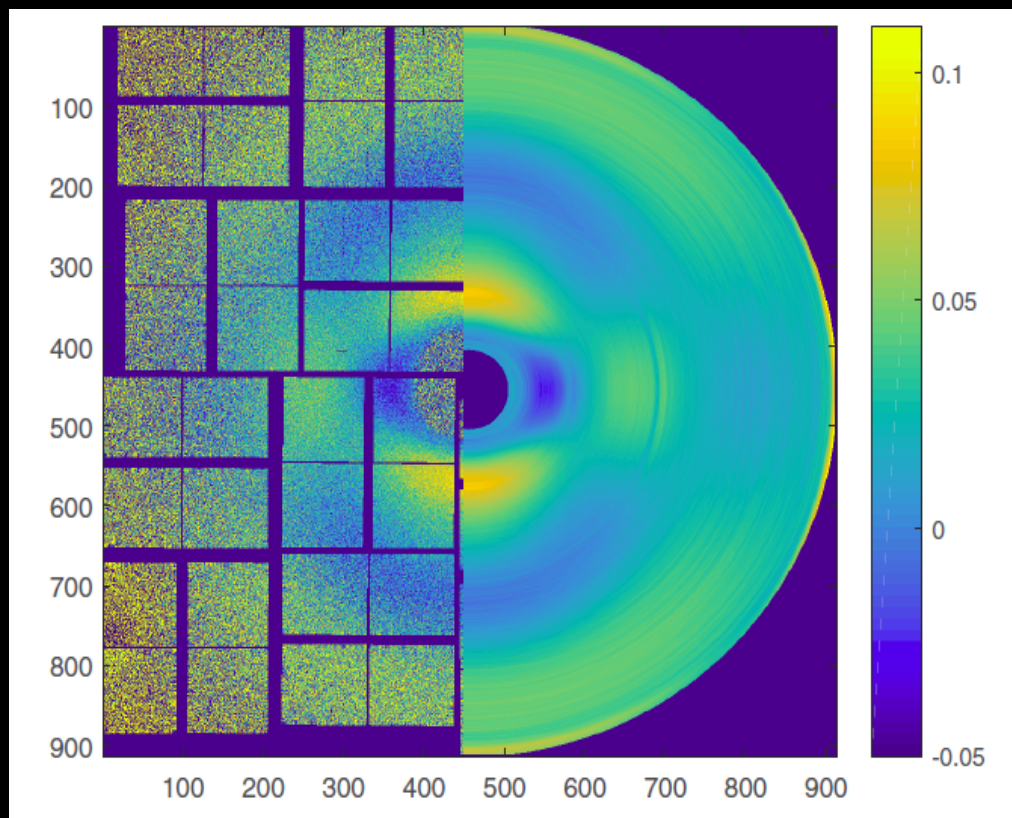
Challenge: “See” (image) molecules in *x-ray* Light



One Goal: “See” (image) molecules in *x-ray* Light



ensemble
of gas-phase
iodine (I₂)
molecules

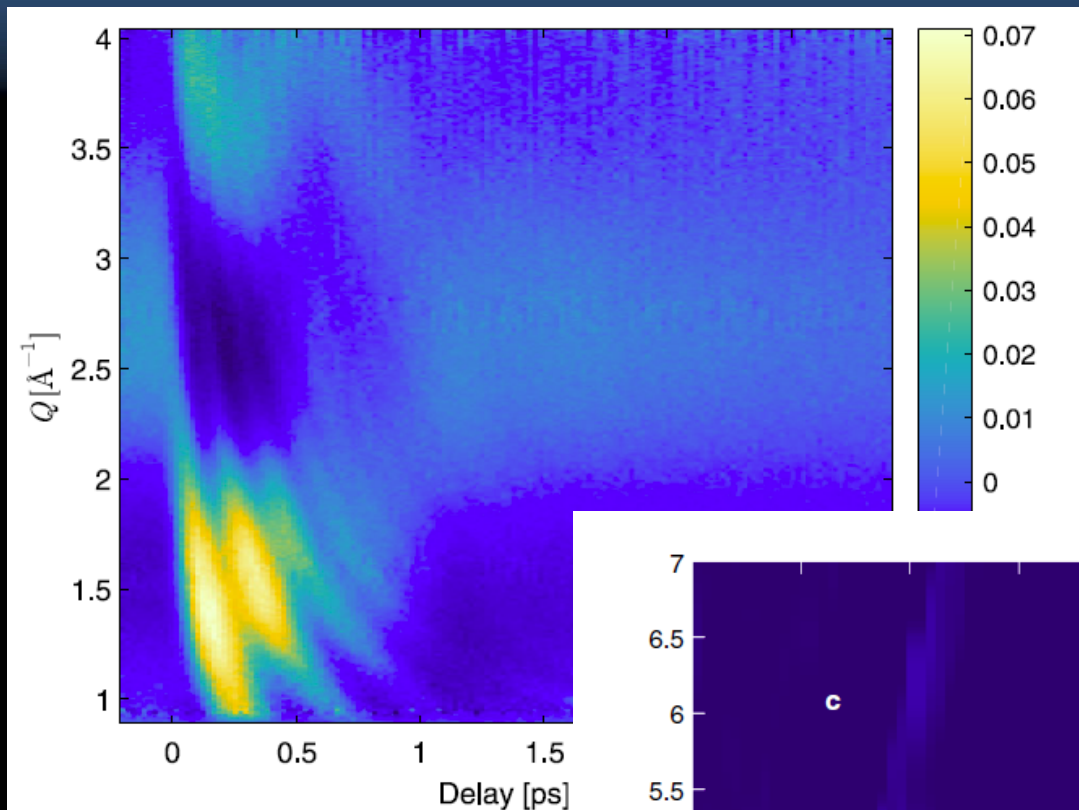


Glownia et al. PRL **117**, 153003 (2016)
(Bucksbaum group@SLAC)

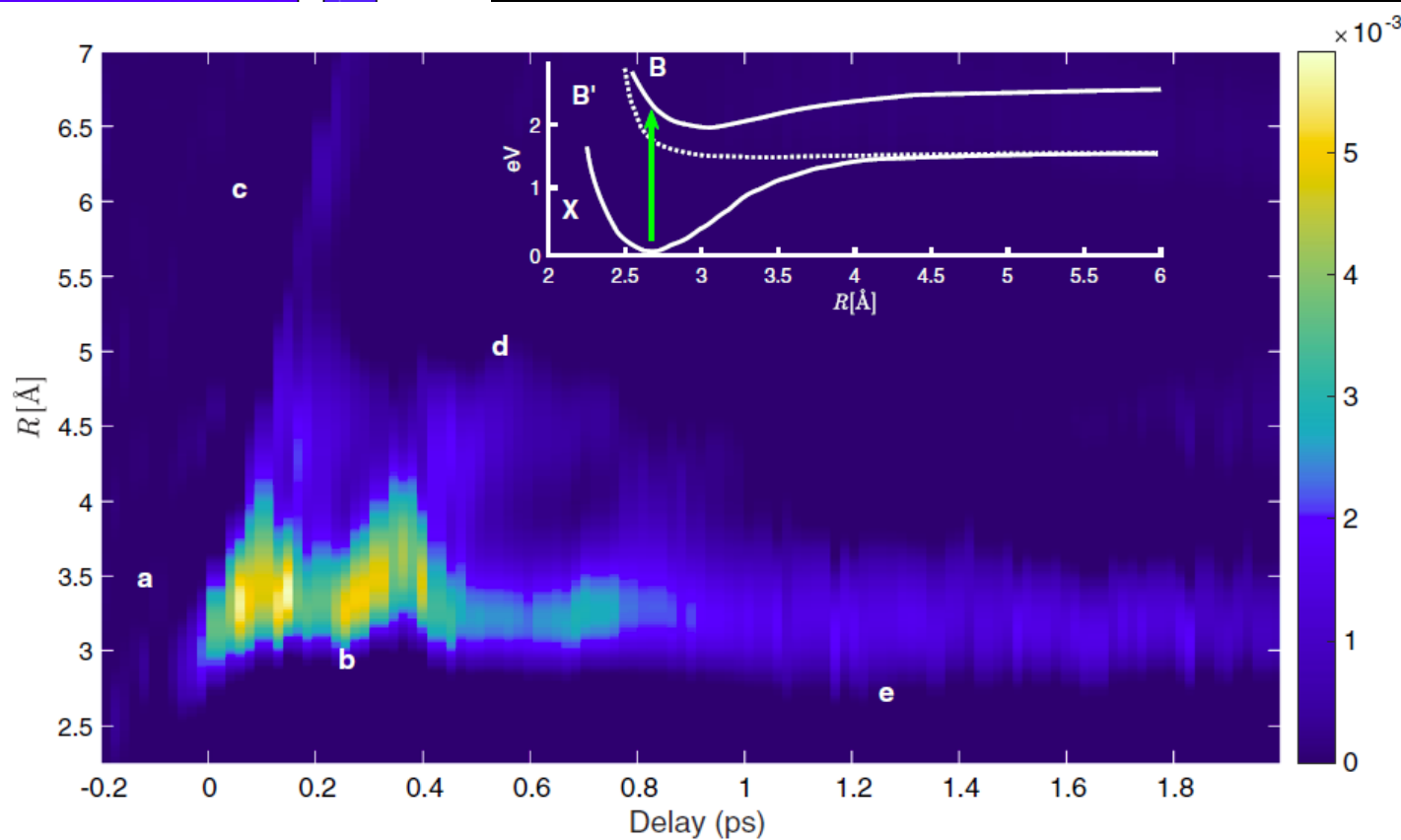
Time-resolved dynamics

Scattering image acquired for various delay times

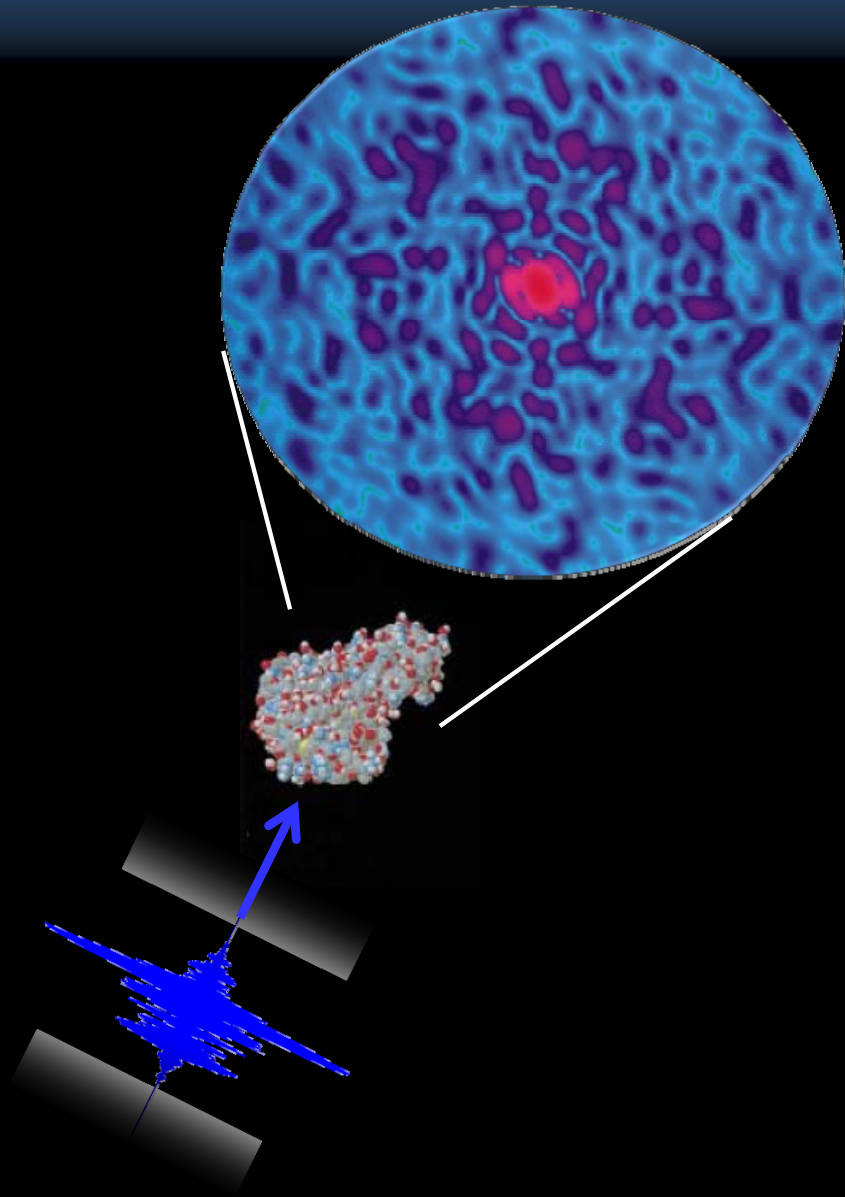
Glowia et al. PRL **117**, 153003 (2016)
(Bucksbaum group@SLAC)



Reconstructed molecular wavepacket after single-photon excitation

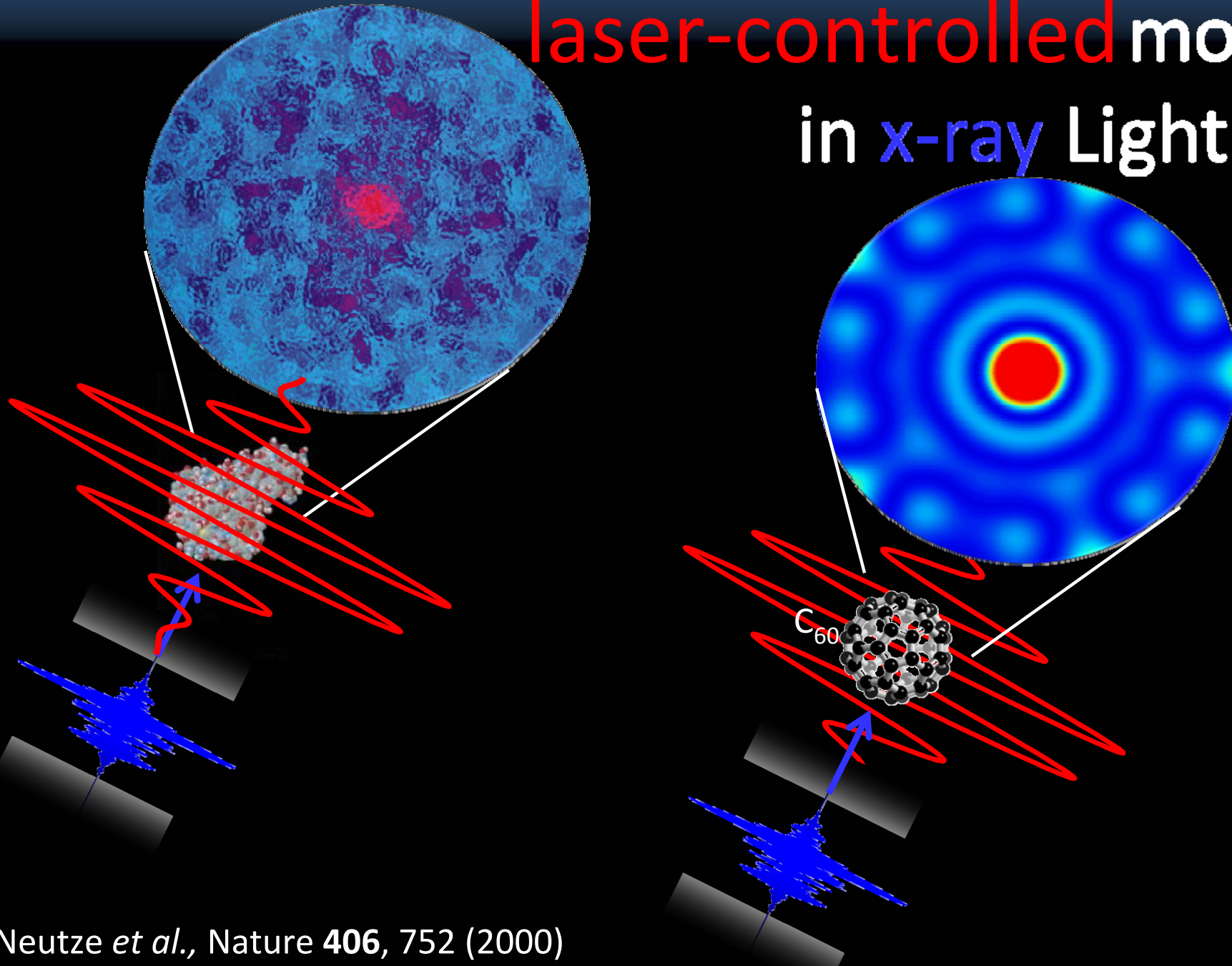


One Goal: “See” (image) molecules in *x-ray* Light



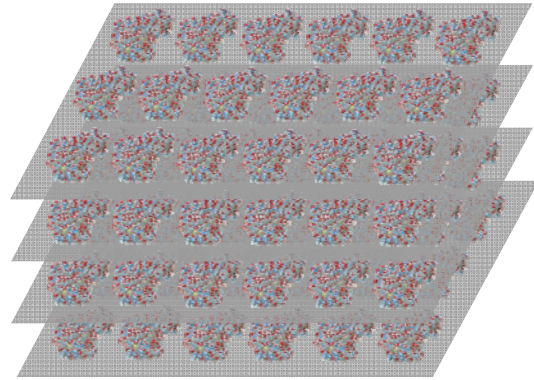
Next Step: "See" (image)

laser-controlled molecules
in x-ray Light



X-ray diffraction crash course/refresher

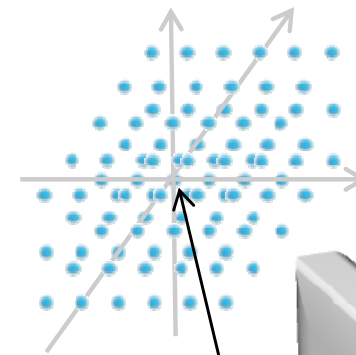
real space



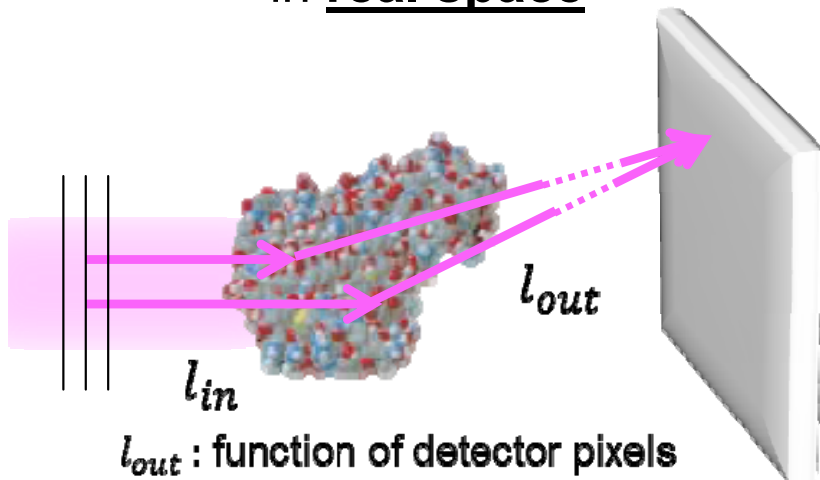
Fourier Transform



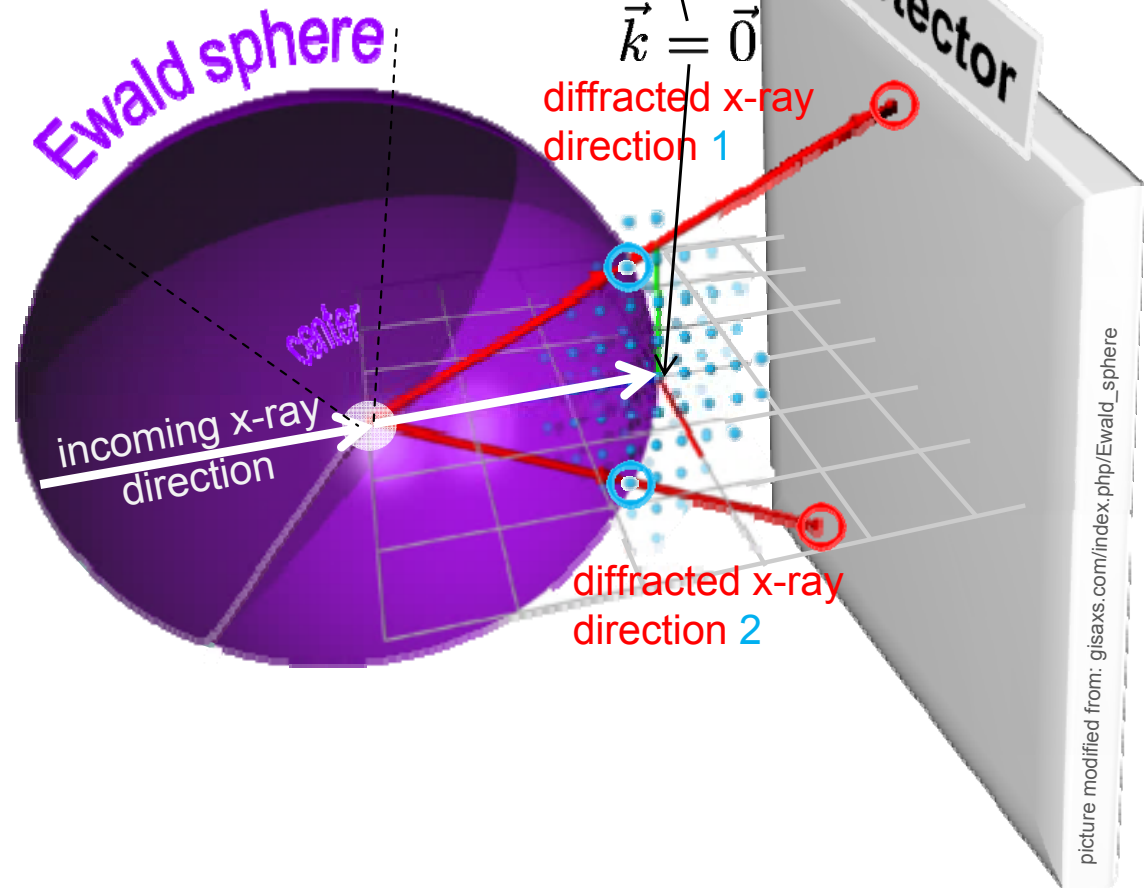
k/momentum space



Alternative route for calculation:
-> addition of geometric paths
in **real space**



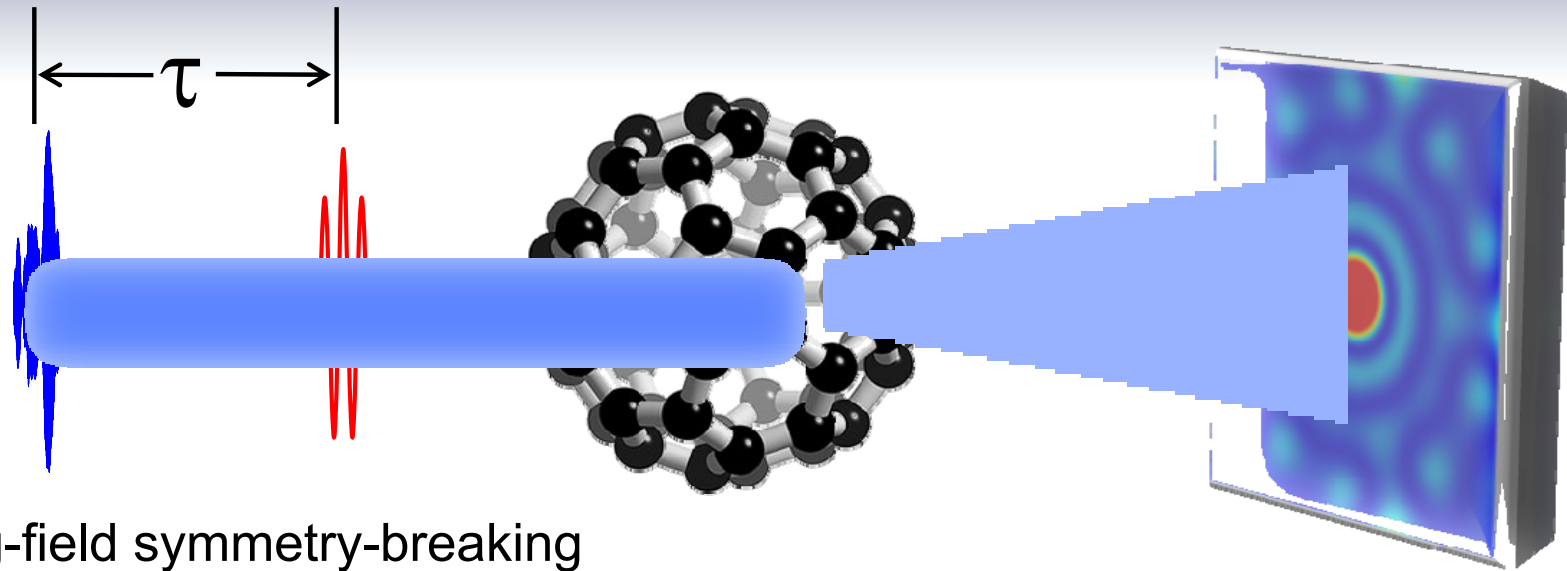
l_{out} : function of detector pixels
and atomic positions x, y, z
(simplified by letting $l_{out} \gg \Delta x, \Delta y, \Delta z$)



picture modified from: gisaxs.com/index.php/Ewald_sphere

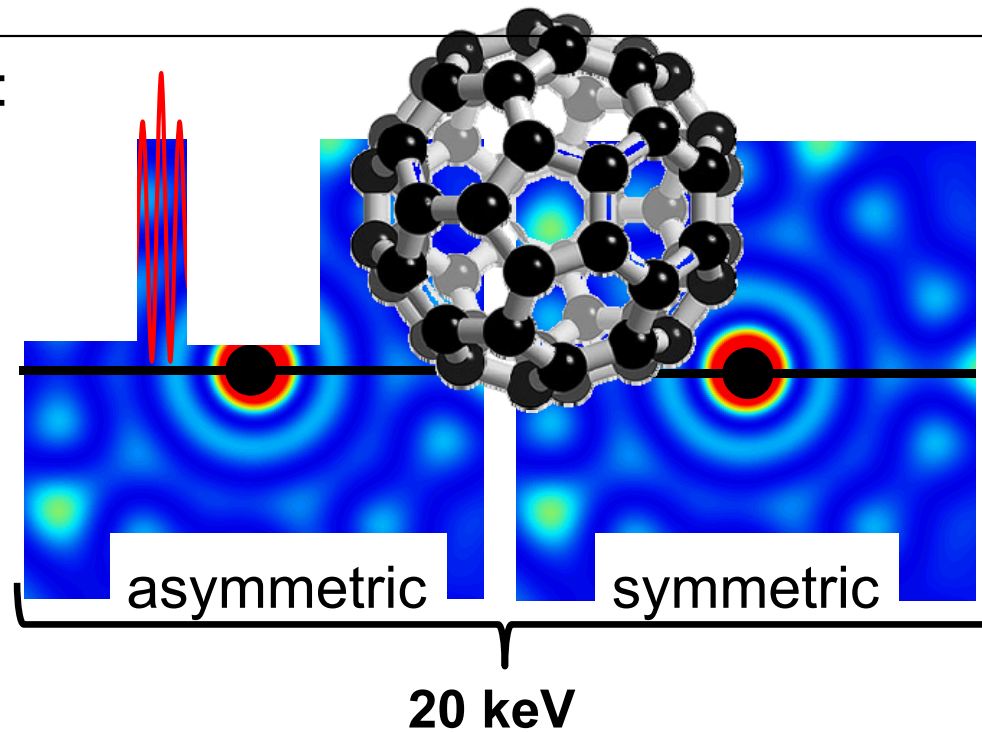
C_{60} vibrational dynamics imaging

C_{60} real-space image
from: 3DChem.com



study strong-field symmetry-breaking

simulation results:



LCLS Proposal

(1st submission 2012)

X-ray imaging of laser-induced coherent dynamics in C₆₀ fullerenes

Claus Peter Schulz*, Ilme Schlichting, Lutz Foucar, Thomas Möller, Christoph Bostedt, Timur Osipov, Arnaud Rouzée, Marc Vrakking, Artem Rudenko, Katharina Kubicek, Simone Techert, Nora Berrah, Jochen Küpper, Ulf Saalman, Jan Michael Rost, Rüdiger Schmidt, Kiyoshi Ueda, Louis DiMauro, Robert Moshhammer, Joachim Ullrich, and Thomas Pfeifer*

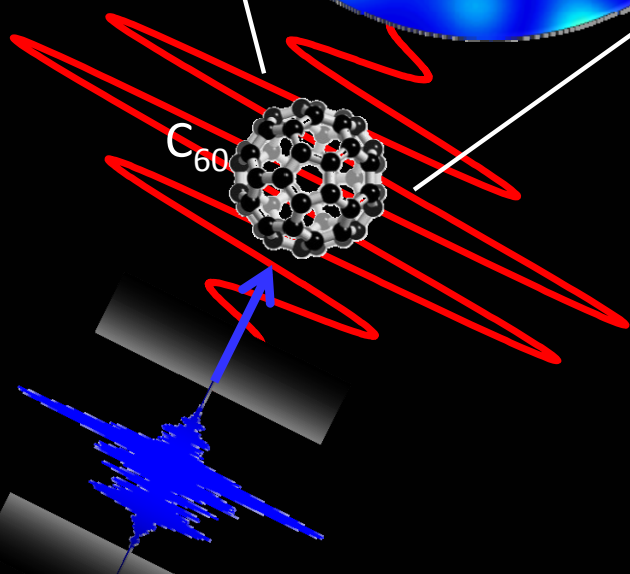
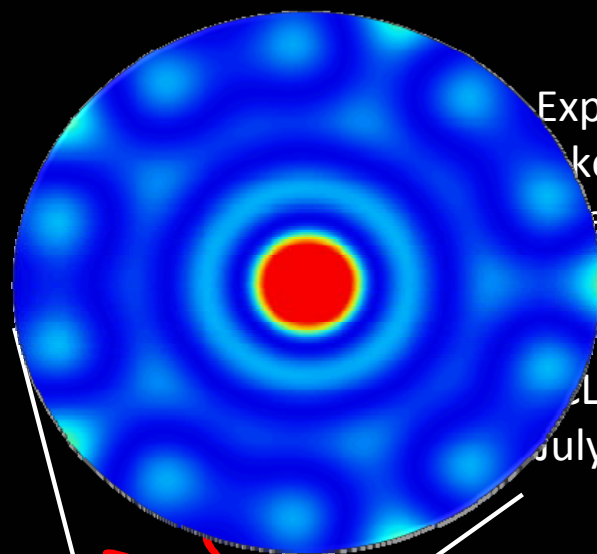
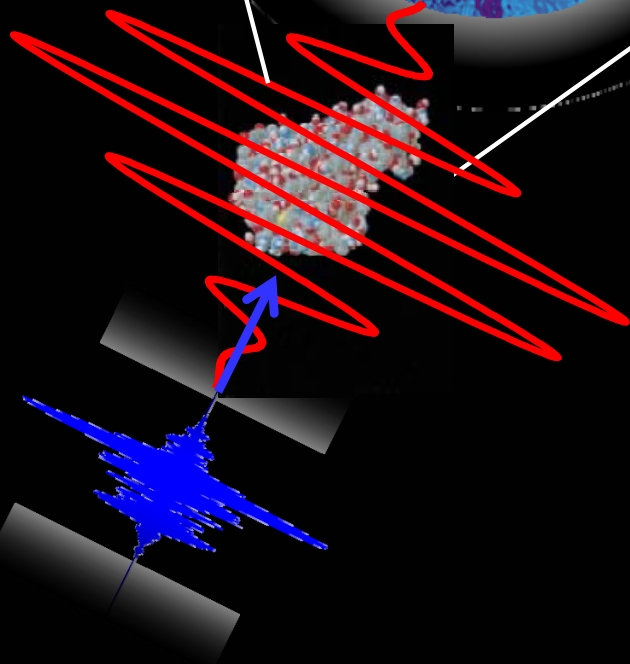
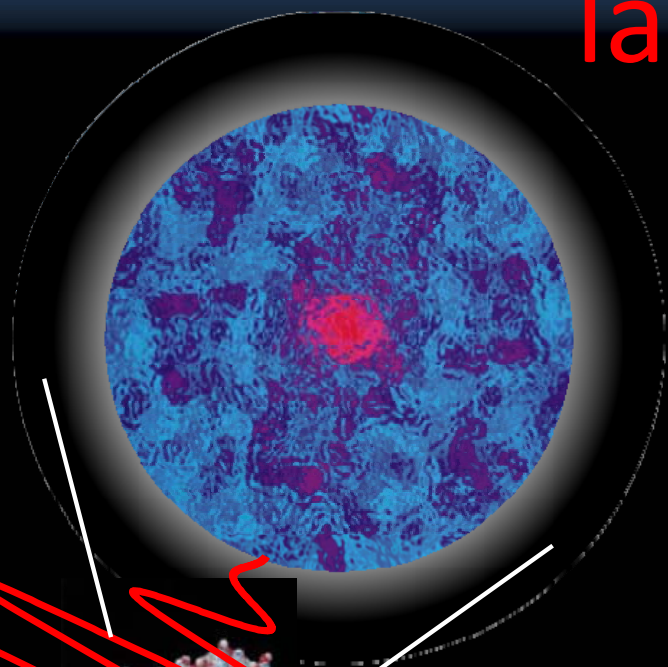
*spokespersons: cps@mbi-berlin.de, tpfeifer@mpi-hd.mpg.de

Abstract: Optical lasers interacting with C₆₀ fullerene molecules induce coherent electronic and vibrational excitations, resulting in the deformation of the originally symmetric structure. The strong-field symmetry-breaking (leading to asymmetric prolate/oblate deformation vs symmetric breathing mode) and the corresponding coherence lifetimes have only been indirectly explored in experiments, while direct imaging could provide answers to these scientific questions. This experiment can only be conducted at the LCLS due to its specific attributes as explained in the experimental section. C₆₀ represents an important model system at the interface between clusters and molecules, and also between inorganic and organic medium-size molecules. The results of this proposed experiment are also expected to have benchmark character for future studies in dynamical x-ray diffractive imaging of molecules.

“Can we >see<  C₆₀ changing shape in a strong field”

One Goal: "See" (image)

laser-controlled molecules
in **x-ray** Light



Experimental Data
keV photon
energy
reactive imaging
XLS Experiment
July 2016

C₆₀ dynamical transitions

delay → IR early

delay
↓
IR early

IR pulse energy: 2.57 mJ

0.76 mJ

0.35 mJ

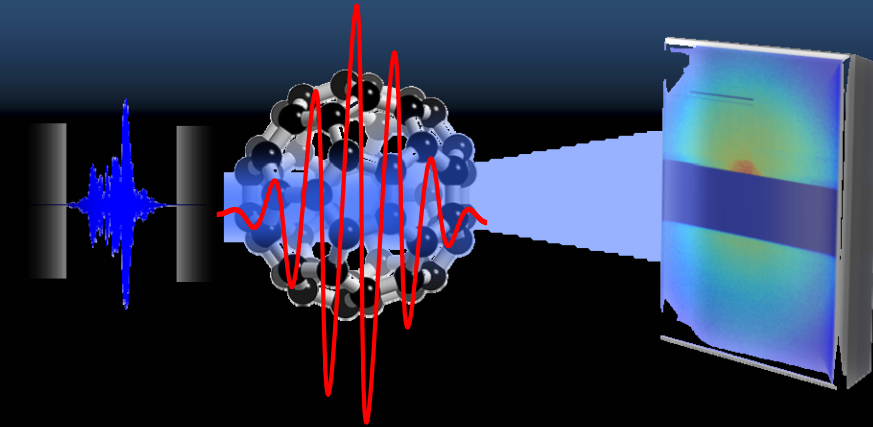
Imaging time-resolved laser-induced symmetry breaking?

horizontal width
vertical width

total diffraction signal

IR-Laser polarization: horizontal

Opening new science opportunities



of particular interest:
- for LCLS II
- for European XFEL
(higher rep rates)

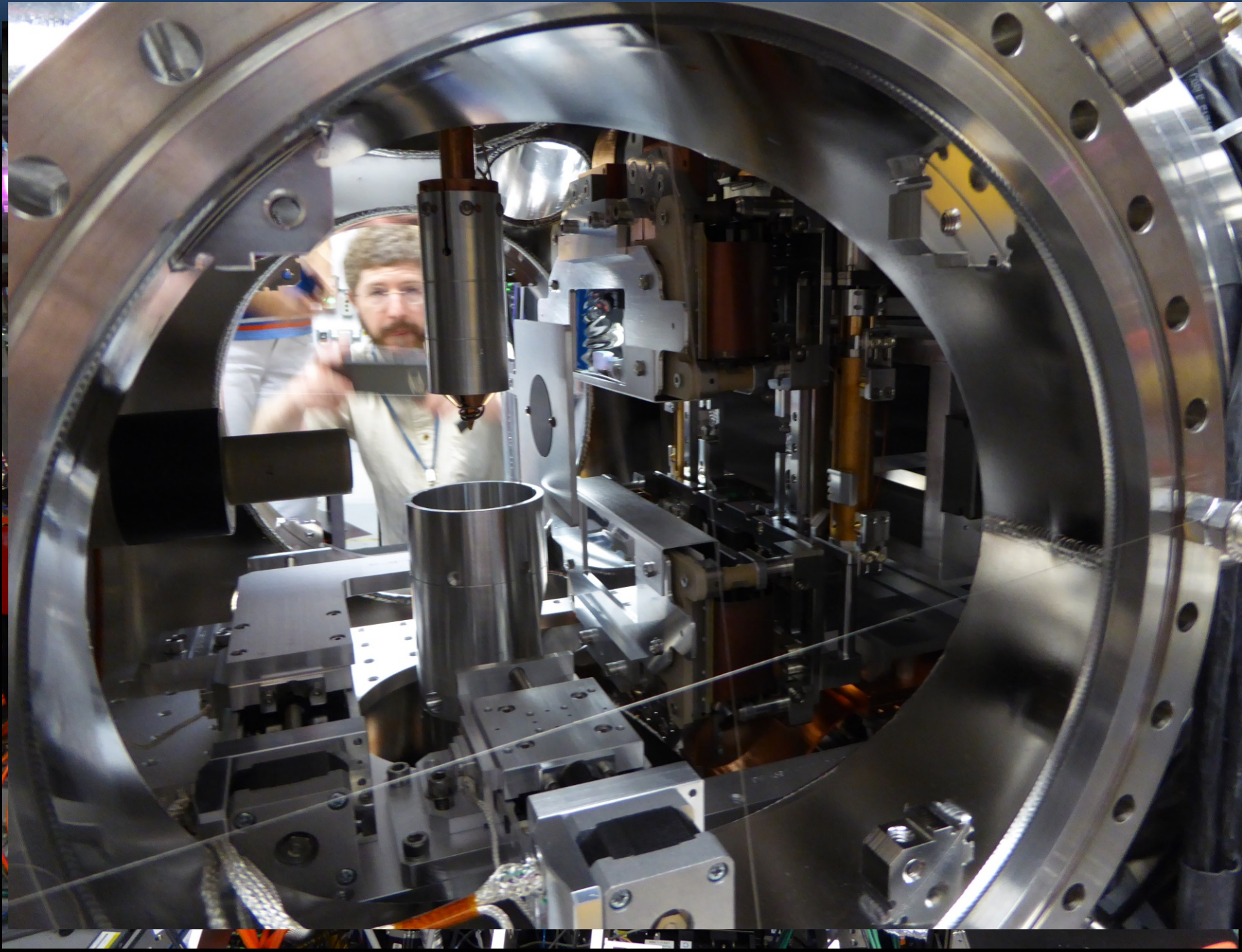
- femtosecond time-resolved x-ray imaging at $10^{11} / \text{cm}^3$ gas-phase (\sim molecular jet) densities is **possible** (for **ensemble** measurements)
- in C_{60} : observation of different **dynamical regimes** depending on (optical) laser **intensity**

high
intensity
 $10^{15} \frac{\text{W}}{\text{cm}^2}$

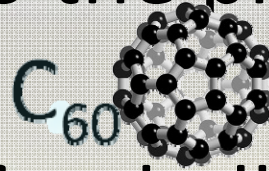
low
intensity
 $10^{14} \frac{\text{W}}{\text{cm}^2}$

Can we steer, and watch, **laser-driven chemistry** in real time?

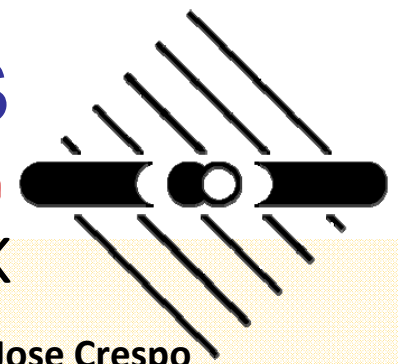
C_{60} imaging: AMO@LCLS



Thanks to the players of the
quantum football/soccer team



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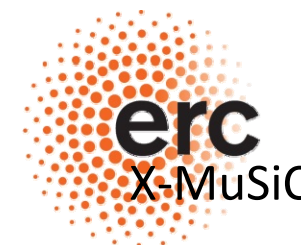
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Heidelberg Center for
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Take-home messages

Atoms and molecules **change** their spectral (line) shape

as a function of **laser intensity**

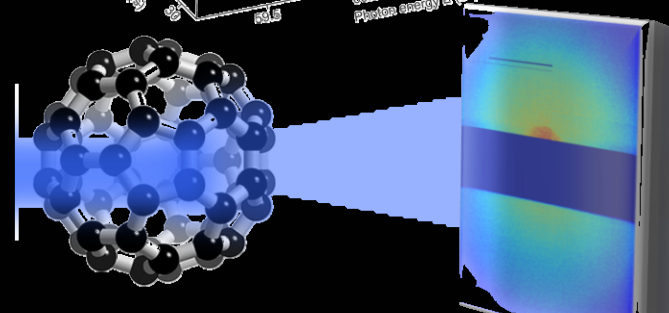
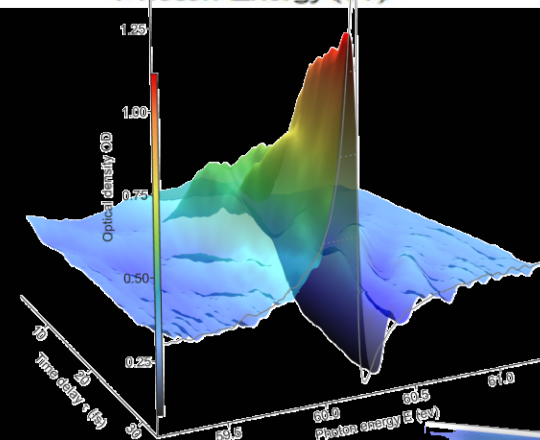
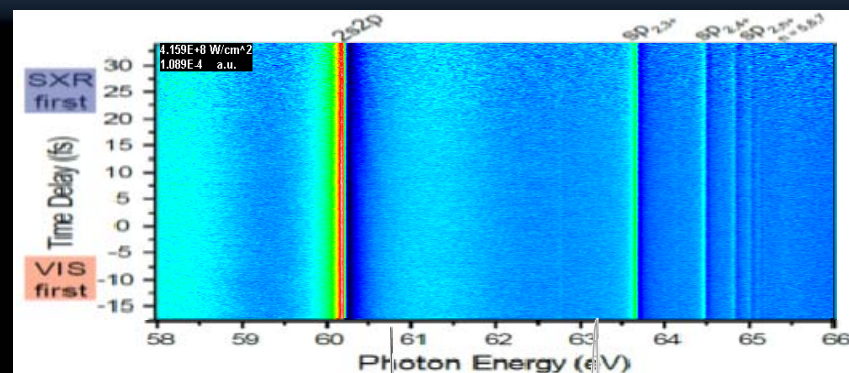
quantum-state
phase $\varphi = \Delta E \Delta t$

$$q(\varphi) = -\cot\left(\frac{\varphi}{2}\right)$$

- **Phase** control tunes the Fano q parameter
- **Amplitude** gate resolves the Fano resonance

... and these mechanisms are general...

- **See/image** molecules **changing** their spatial shape
Intensity dependence of structural dynamics in C_{60}



Understand & Control Matter in Strong laser fields